

chemistry meets topology

what oxidation numbers are all about

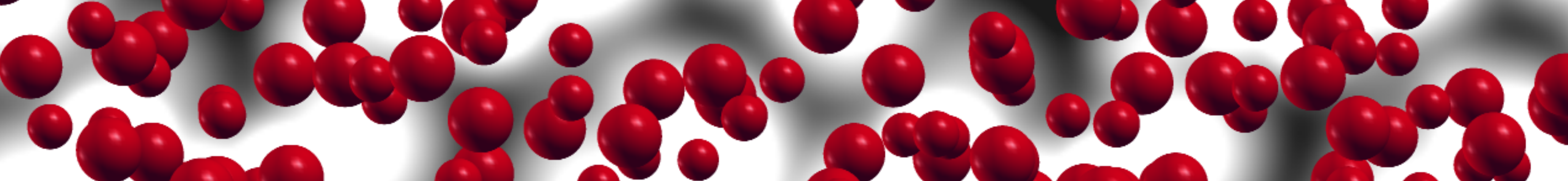
Stefano Baroni

Scuola Internazionale Superiore di Studi Avanzati
Trieste — Italy

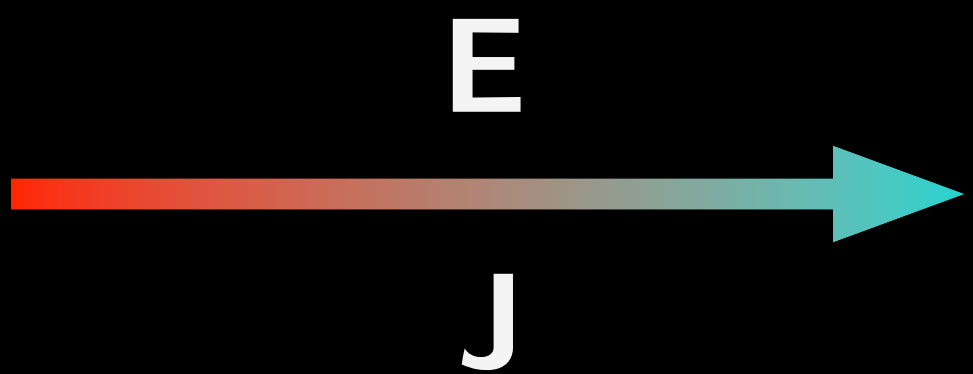


a prequel



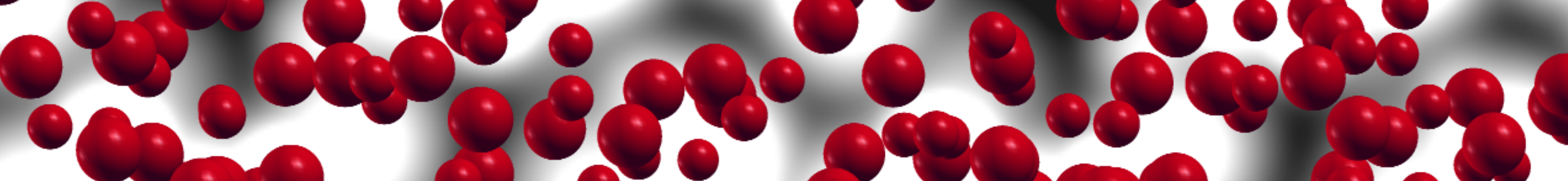


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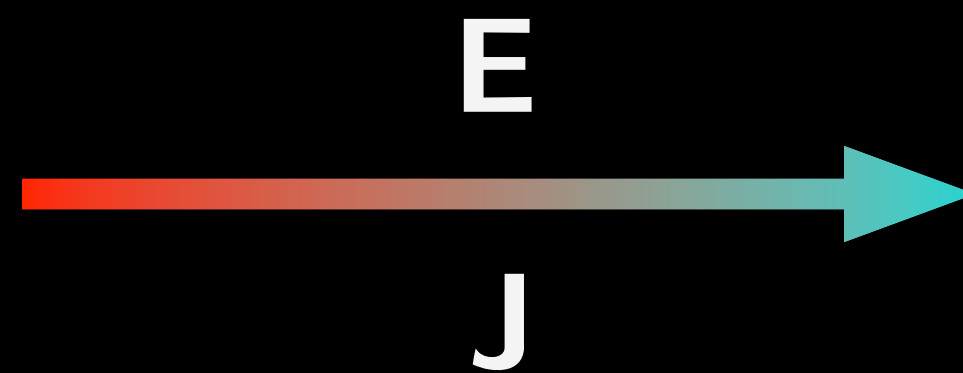
$$J = \sigma E$$



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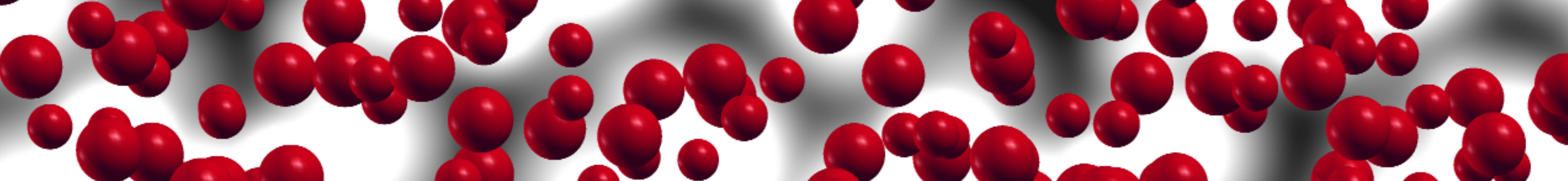
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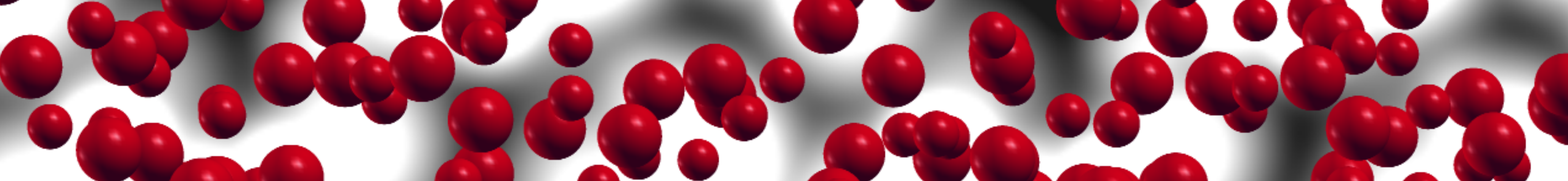
$$\mathbf{J} = \sigma \mathbf{E}$$

$$\mathbf{J} = \frac{1}{\Omega} \dot{\boldsymbol{\mu}}$$



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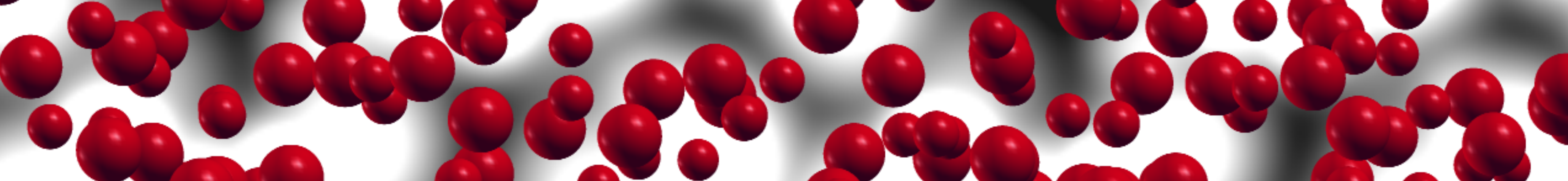
$$\begin{aligned} \mathbf{J} &= \frac{1}{\Omega} \dot{\boldsymbol{\mu}} \\ &= \frac{1}{\Omega} \sum_i \mathbf{z}_i^* \cdot \mathbf{v}_i \end{aligned}$$



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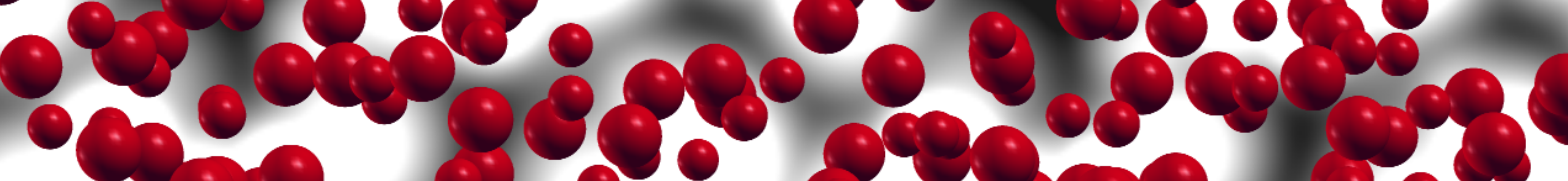
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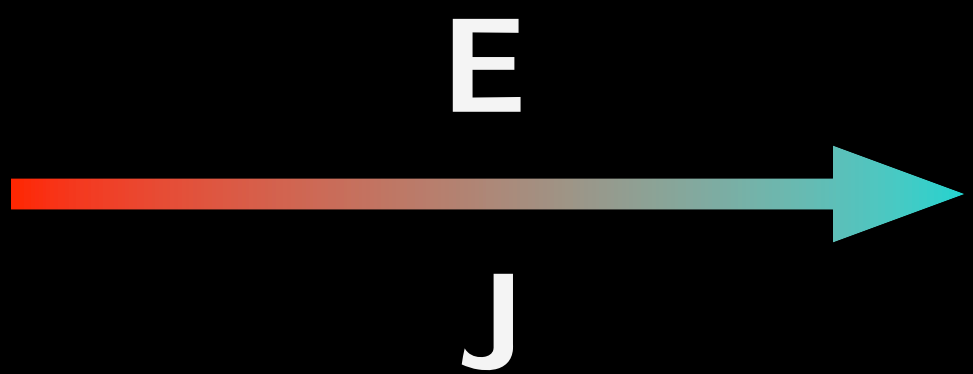
$z_{i\alpha\beta}^* = \frac{\partial \mu_\alpha}{\partial u_{i\beta}}$

$$\sigma = \frac{\Omega}{3k_B T} \int_0^\infty \langle \mathbf{J}(t) \cdot \mathbf{J}(0) \rangle dt$$





+



-

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$$= \frac{\Omega}{3k_B T} \lim_{t \rightarrow \infty} \frac{1}{2t} \left\langle \left| \int_0^t \mathbf{J}(t') dt' \right|^2 \right\rangle$$

keep in mind!



a conundrum in transport theory

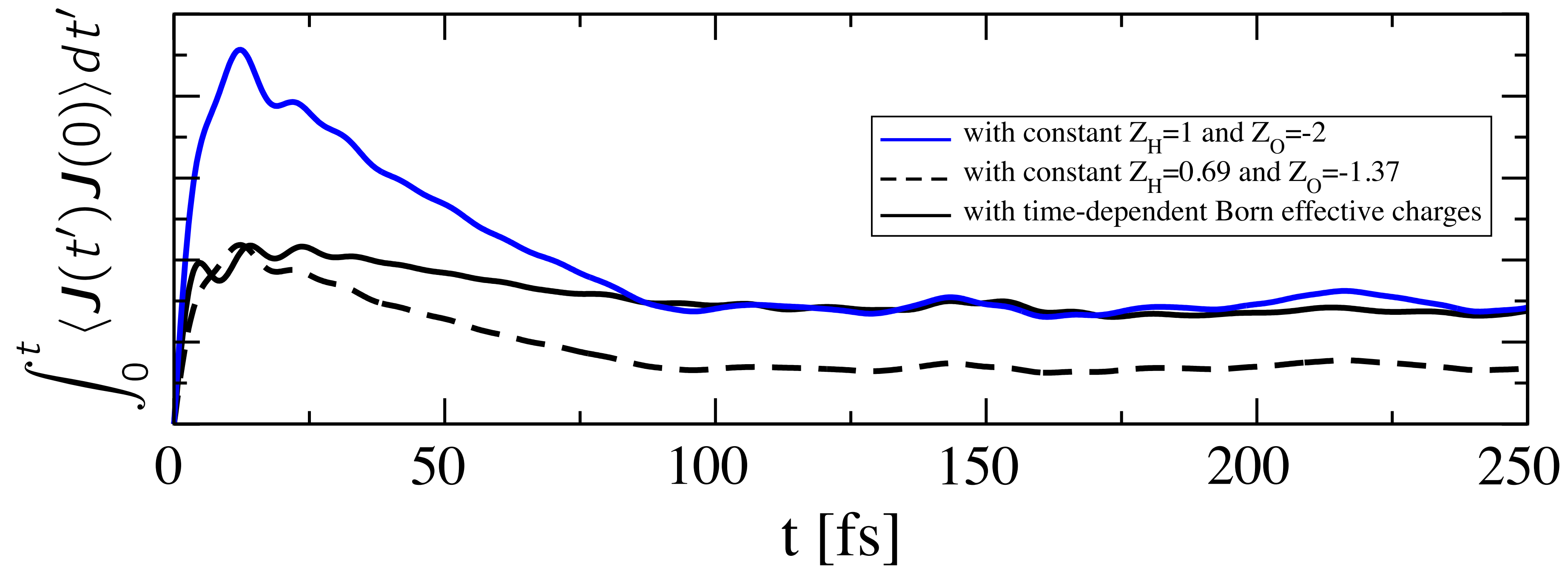
PRL 107, 185901 (2011)

PHYSICAL REVIEW LETTERS

week ending
28 OCTOBER 2011

Dynamical Screening and Ionic Conductivity in Water from *Ab Initio* Simulations

Martin French,¹ Sebastien Hamel,² and Ronald Redmer¹



a conundrum in transport theory

PRL 107, 185901 (2011)

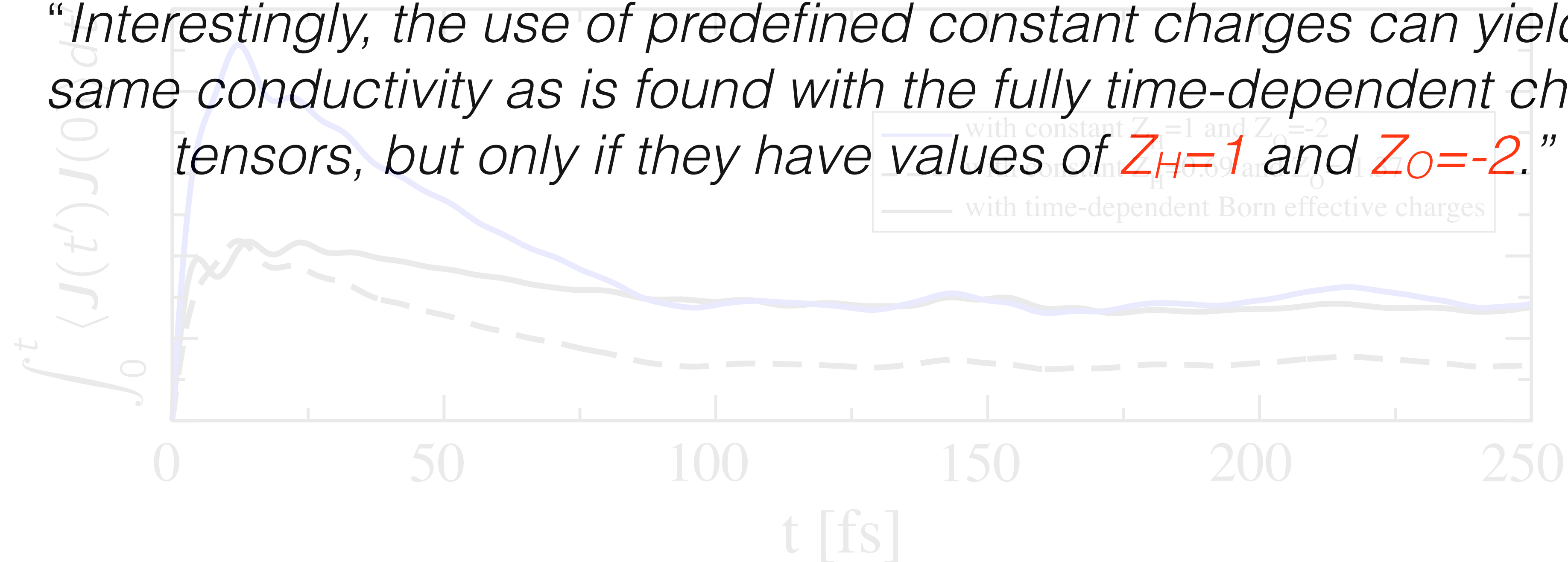
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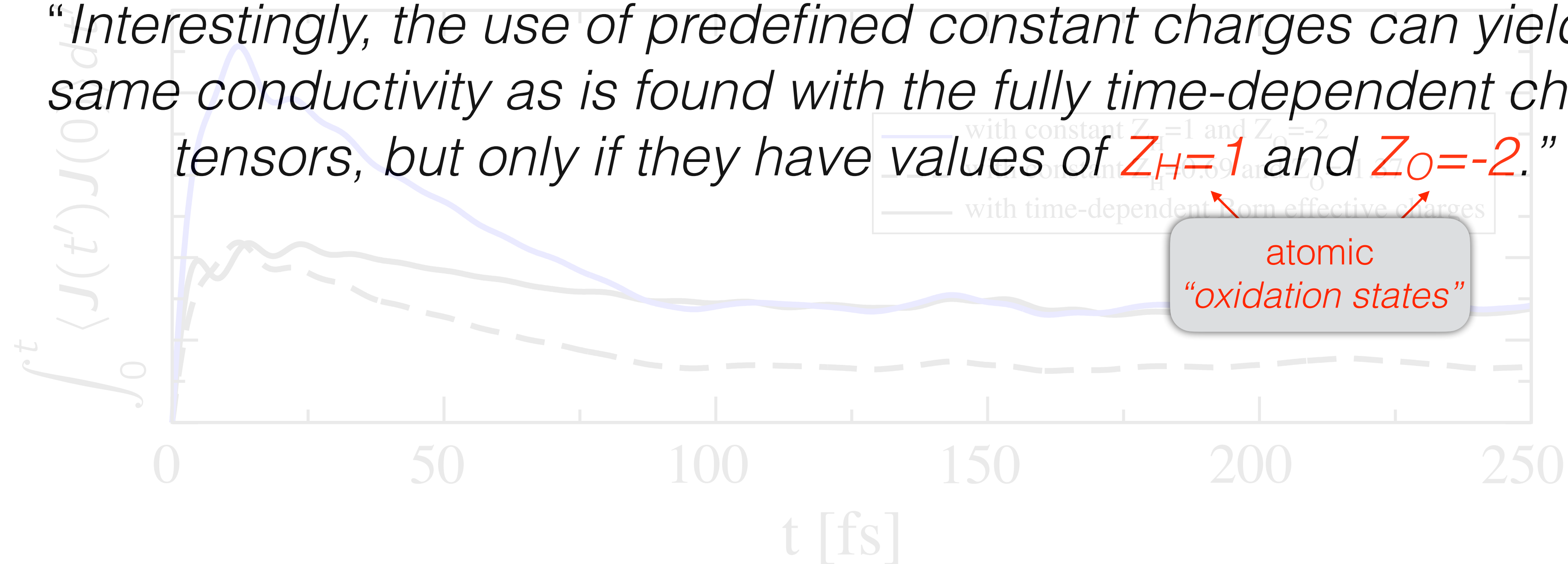
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... but what are
oxidation states,
in the first place?

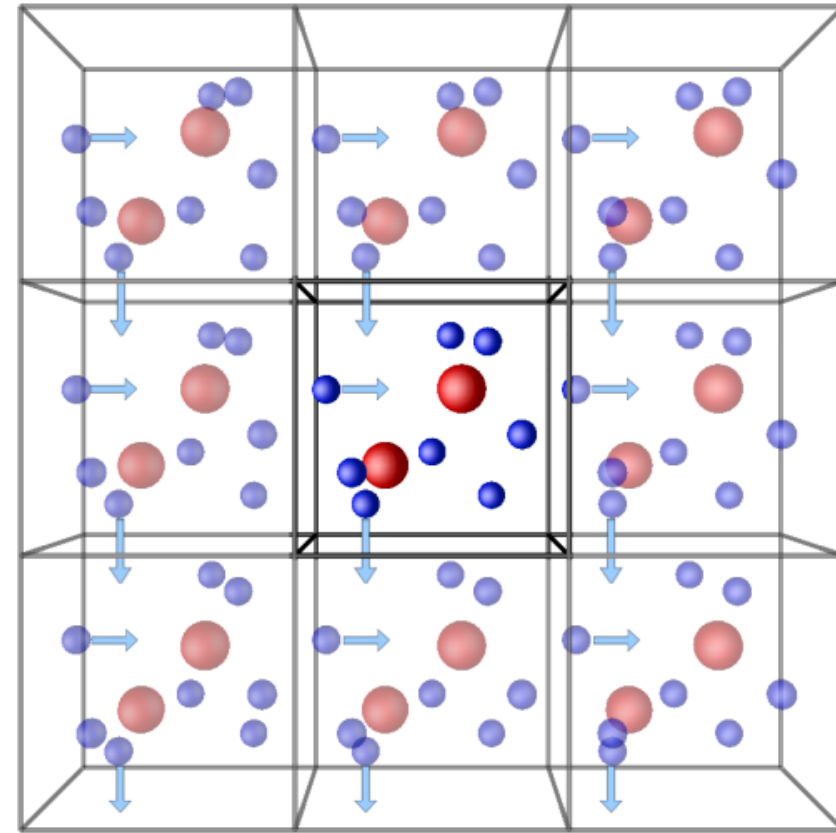


the oxidation state of an atom is the charge of this atom after ionic approximation of its heteronuclear bonds

<https://doi.org/10.1351/goldbook.O04365>

quantisation of adiabatic particle transport

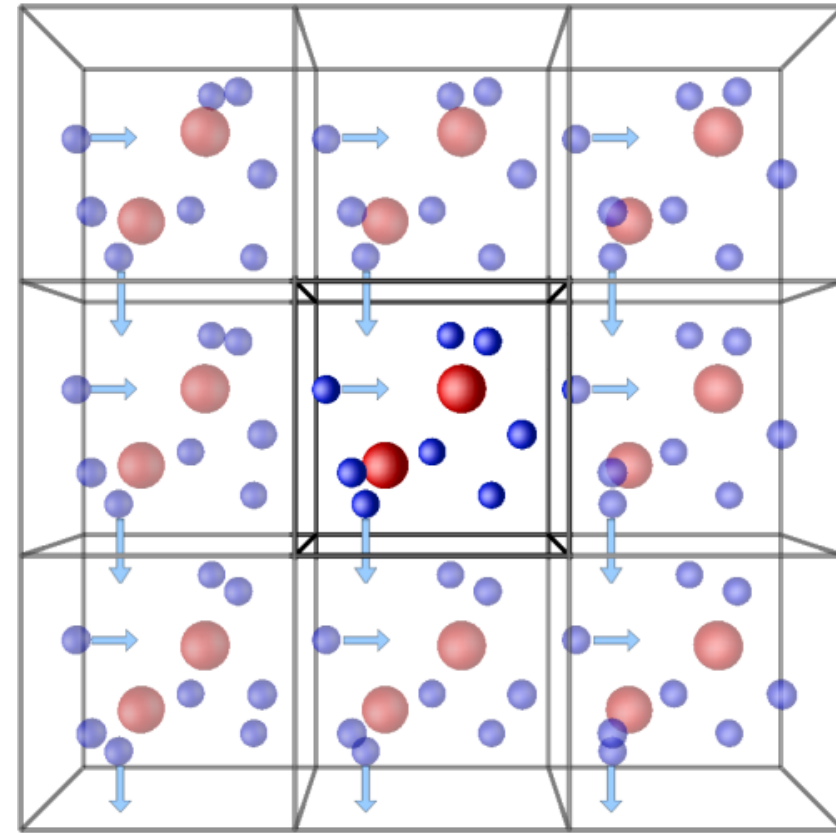
classical PBC



$$V(x + L) = V(x)$$

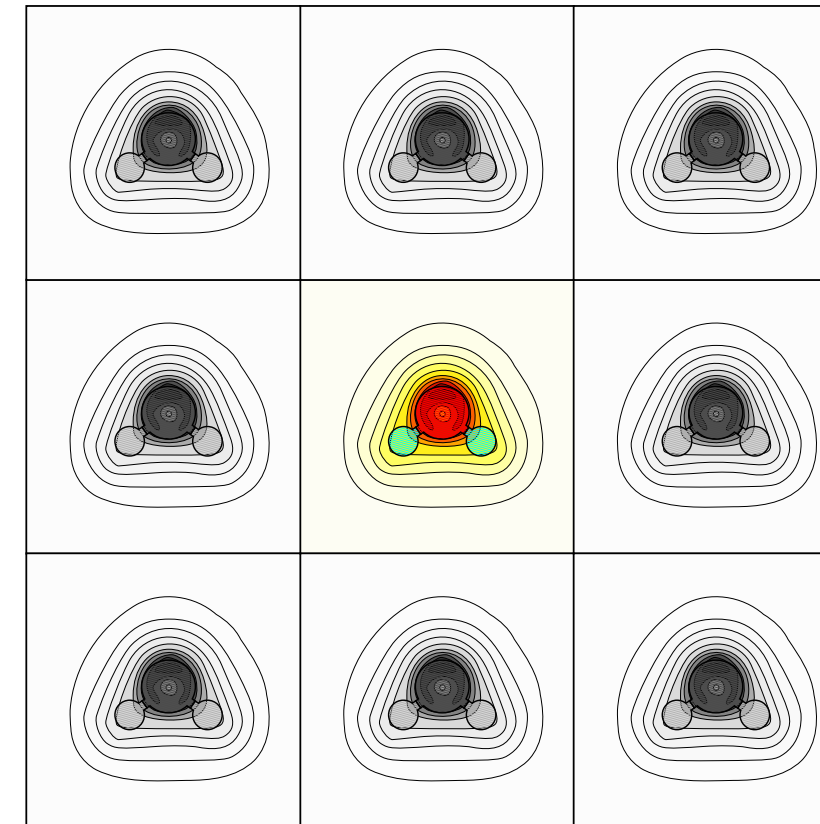
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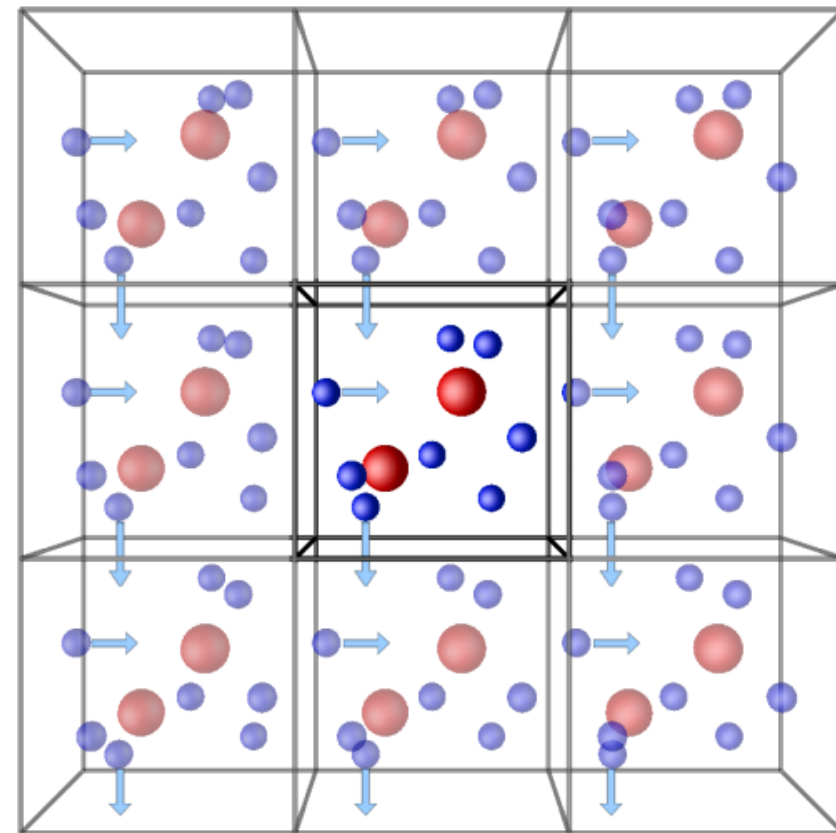
quantum PBC



$$\psi(x + L) = \psi(x)$$

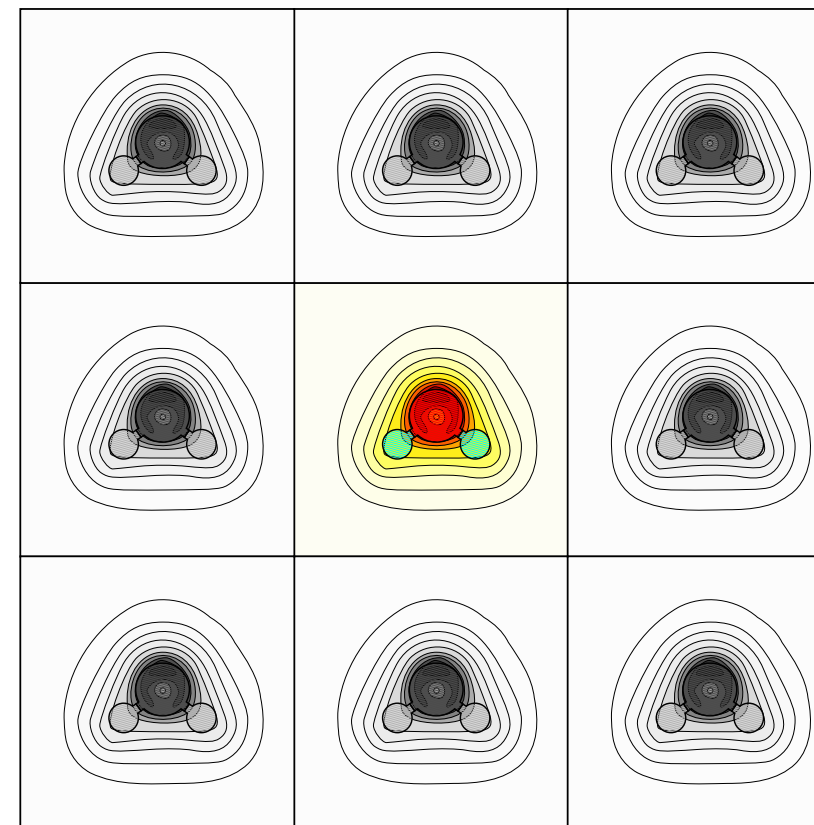
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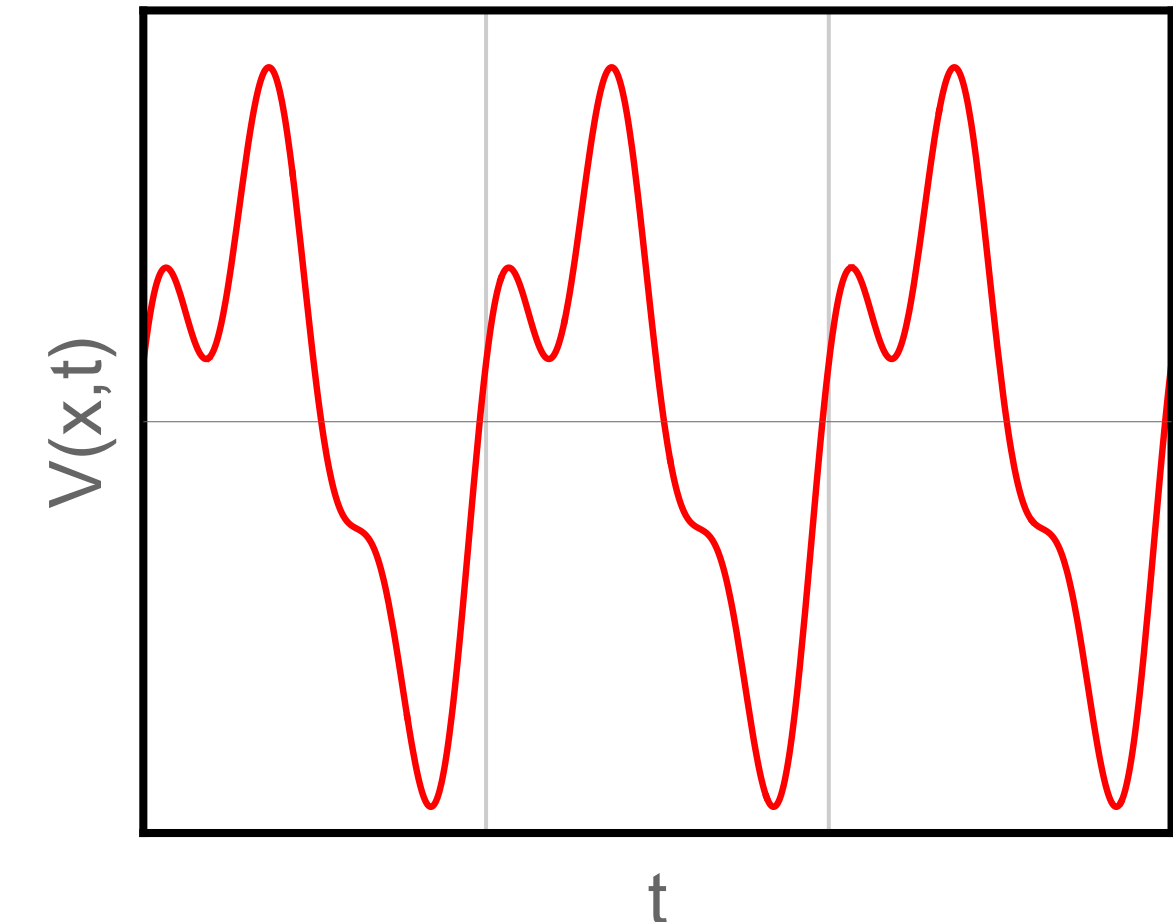
$$V(x + L) = V(x)$$

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$$\psi(x + L) = \psi(x)$$

time periodicity



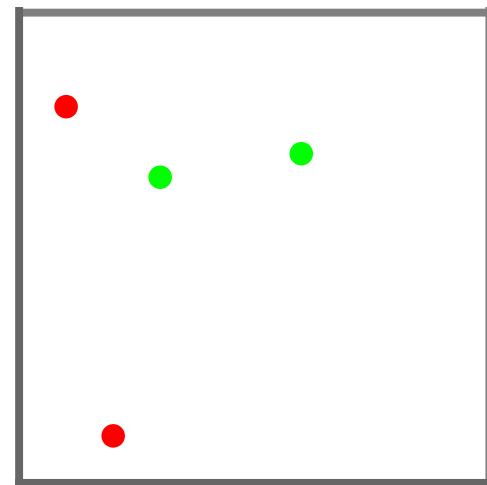
$$V(x, t + T) = V(x, t)$$

$$\frac{L^{d-1}}{e} \int_0^T J_\alpha(t) dt = Q_\alpha \in \mathbb{Z}$$

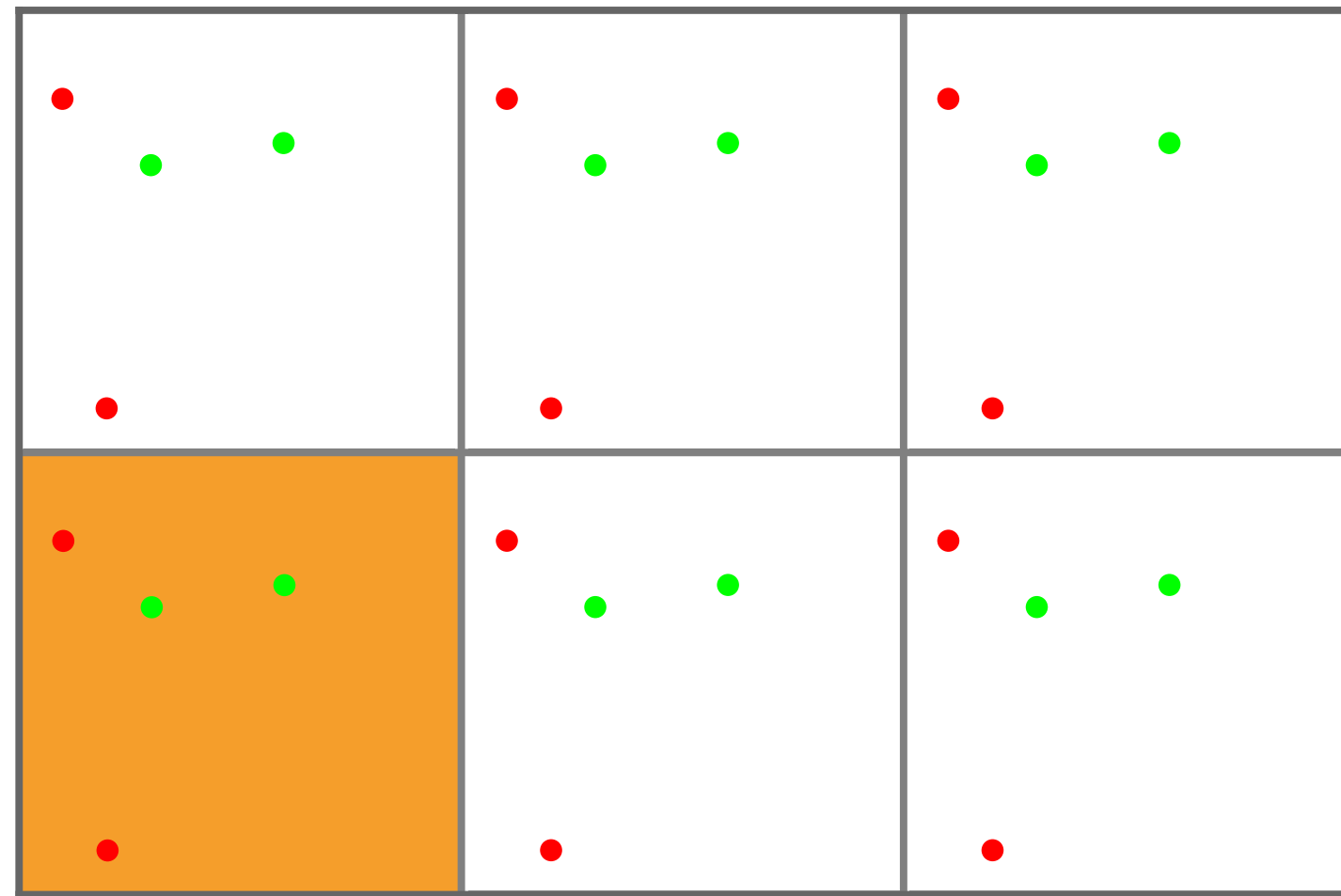
D.J. Thouless, *Quantization of particle transport*, Phys. Rev. B **27**, 2083 (1983)



what are oxydation states, in the first place?

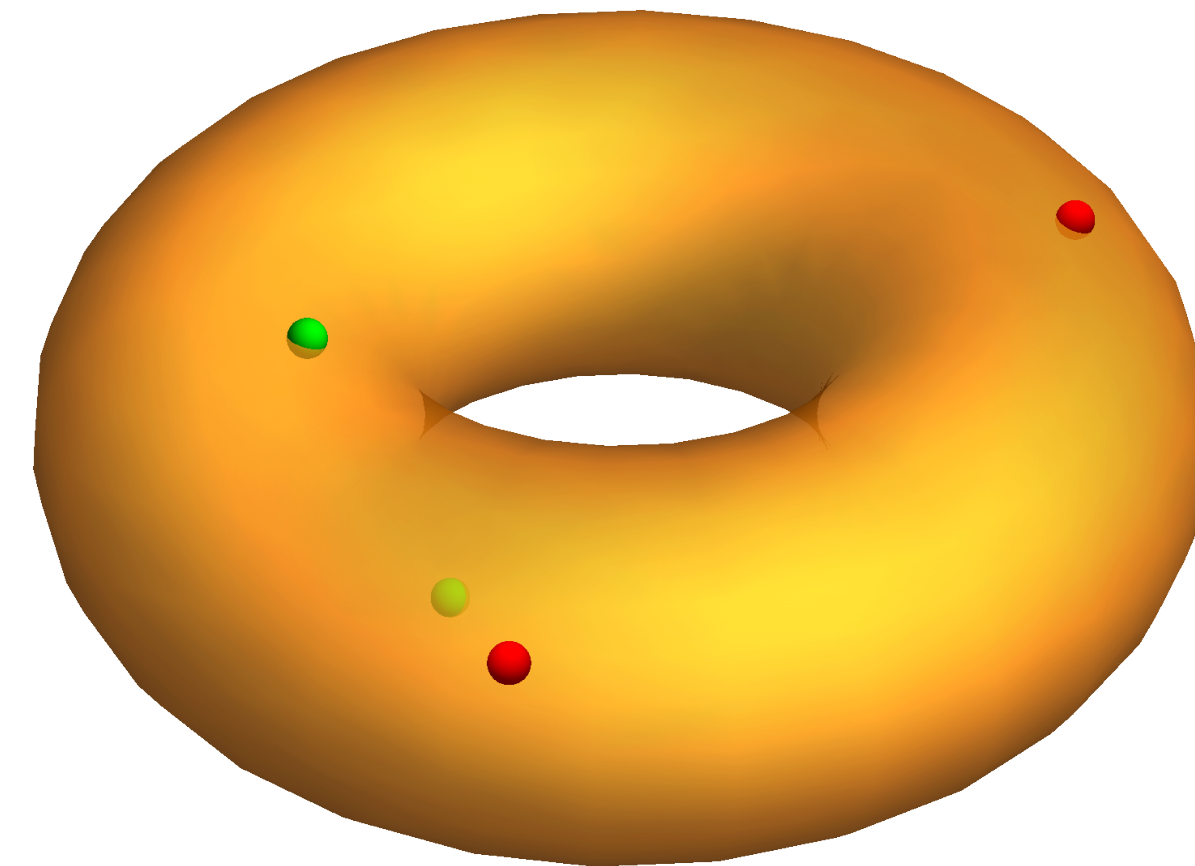
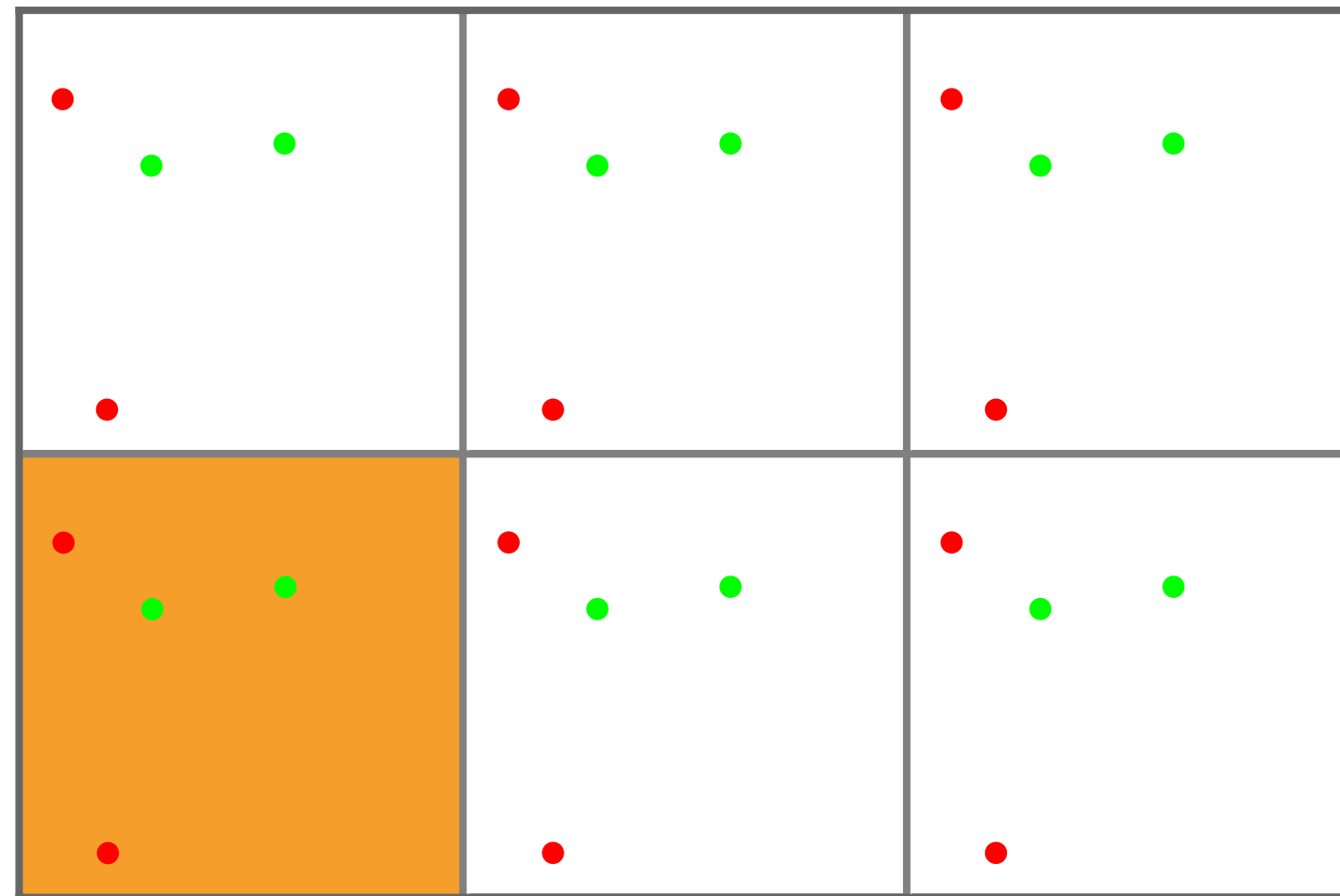


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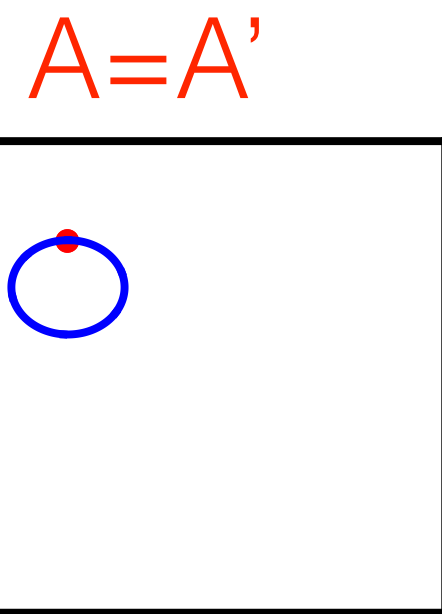


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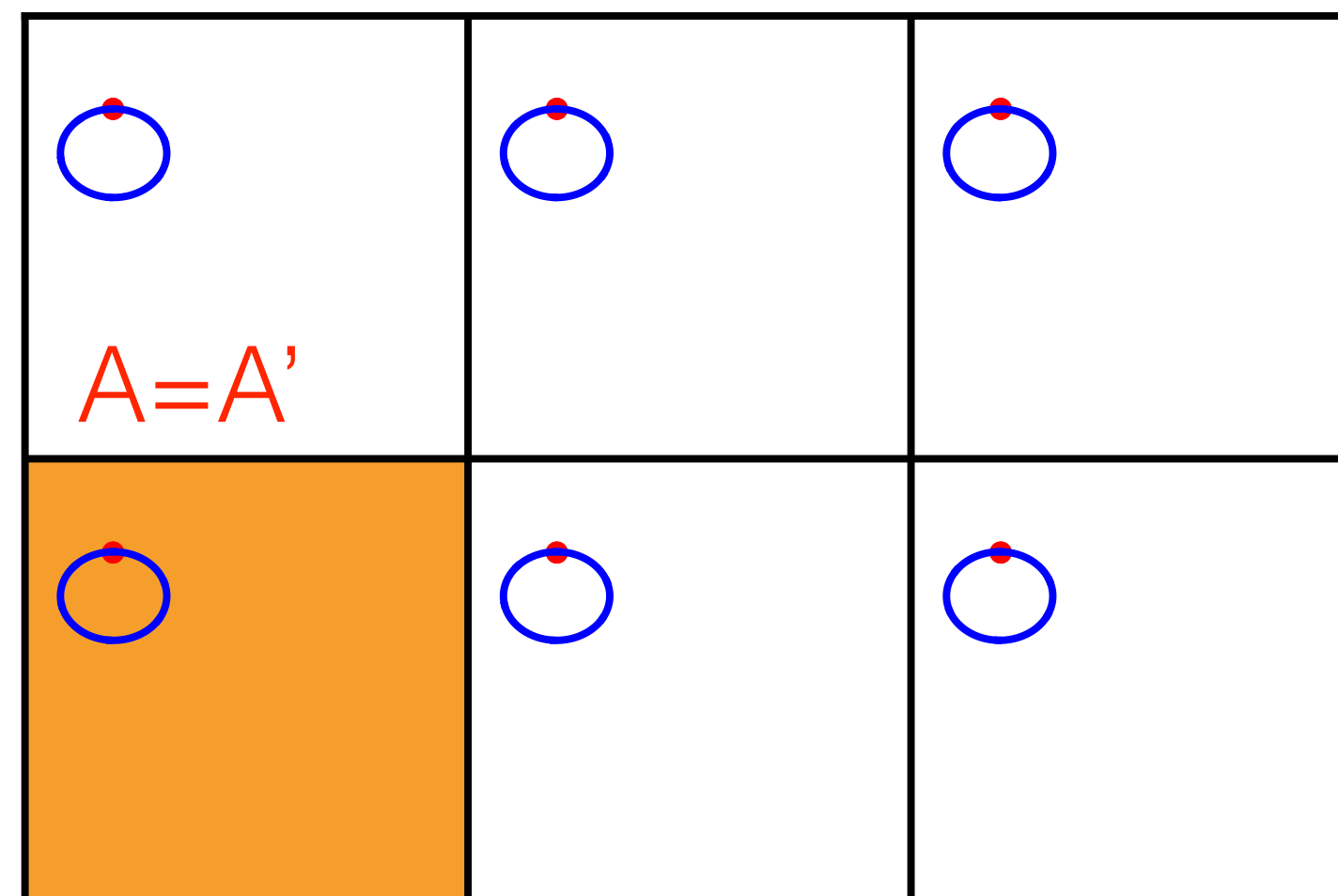
$$[0, L]^{3N} \xrightarrow{\text{PBC}} \mathbb{T}^{3N}$$



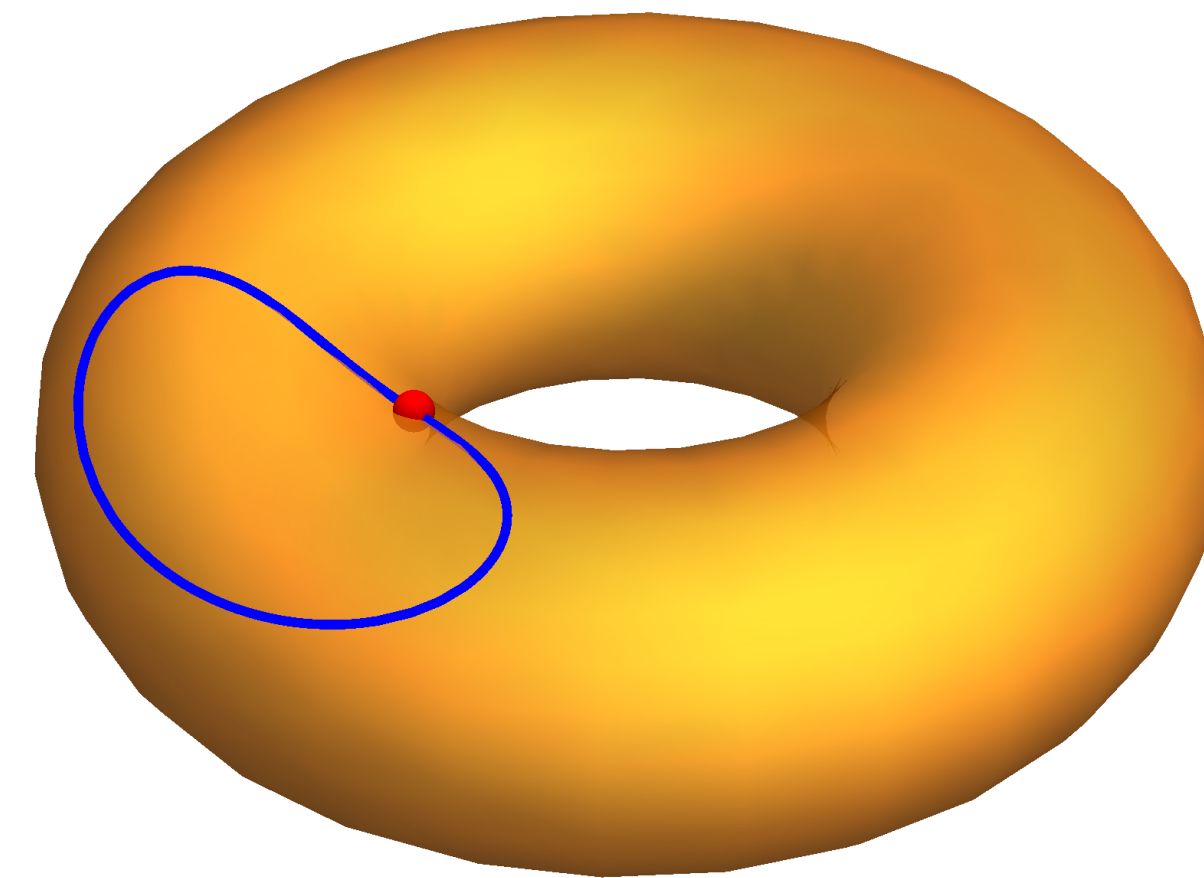
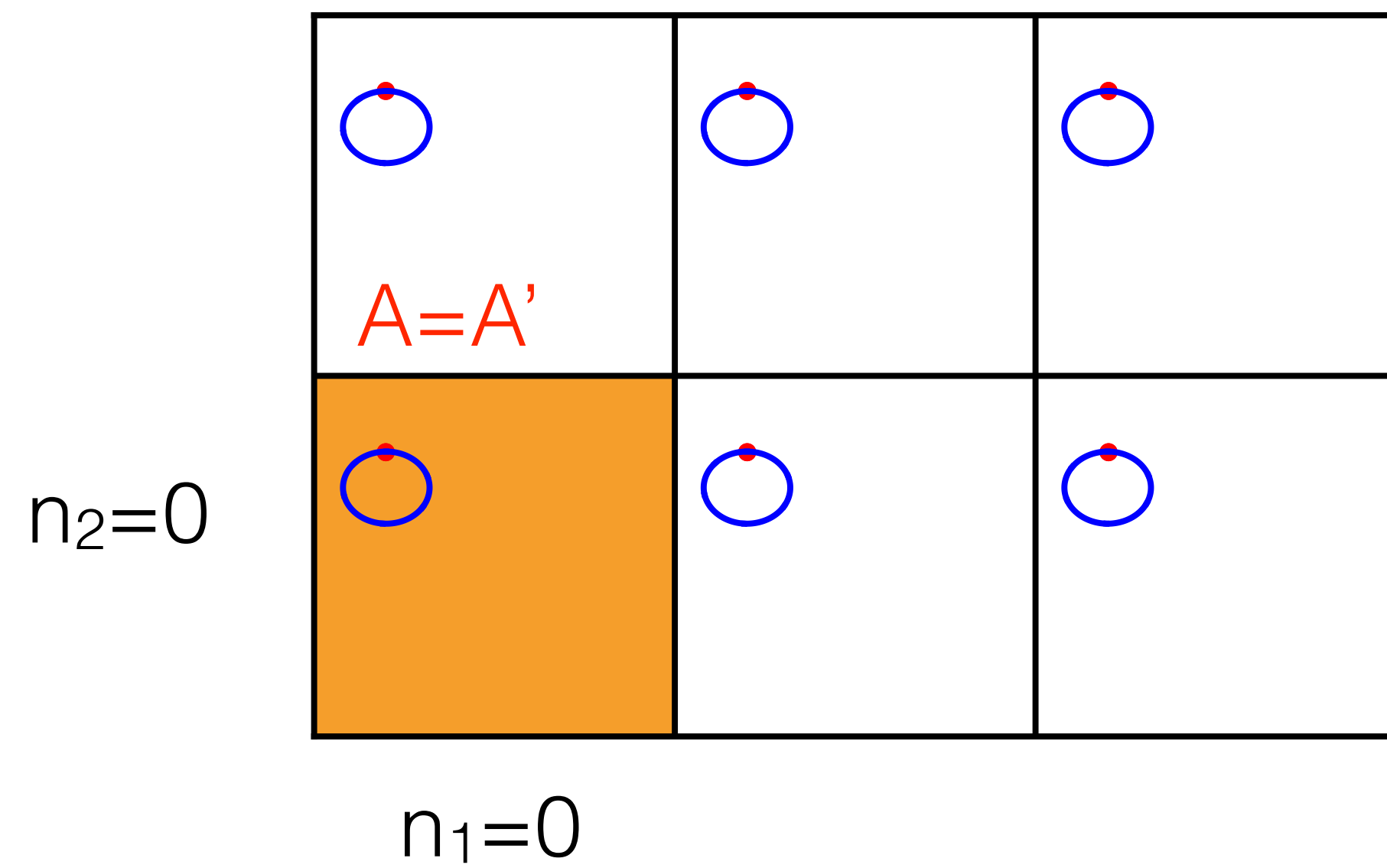
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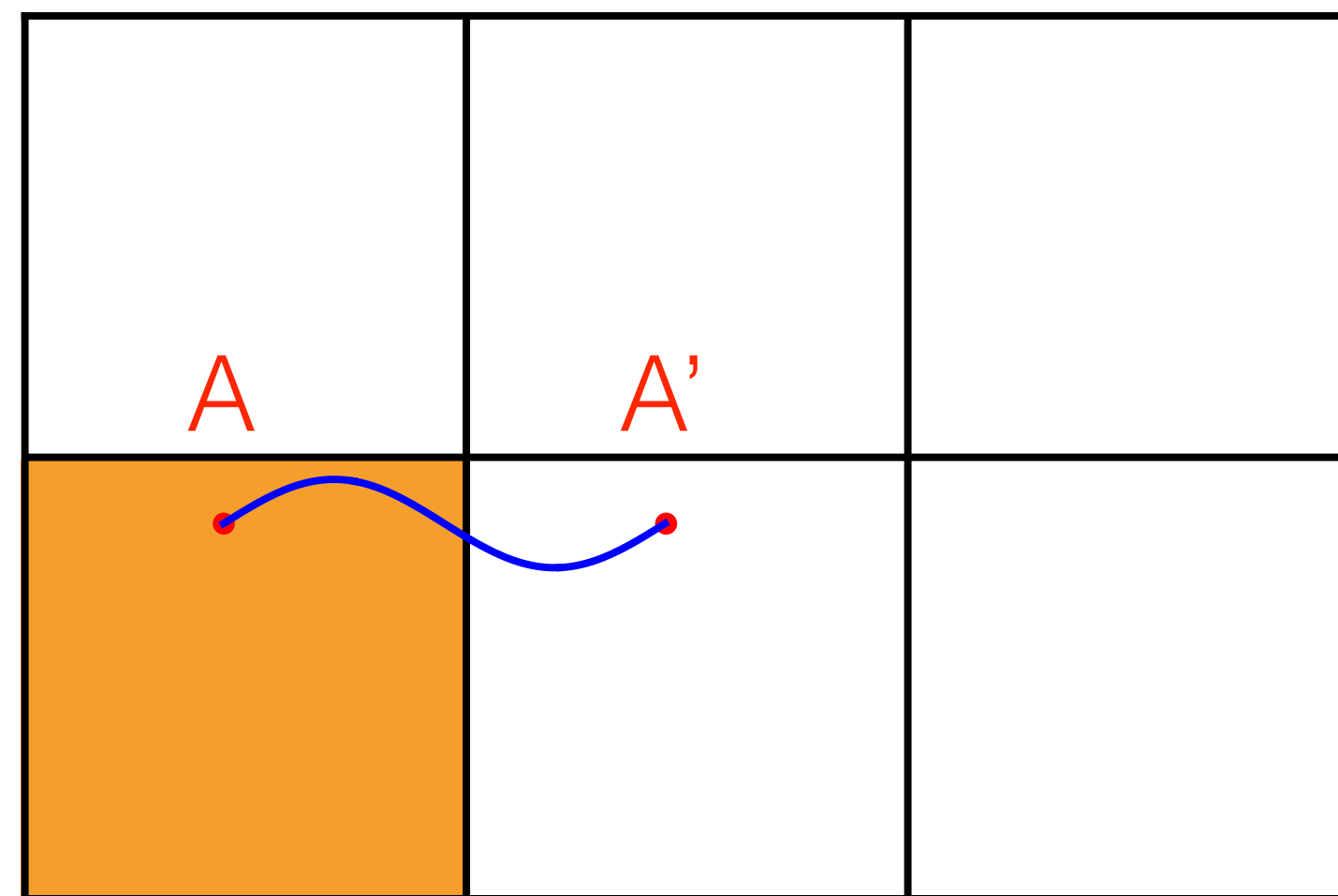
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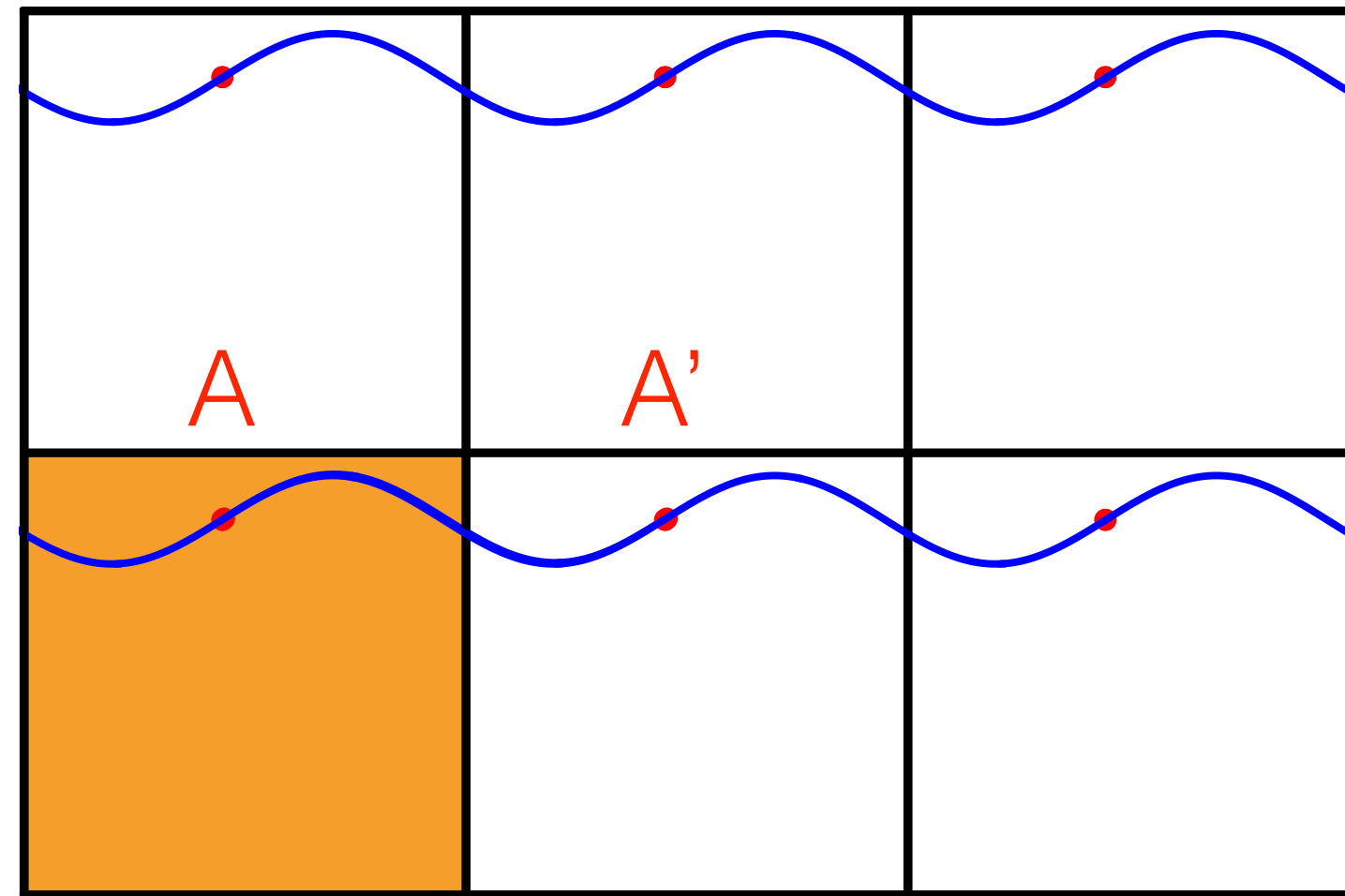
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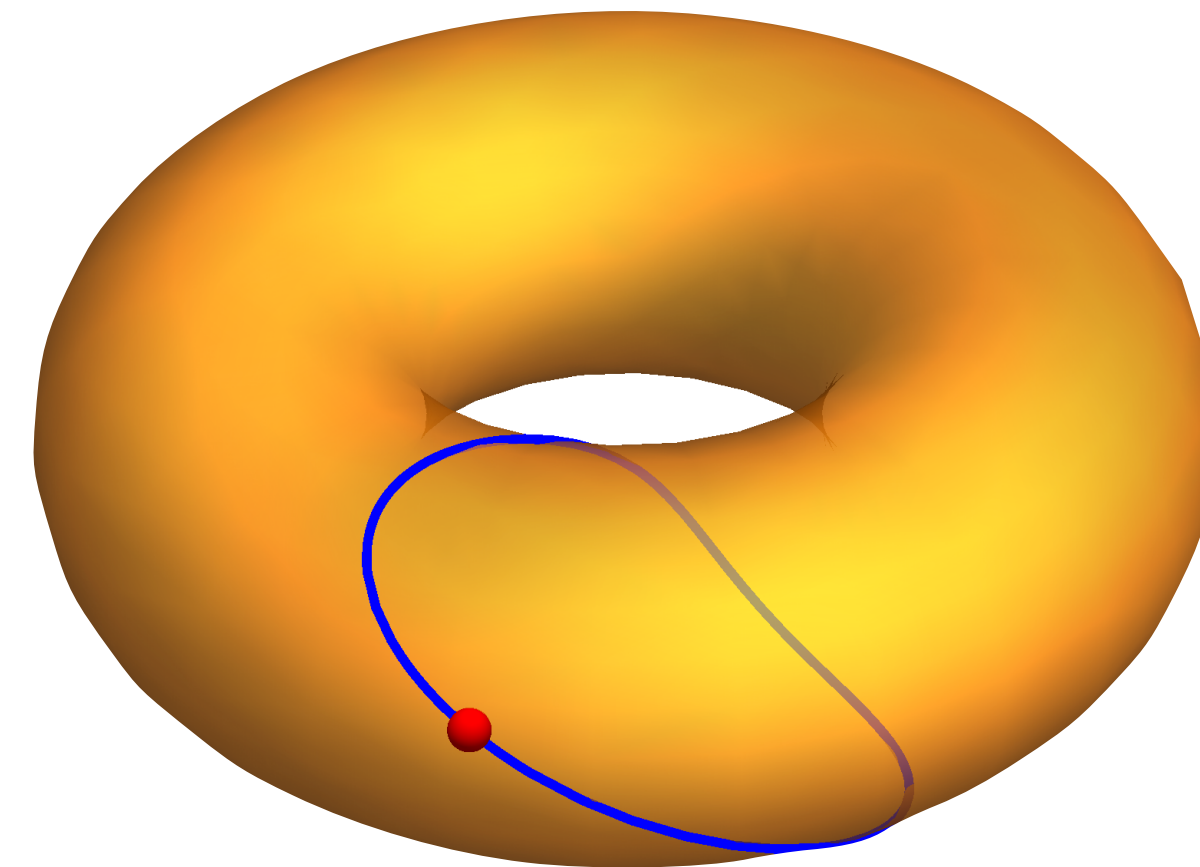
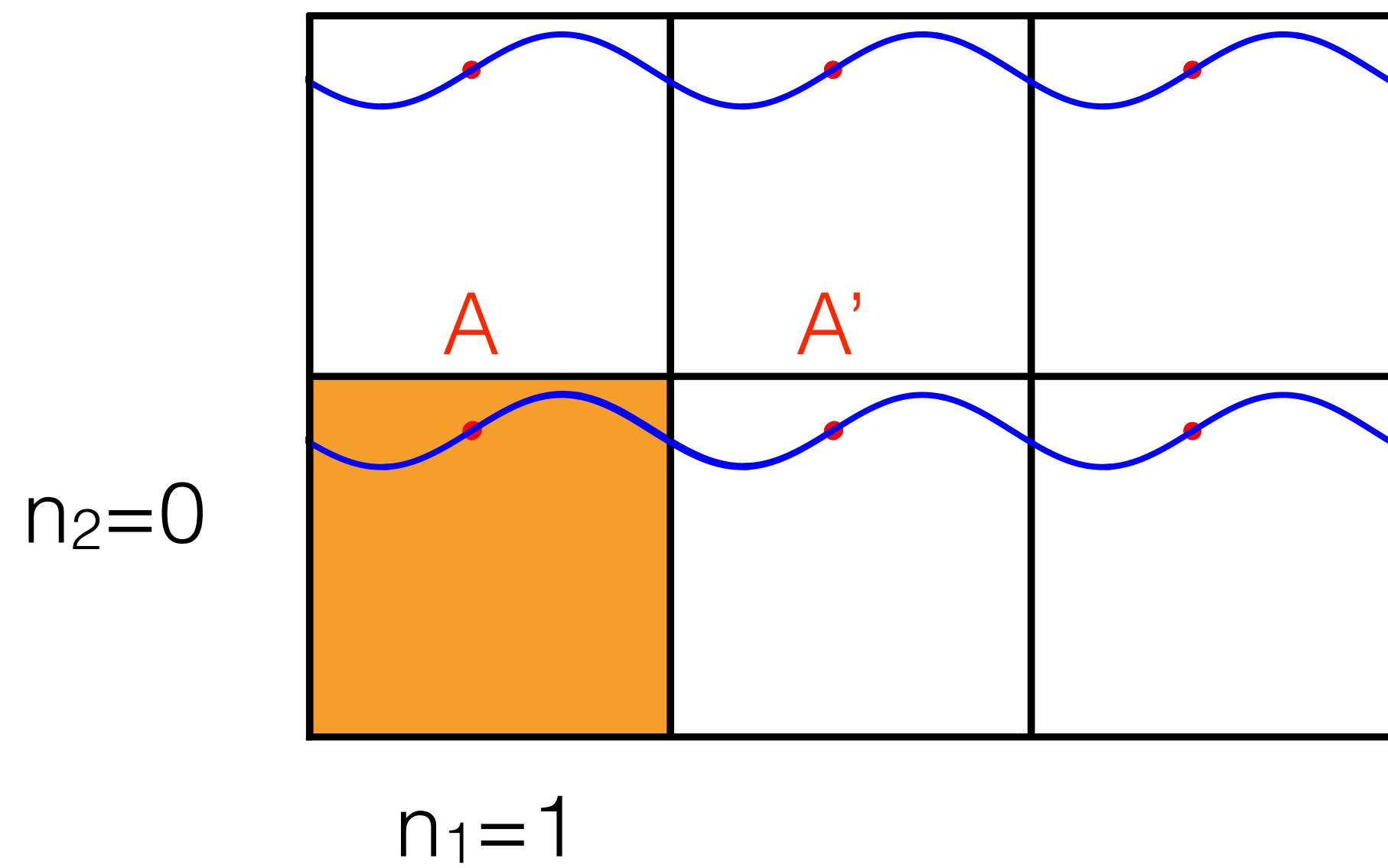
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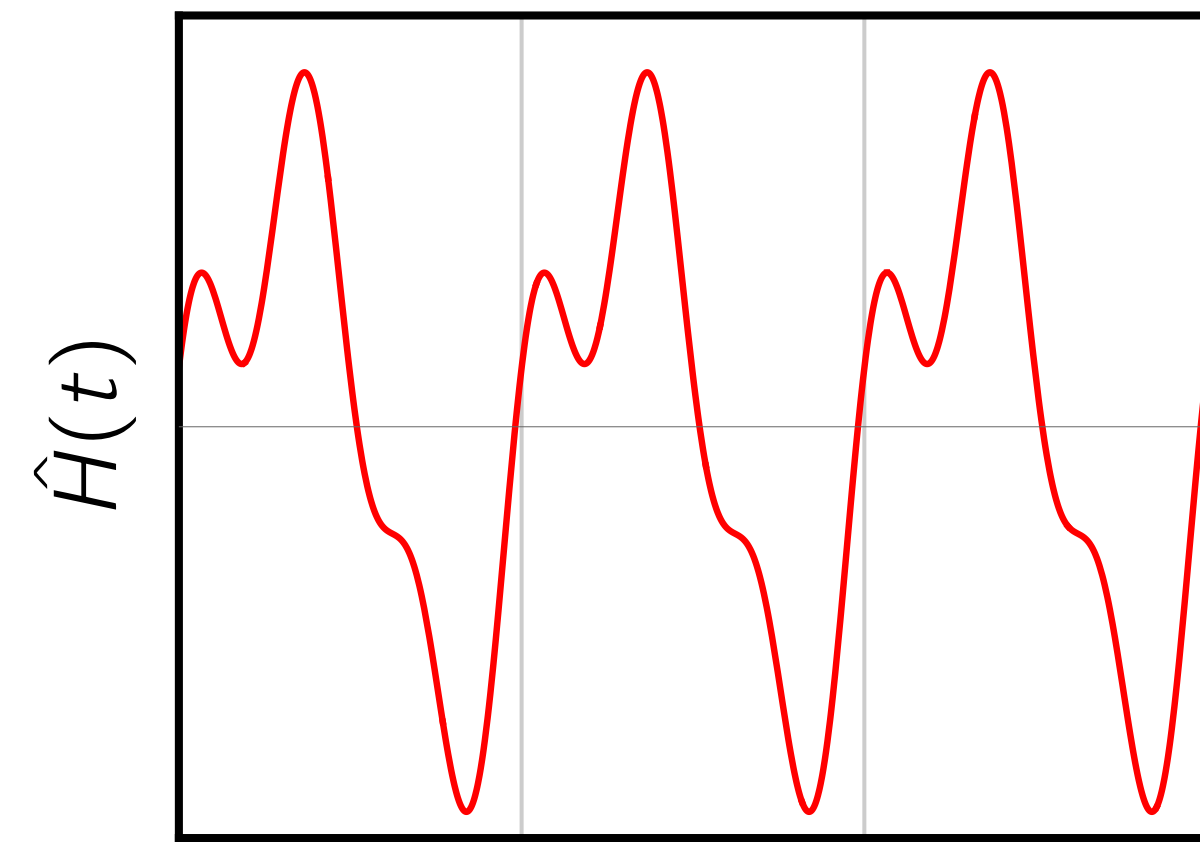
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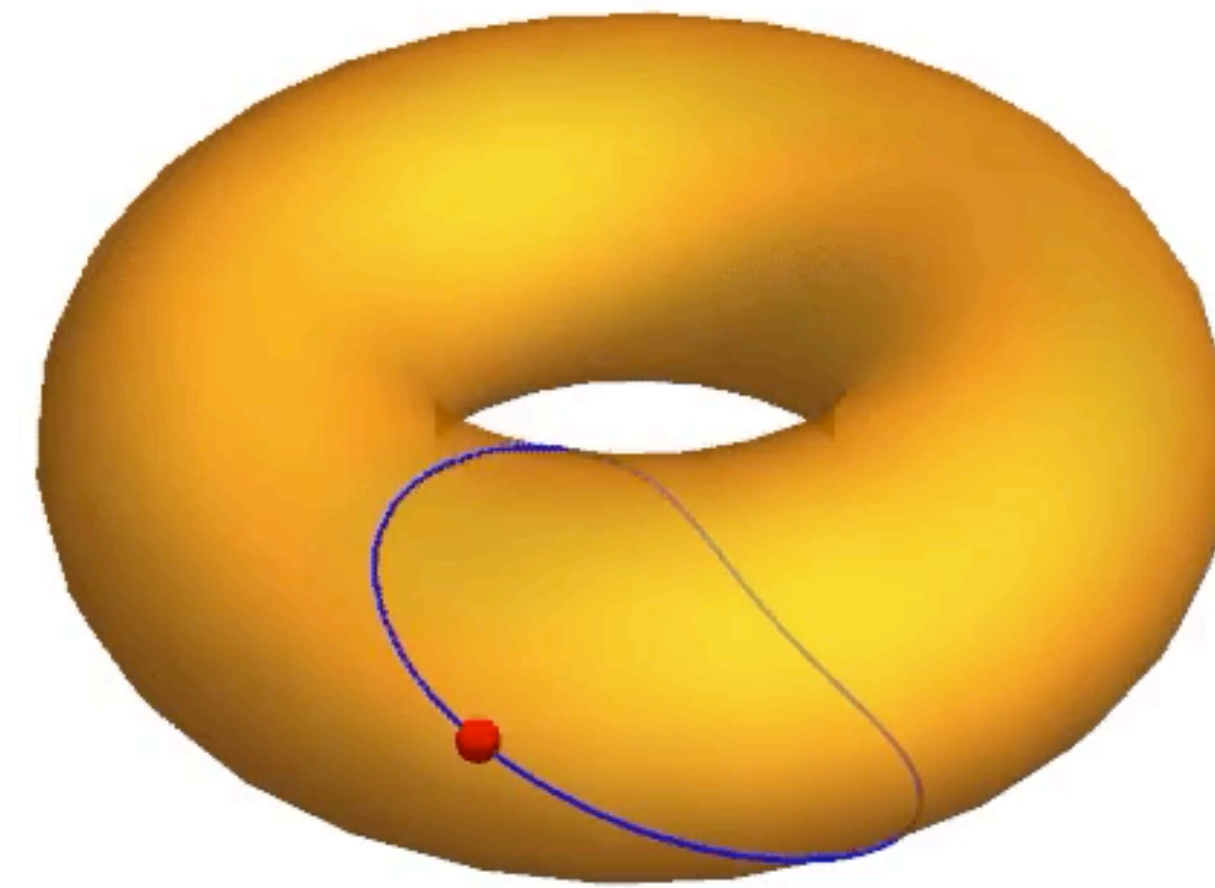
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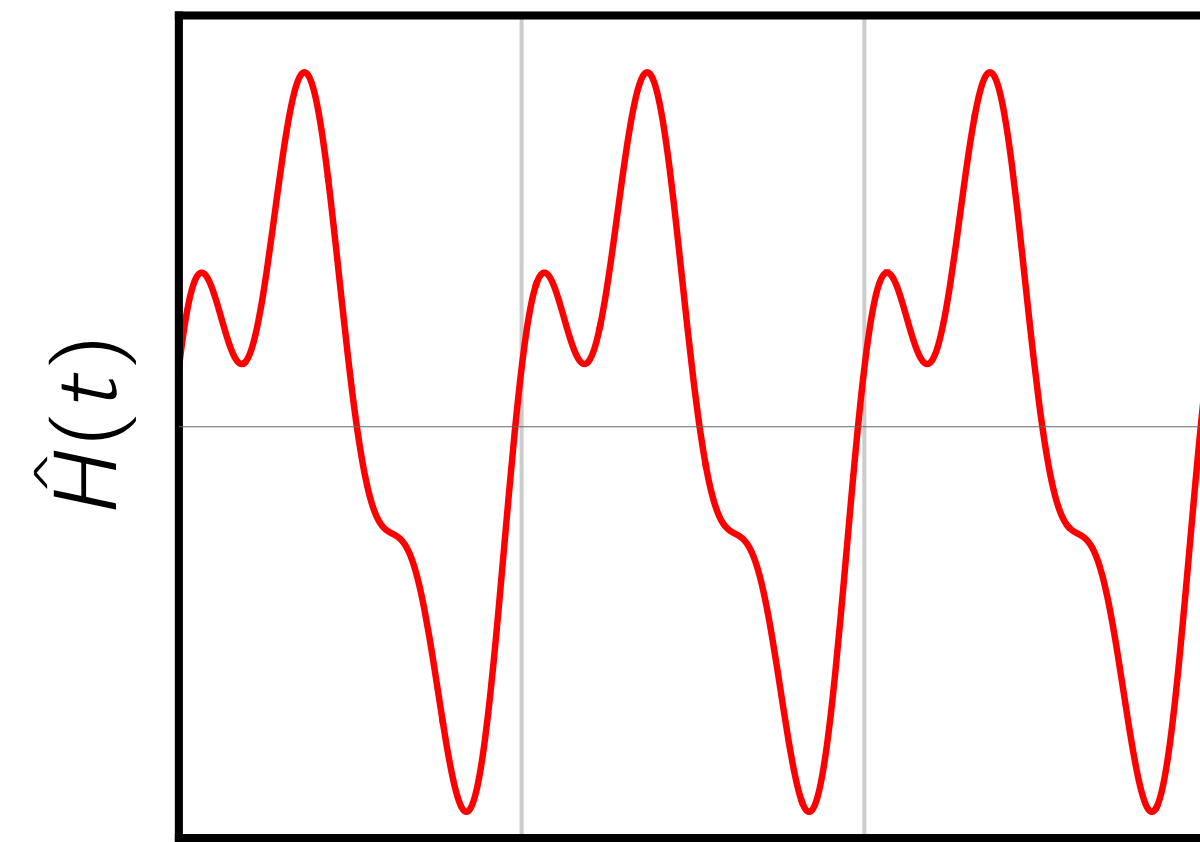
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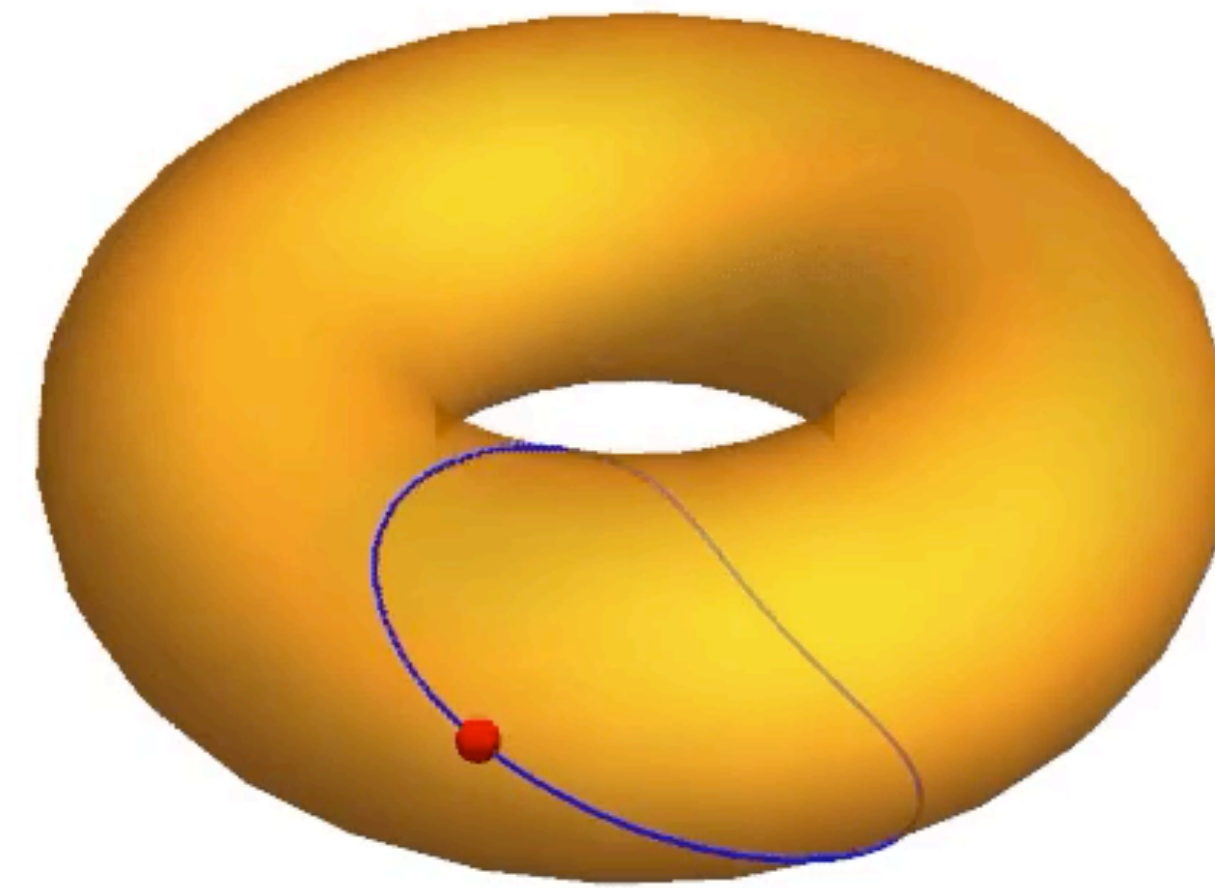
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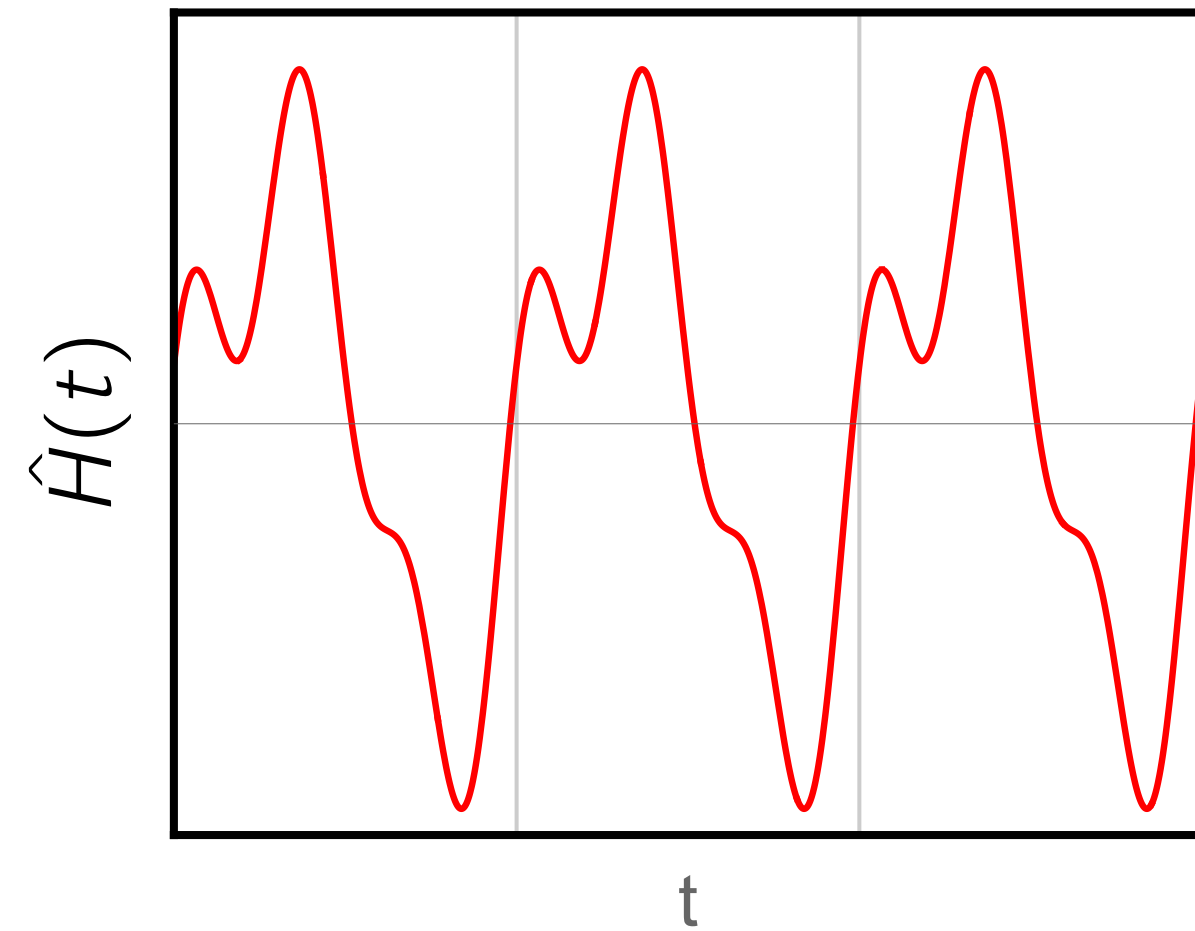
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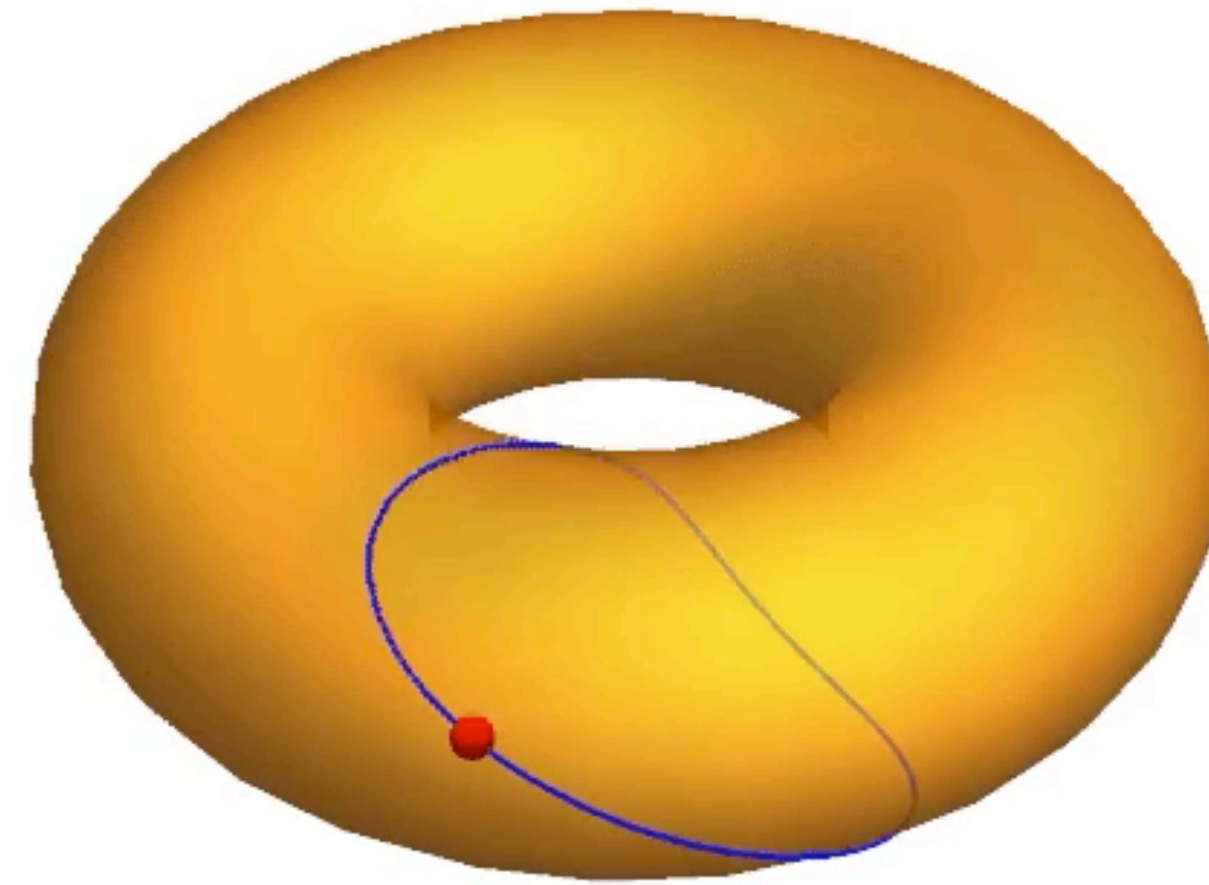
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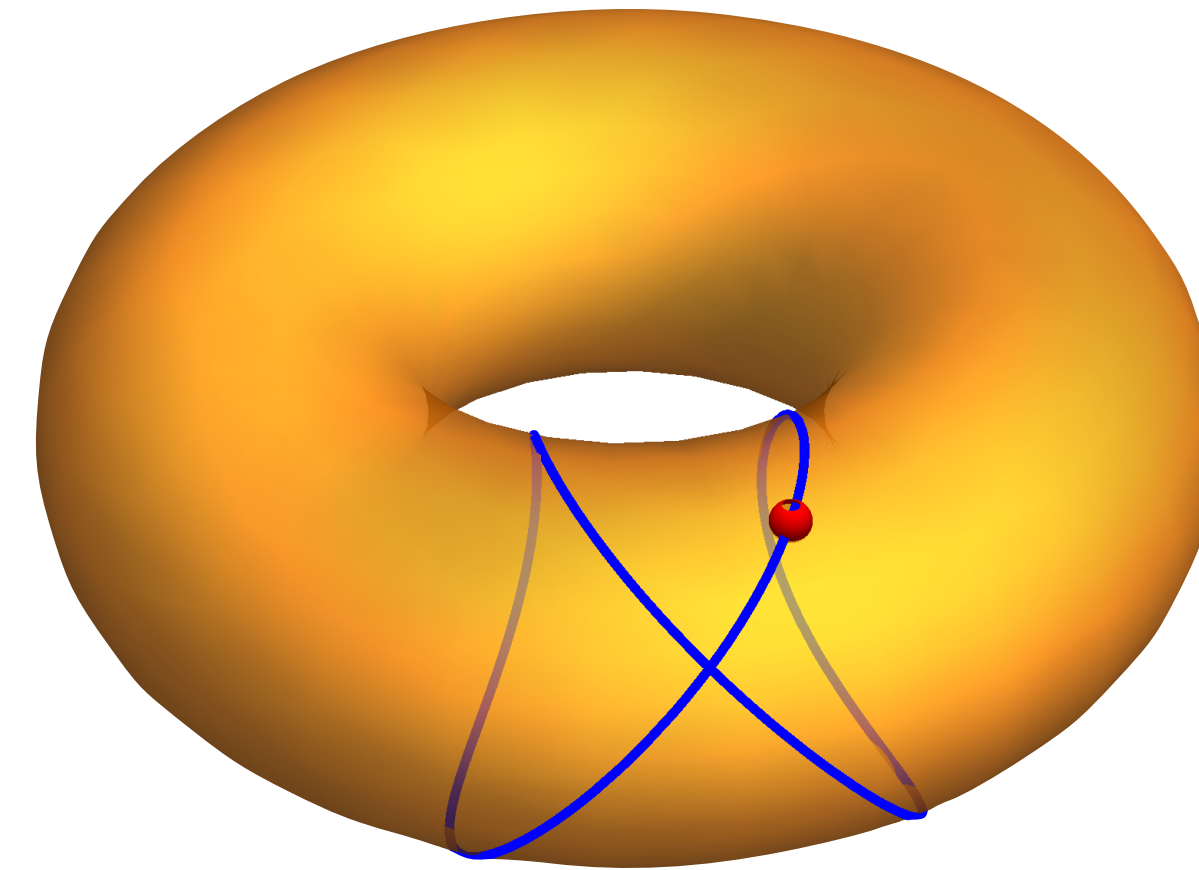
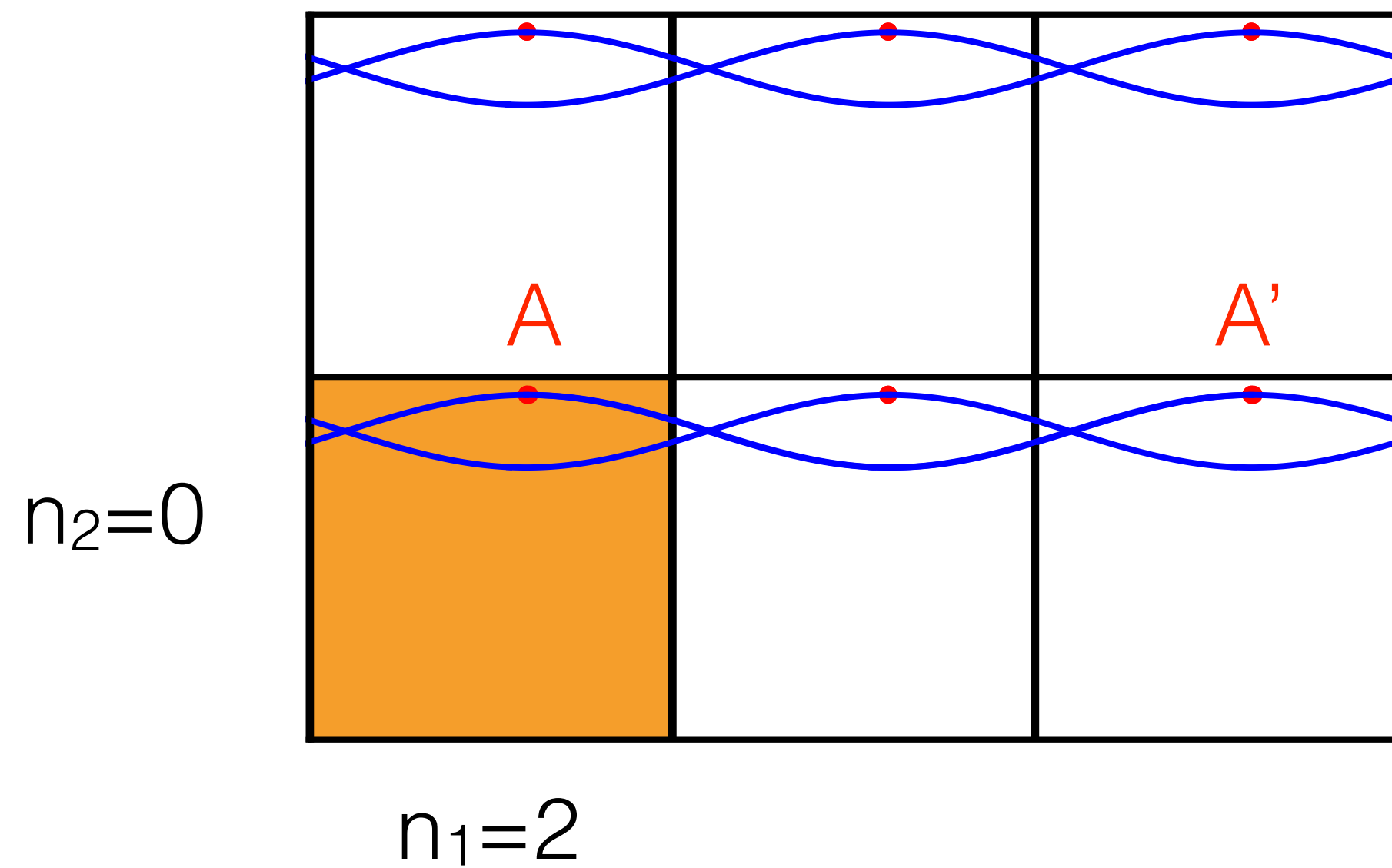
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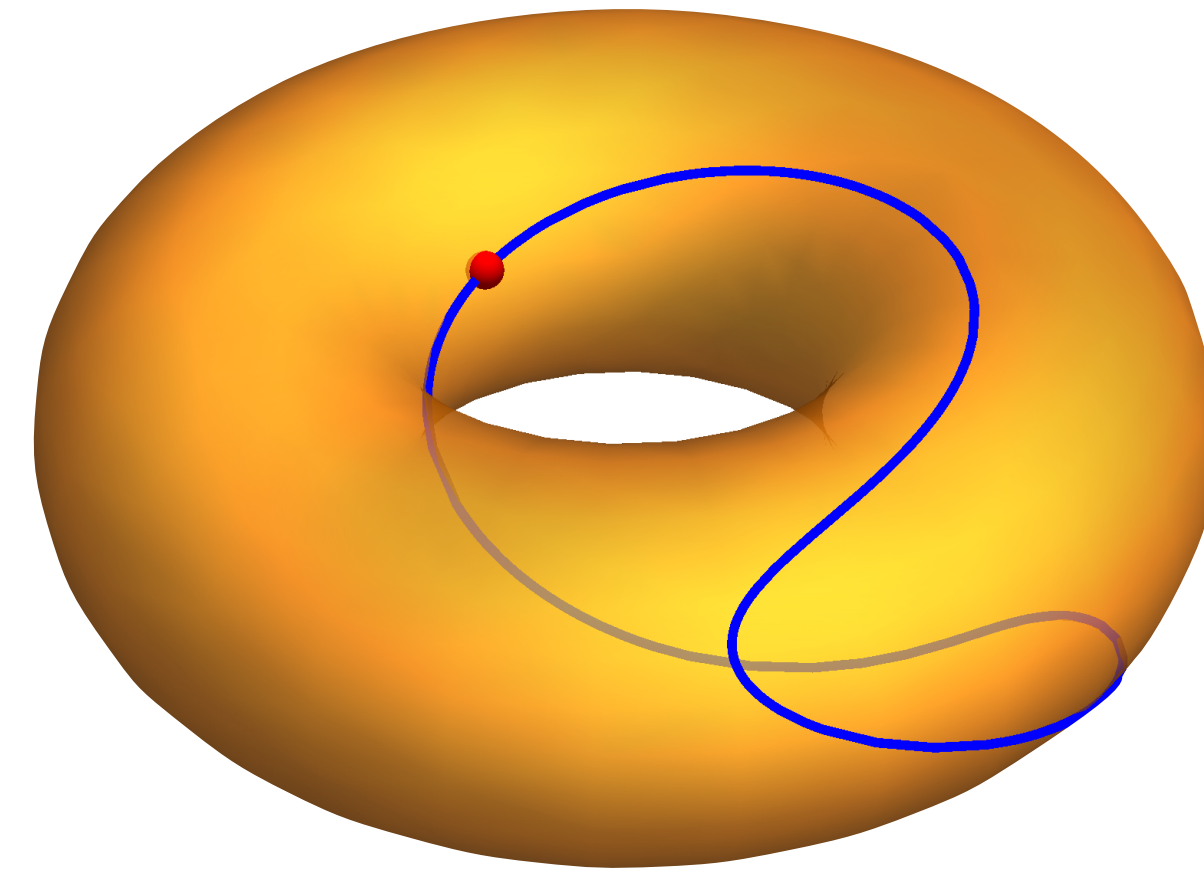
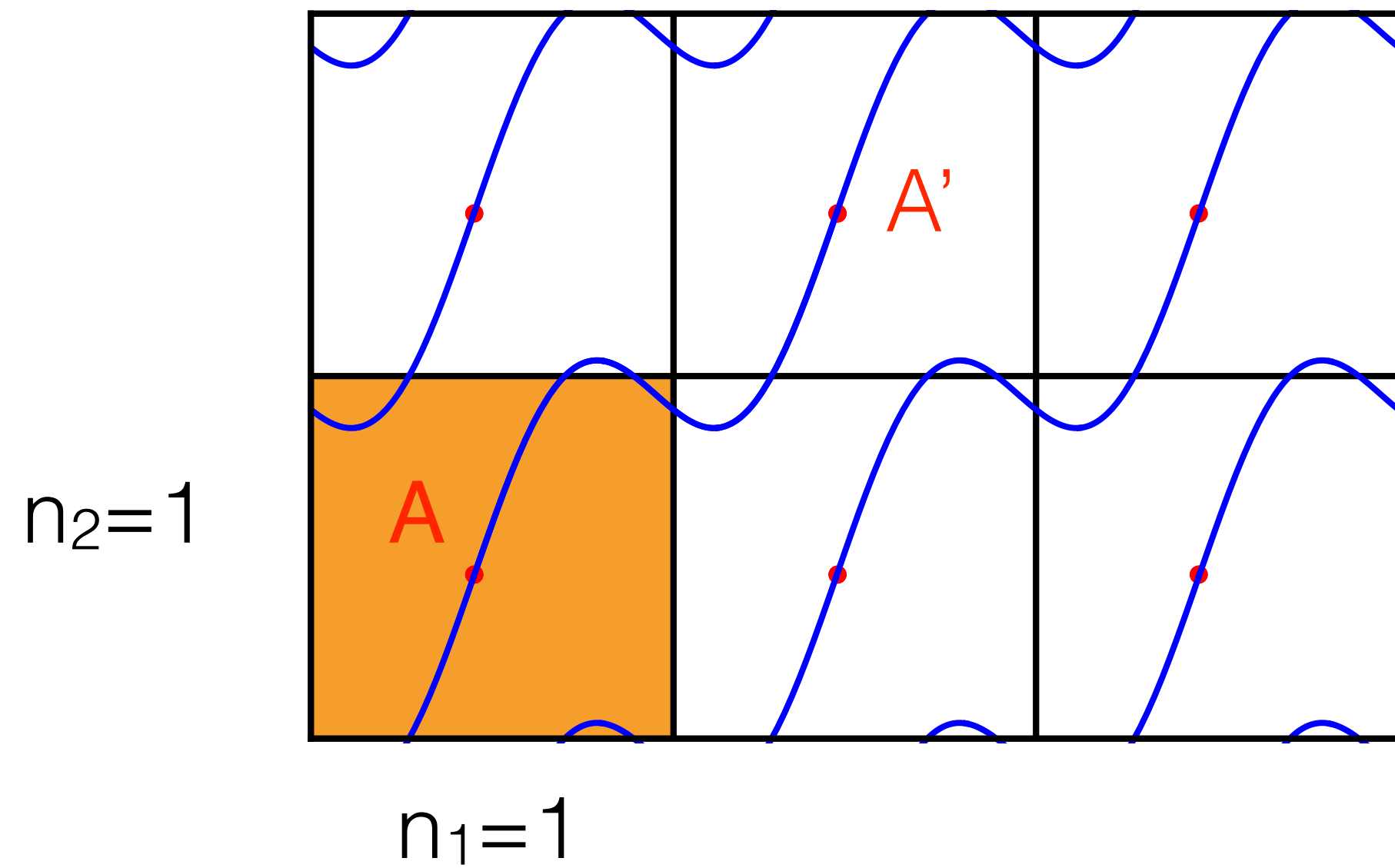


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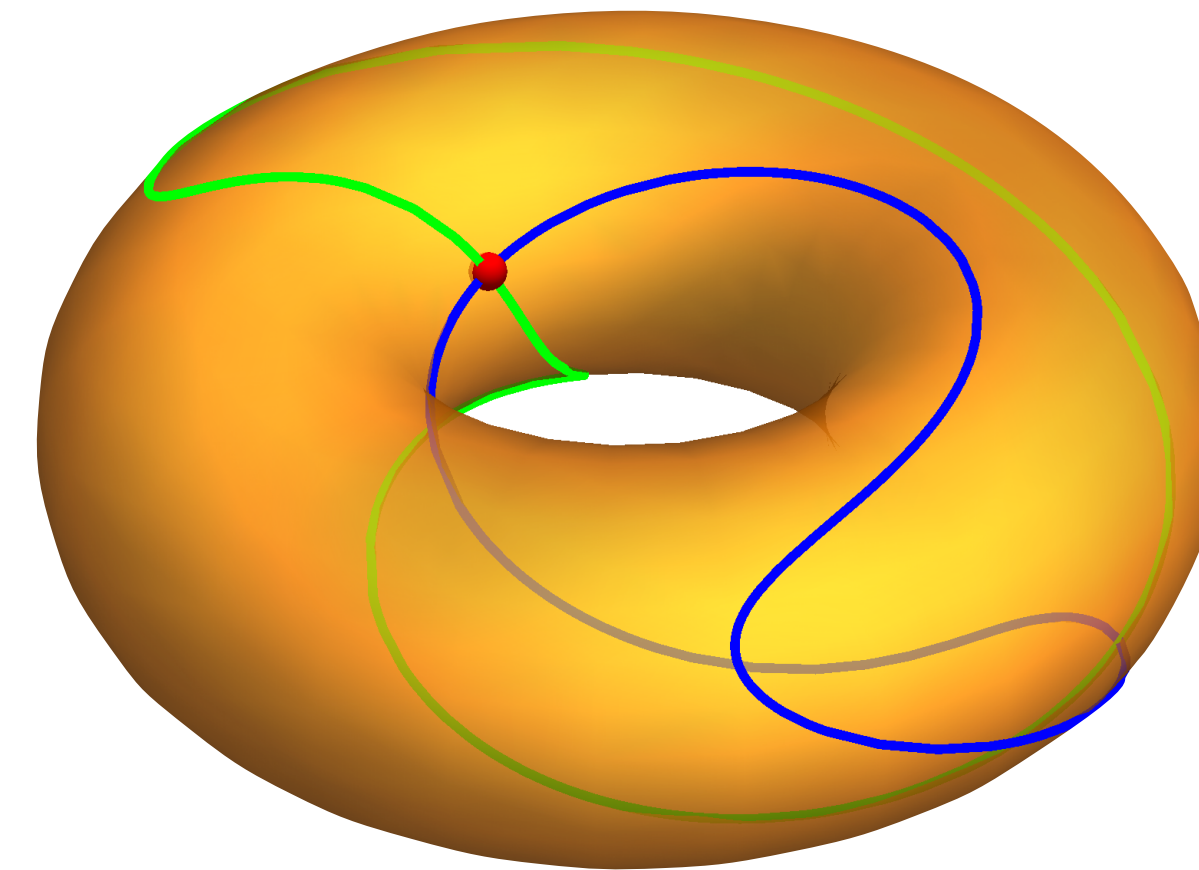
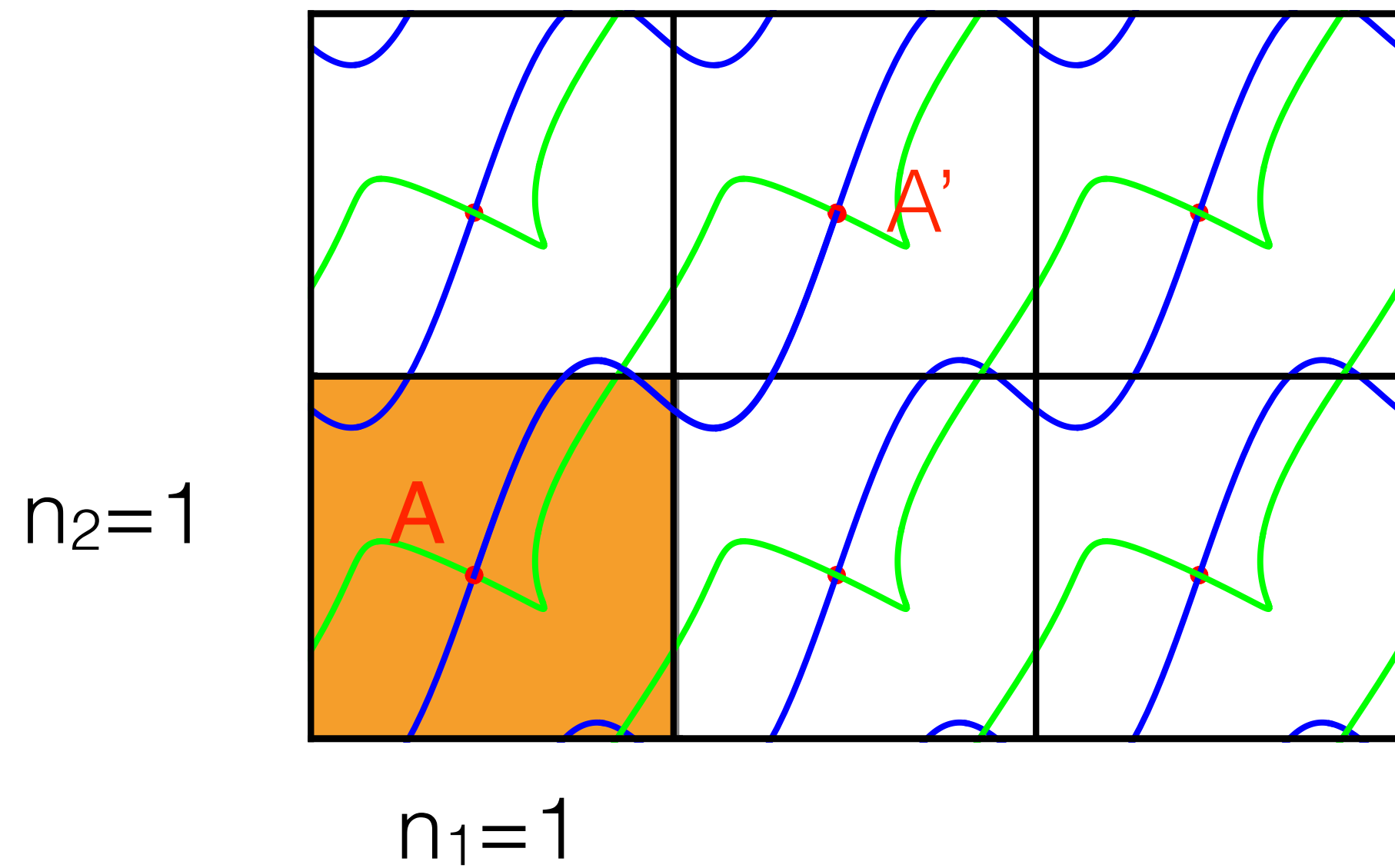
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what are oxydation states, in the first place?



$$Q_\alpha(AA') = Q_\alpha(AA') = Q_\alpha[n_1 = 1, n_2 = 1]$$

what are oxydation states, in the first place?

$$Q_{\alpha}[C] = \frac{1}{\ell} \mu_{\alpha}[C]$$



what are oxydation states, in the first place?

$$\begin{aligned} Q_{\alpha}[\mathcal{C}] &= \frac{1}{\ell} \mu_{\alpha}[\mathcal{C}] \\ &= Q_{\alpha}(n_{1x}, n_{1y}, n_{1z}, \dots, n_{Nz}) \end{aligned}$$



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- All loops can be shrunk to a point without closing the gap (*strong adiabaticity*);
- Any two like atoms can be swapped without closing the gap



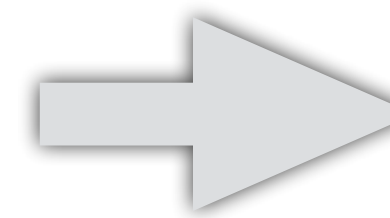
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$$q_{i\alpha\beta} = q_{S(i)} \delta_{\alpha\beta}$$

atomic oxidation state

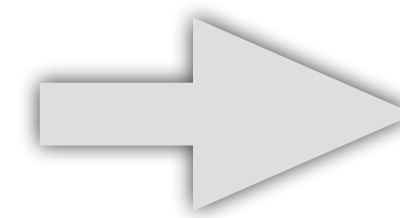
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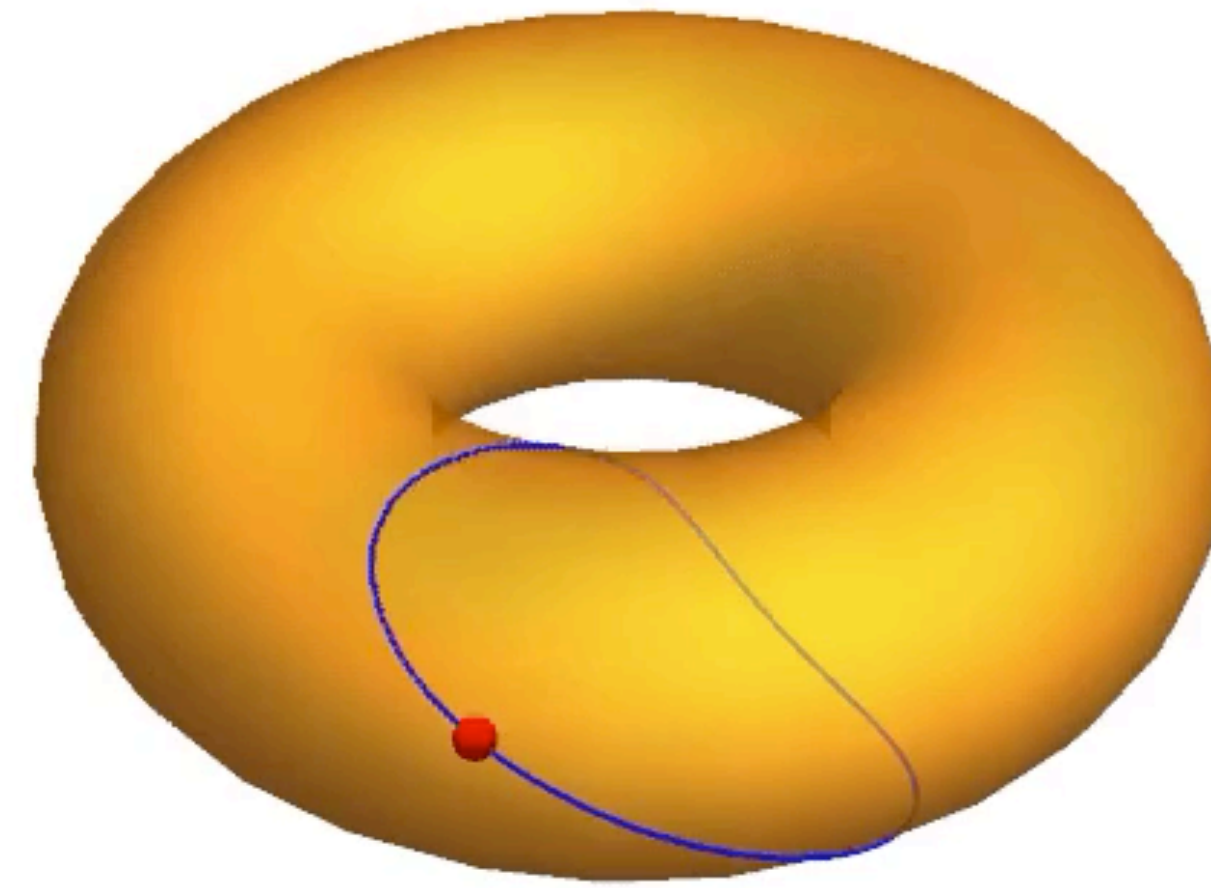
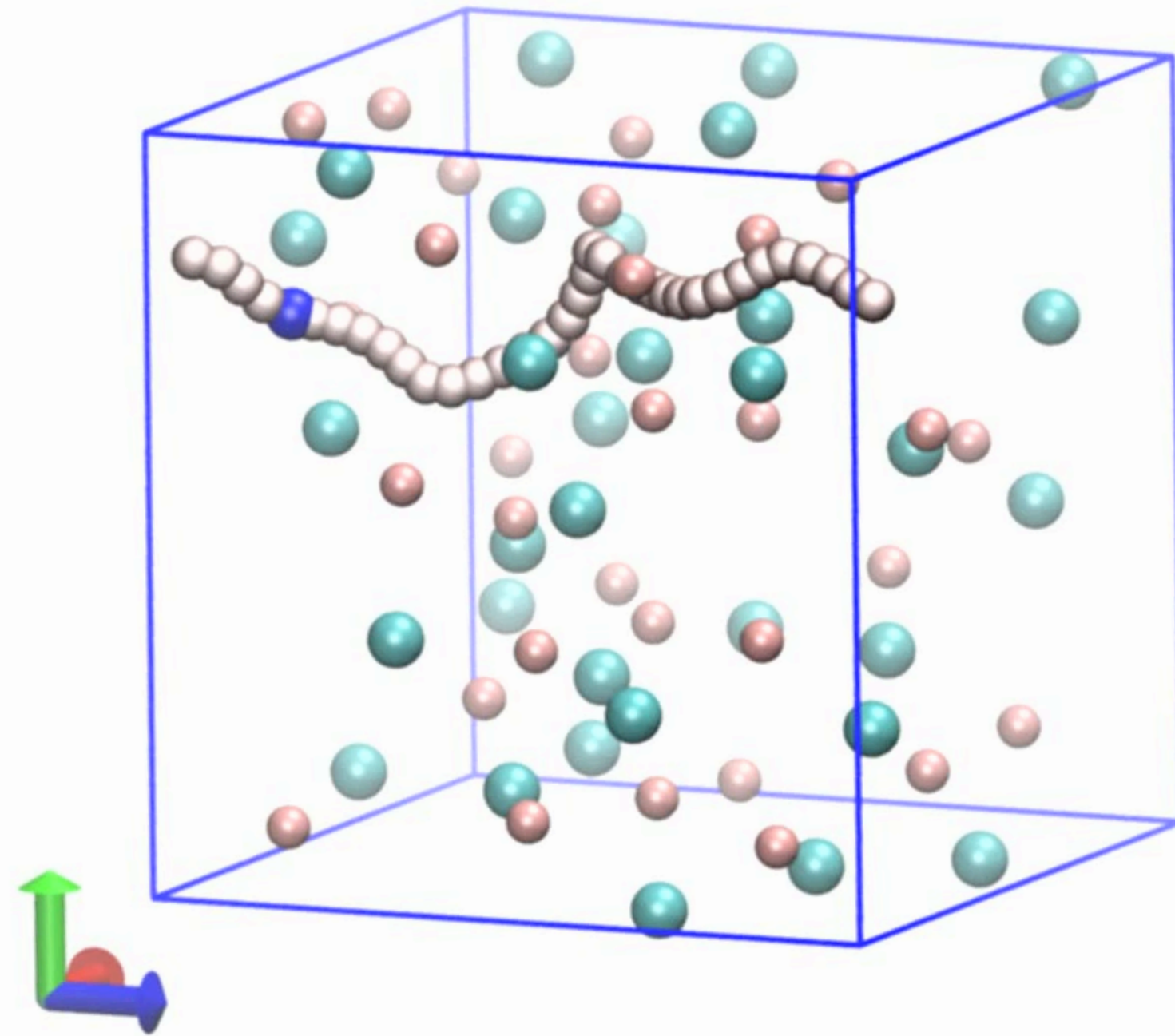
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atomic oxidation state

... they are topological invariants!

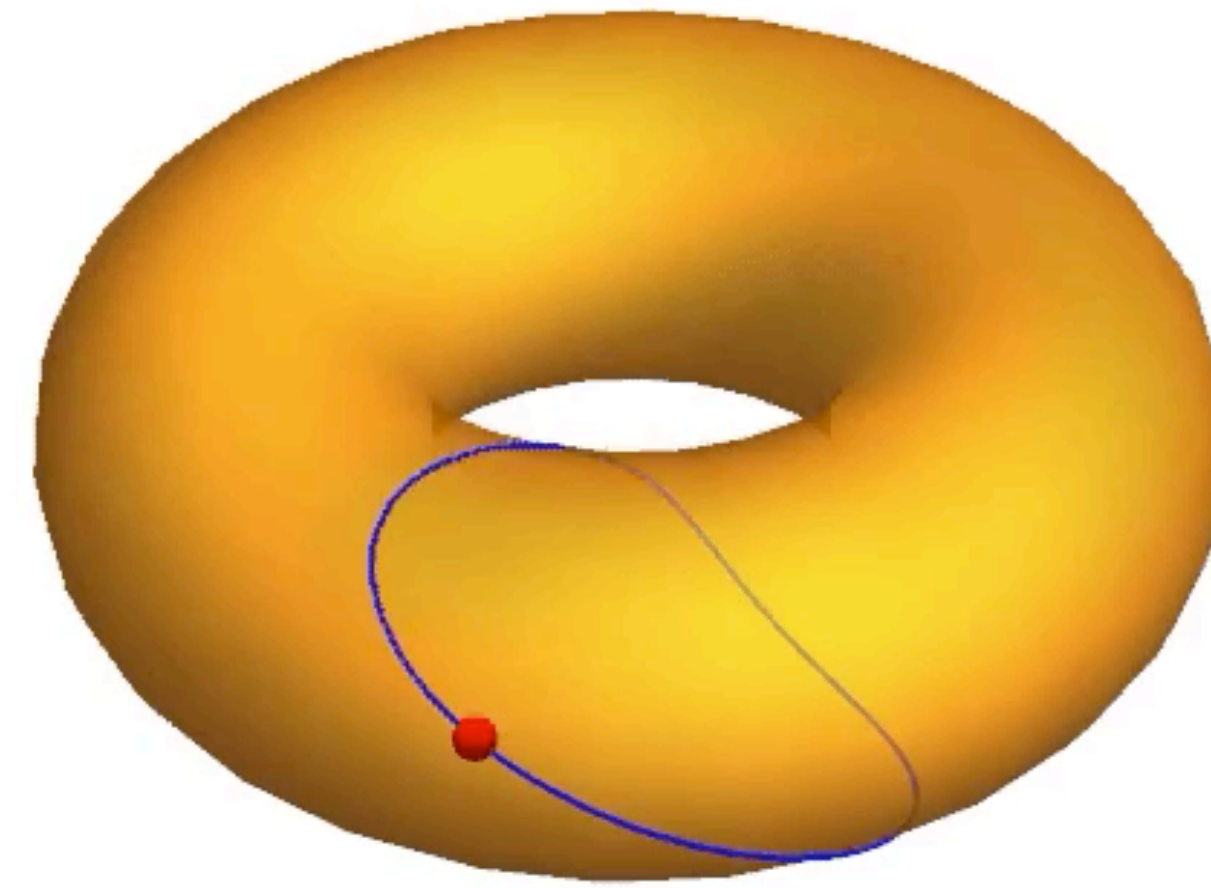
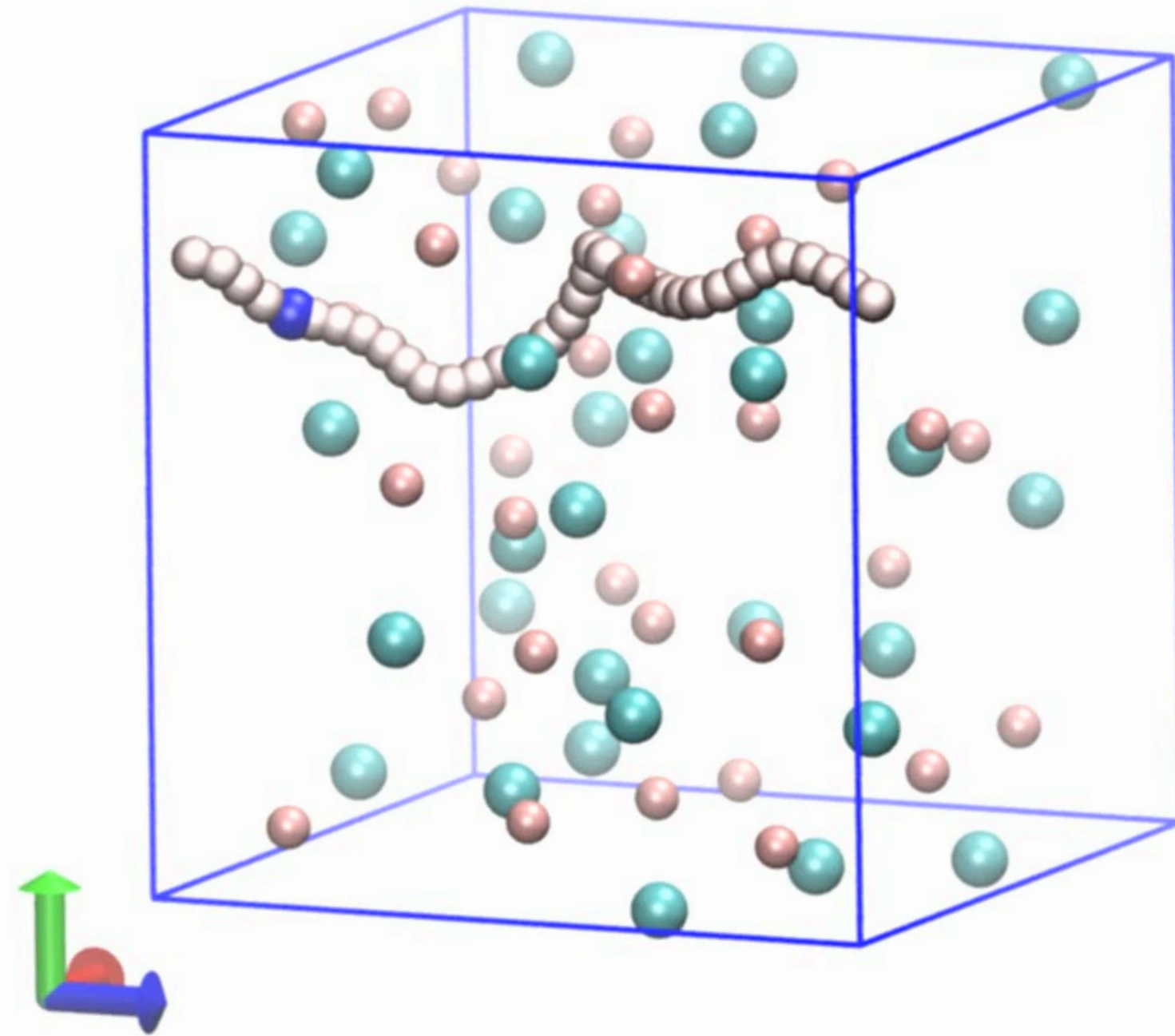


a numerical experiment on molten KCl



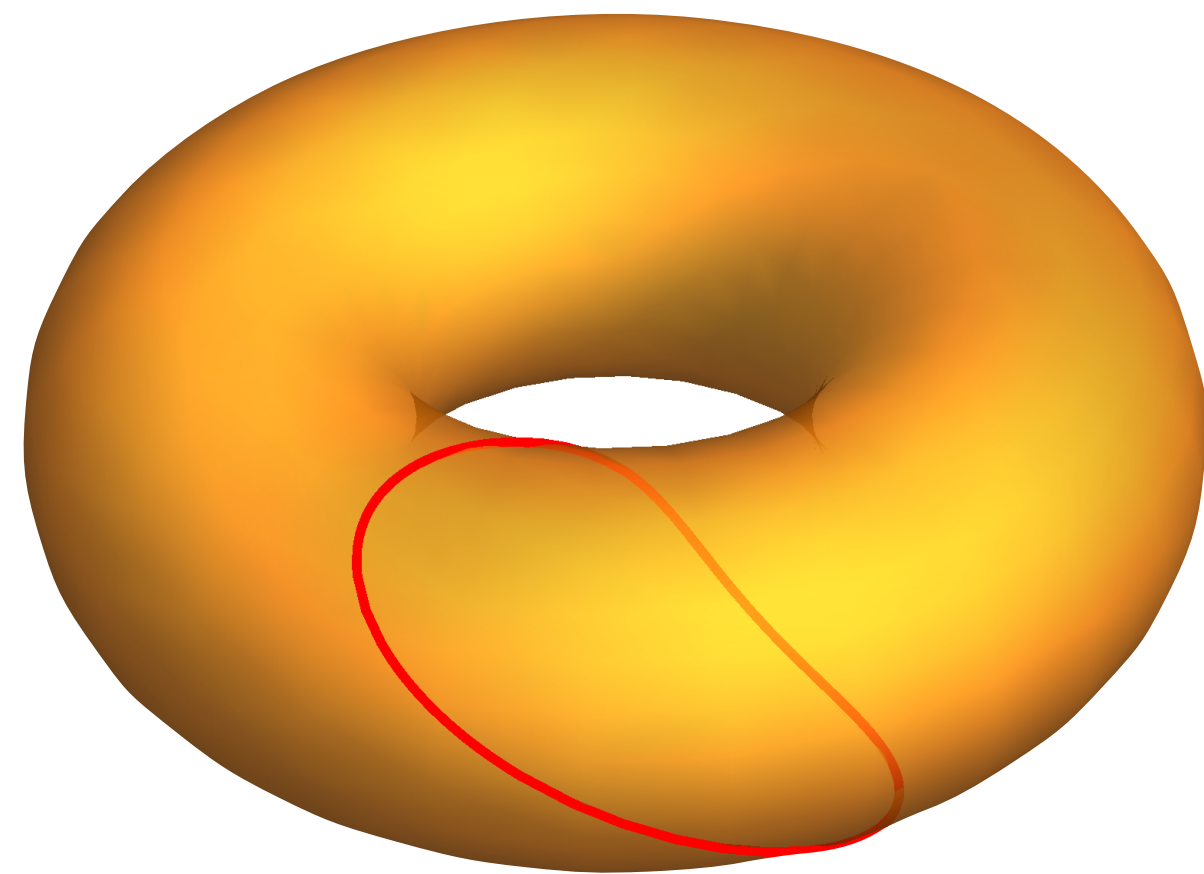
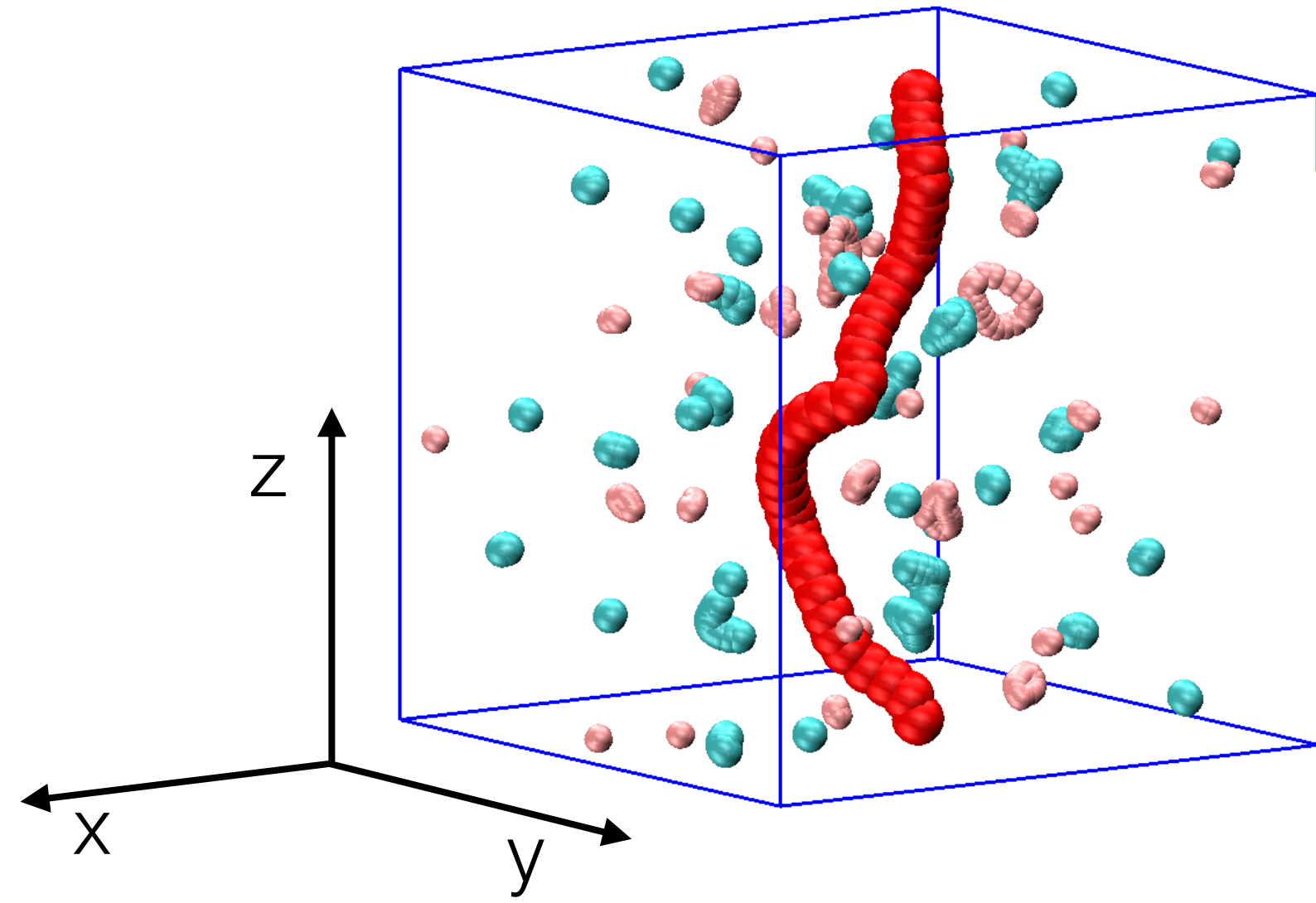
a topologically non-trivial minimum-energy path
connecting two identical configurations of a ionic melt

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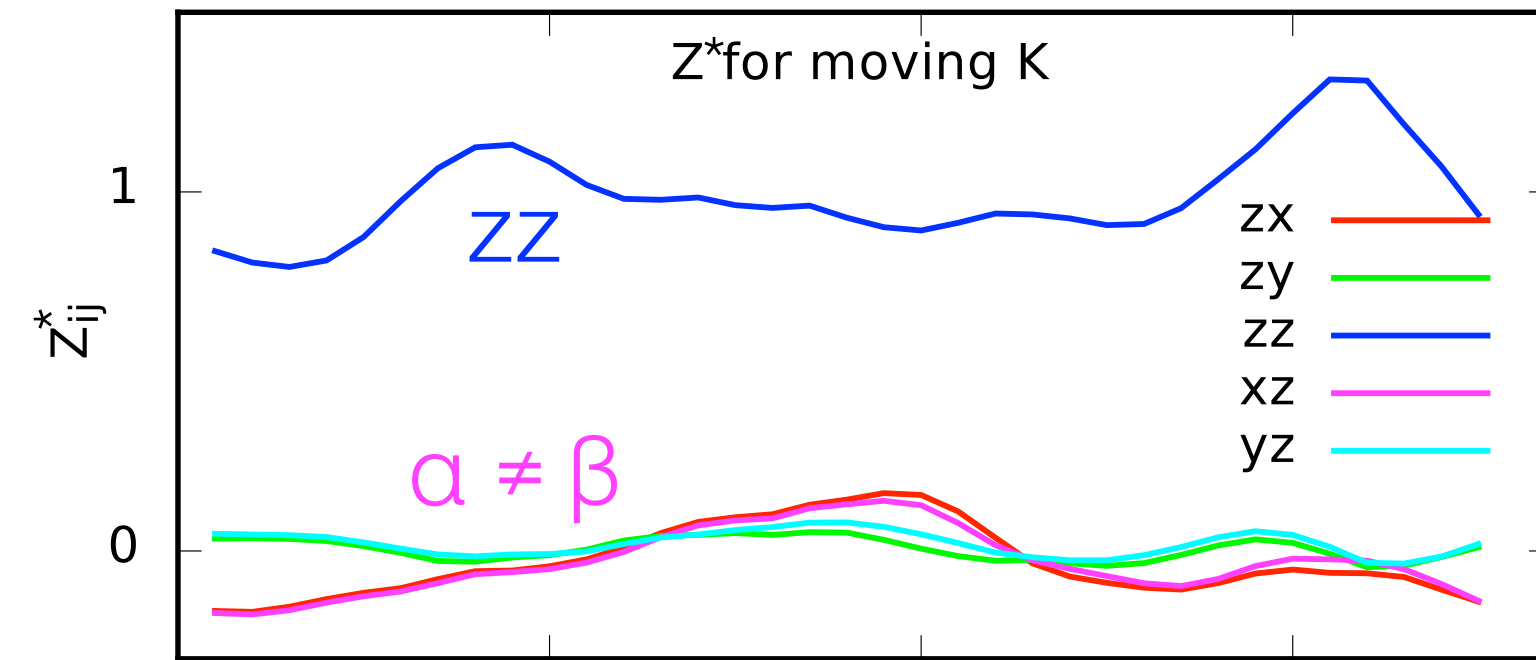
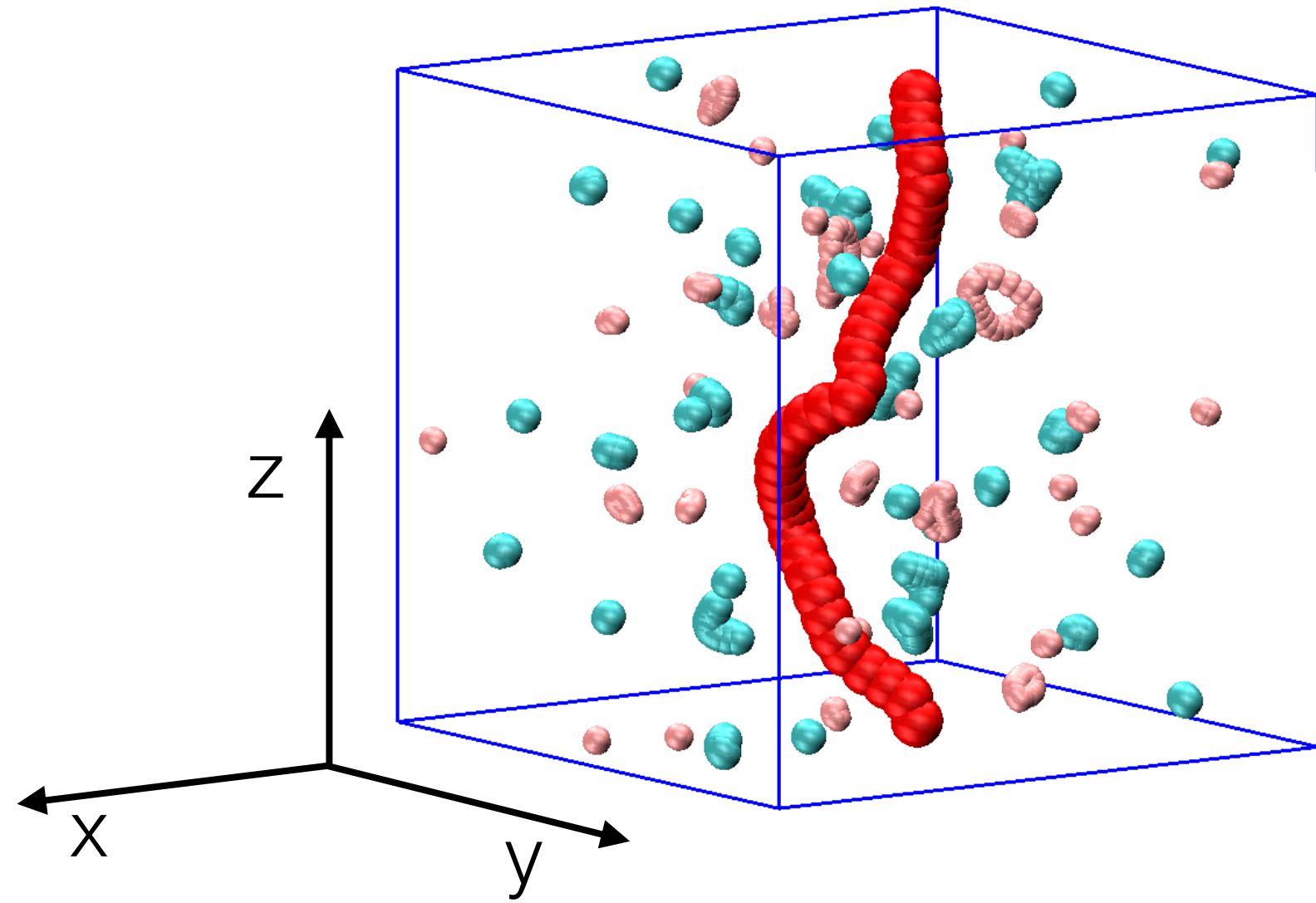


a topologically non-trivial minimum-energy path
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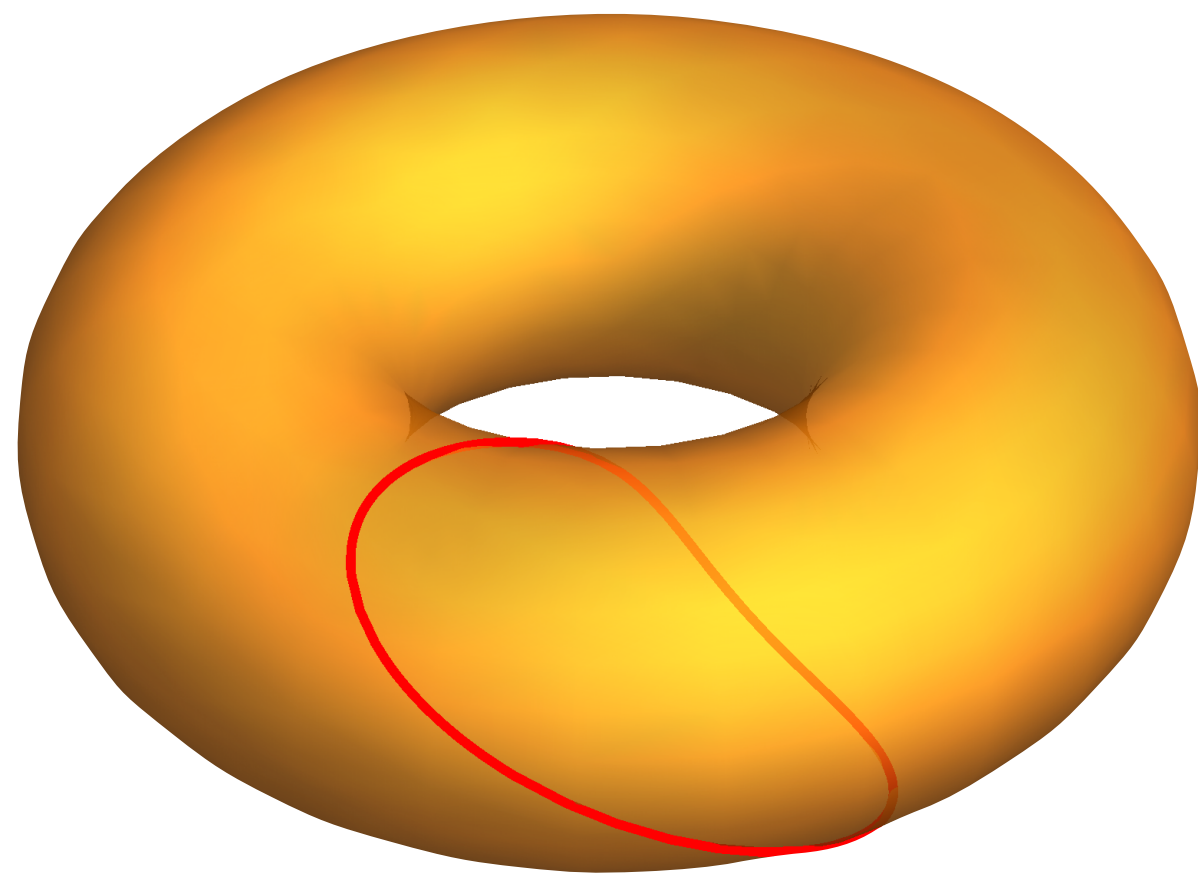
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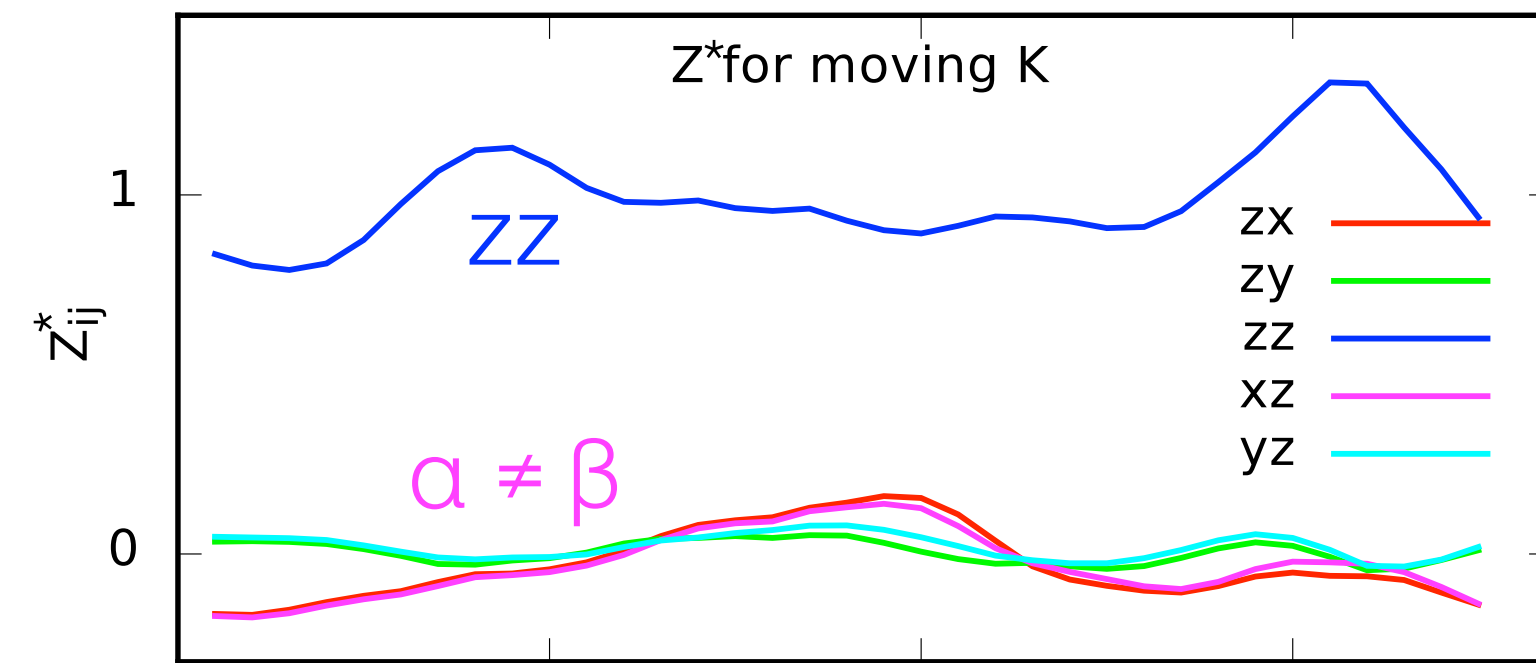
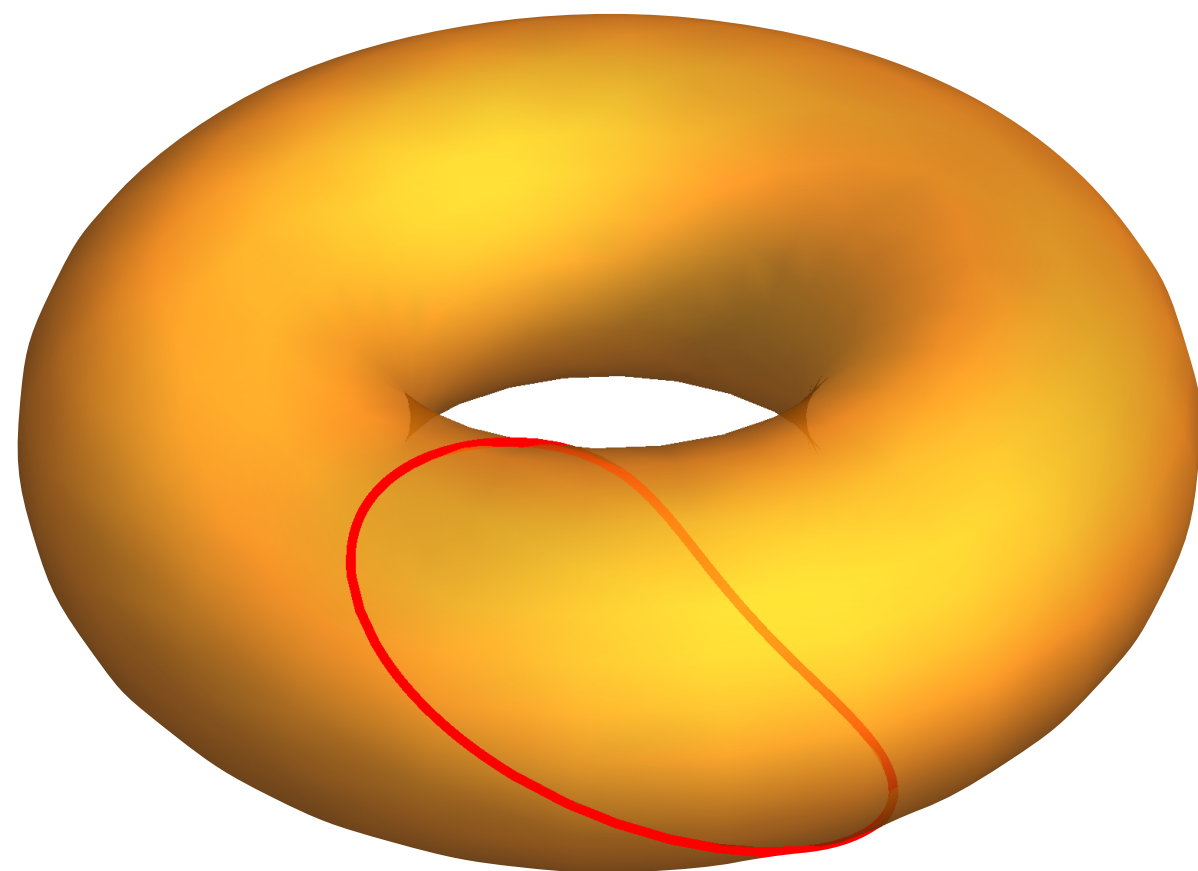
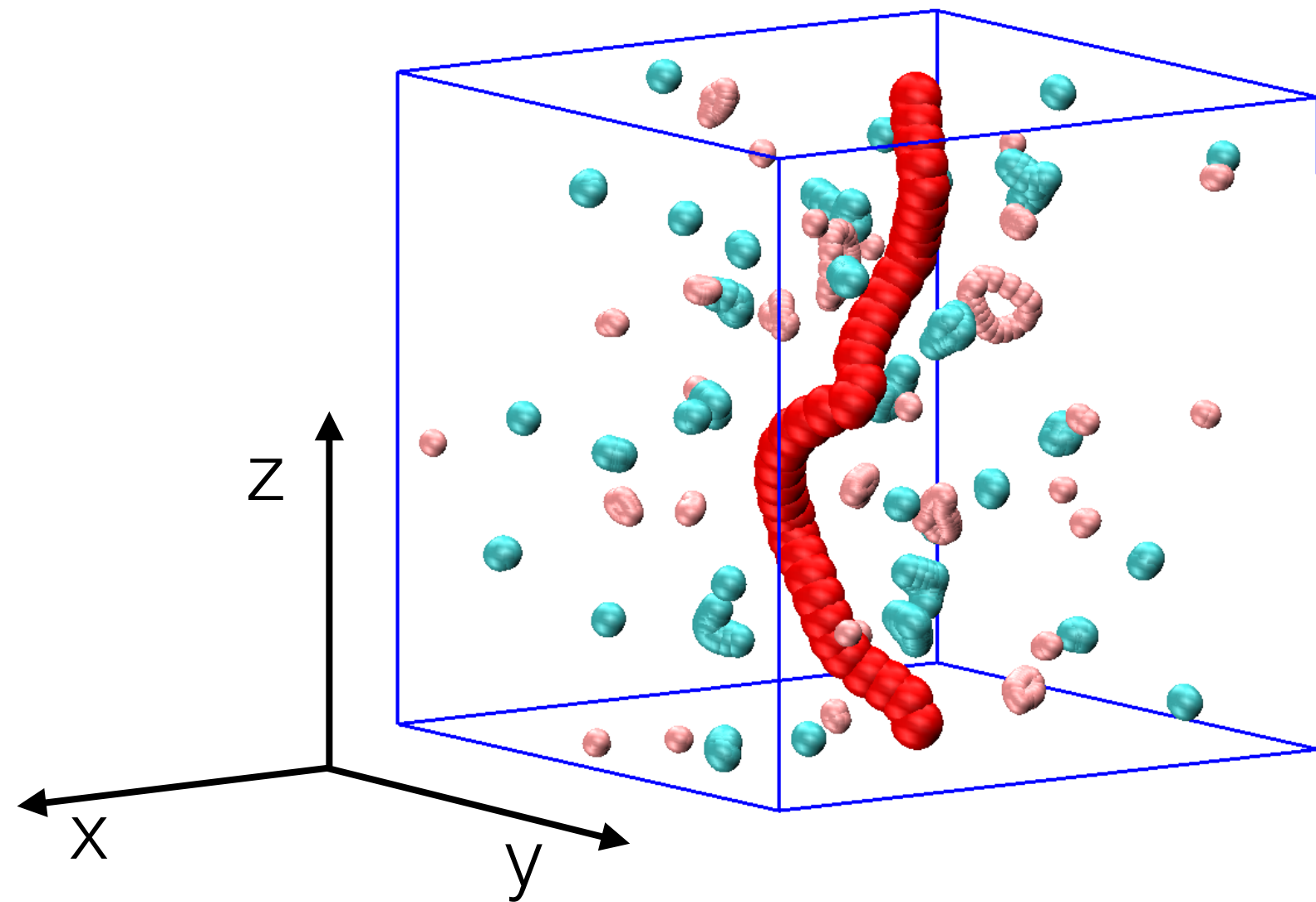
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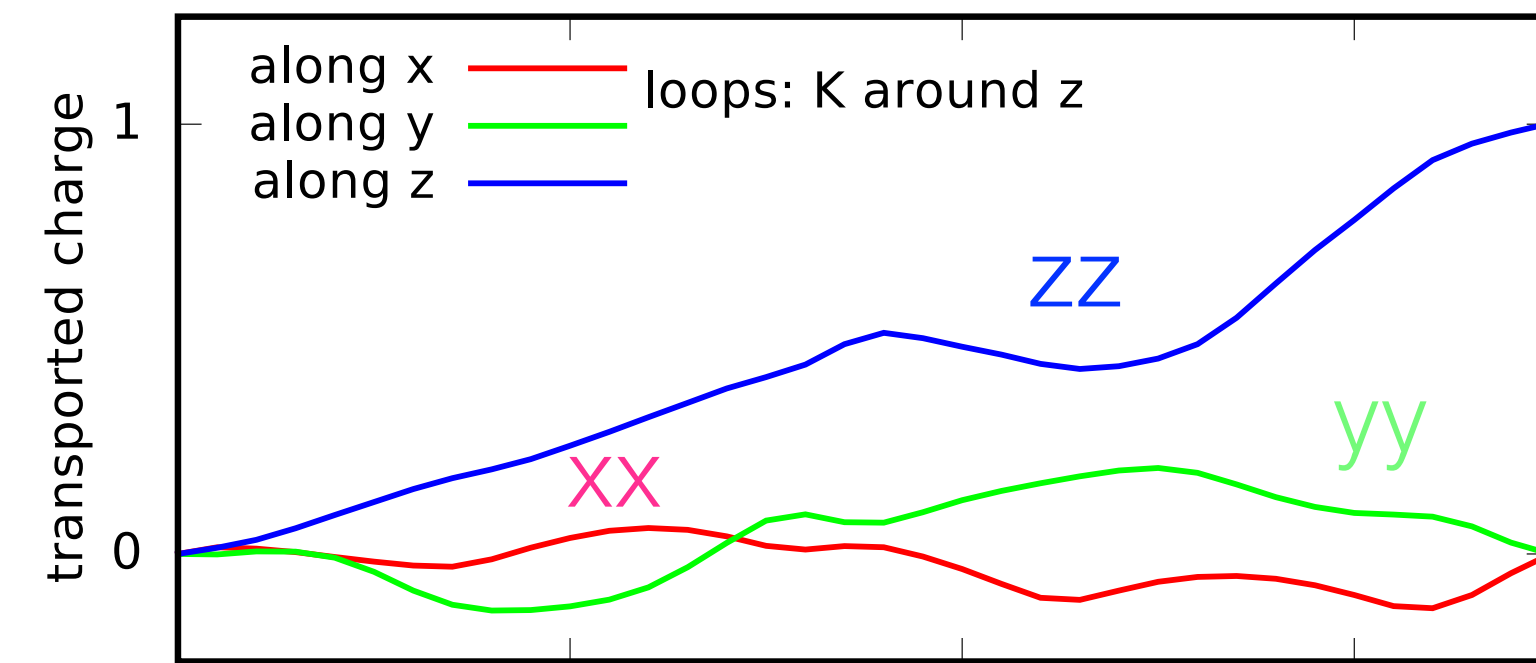
effective charge



a numerical experiment on molten KCl



effective charge

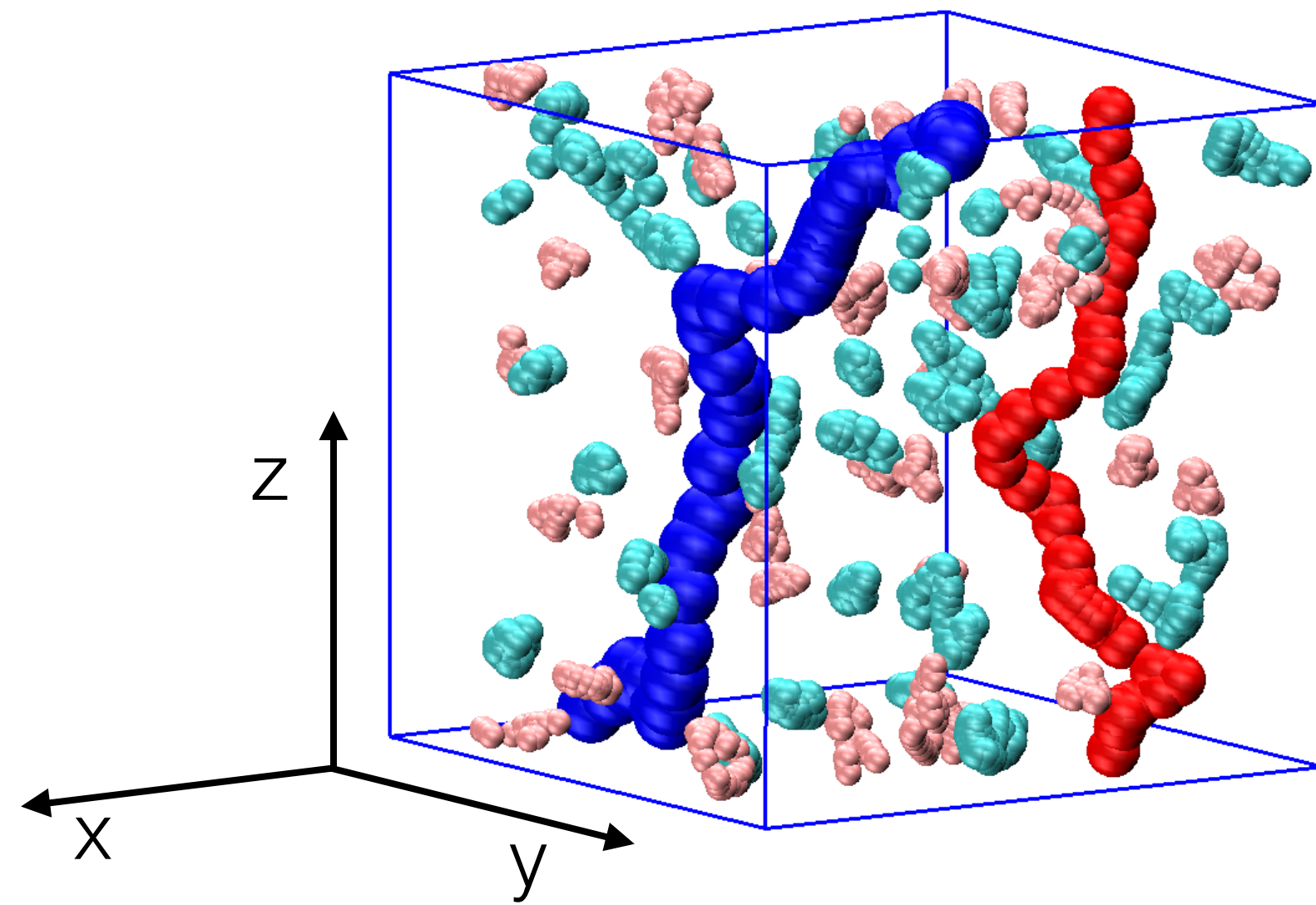


topological charge

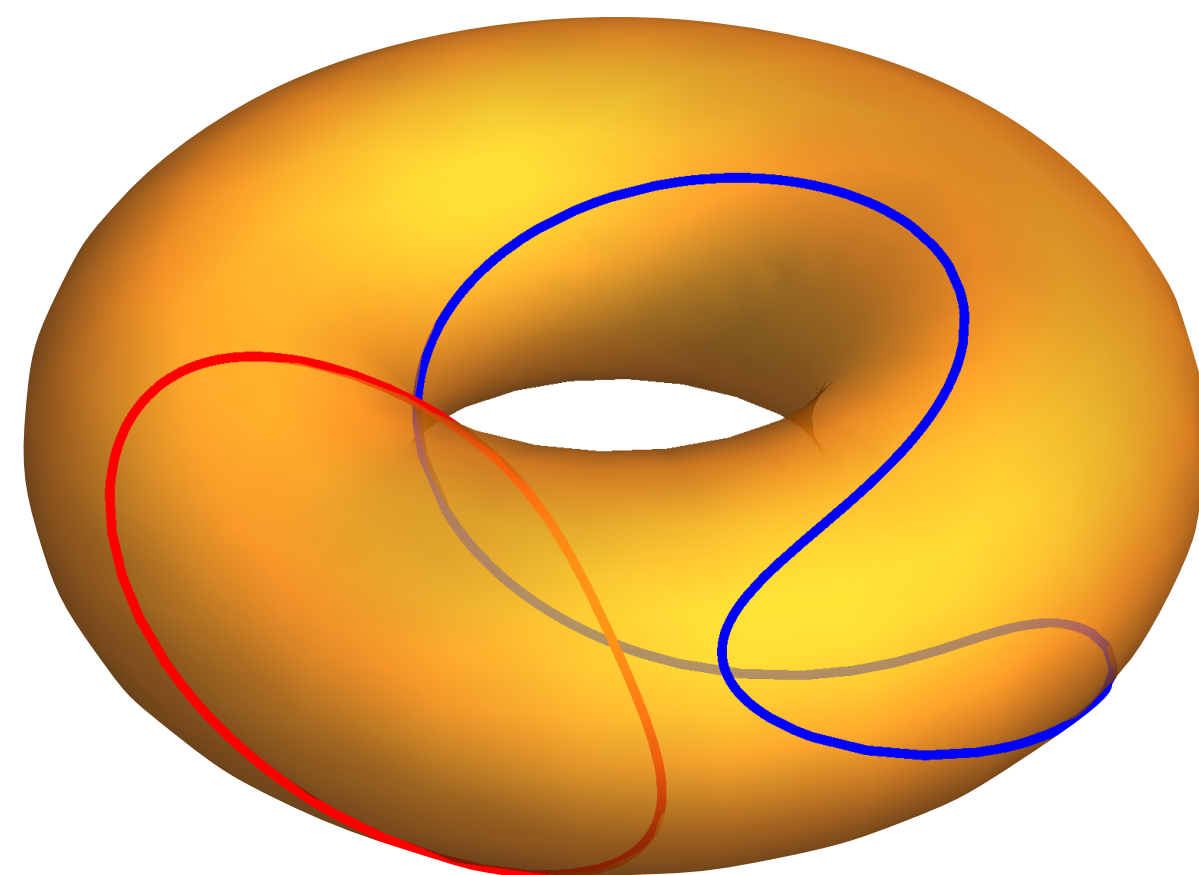
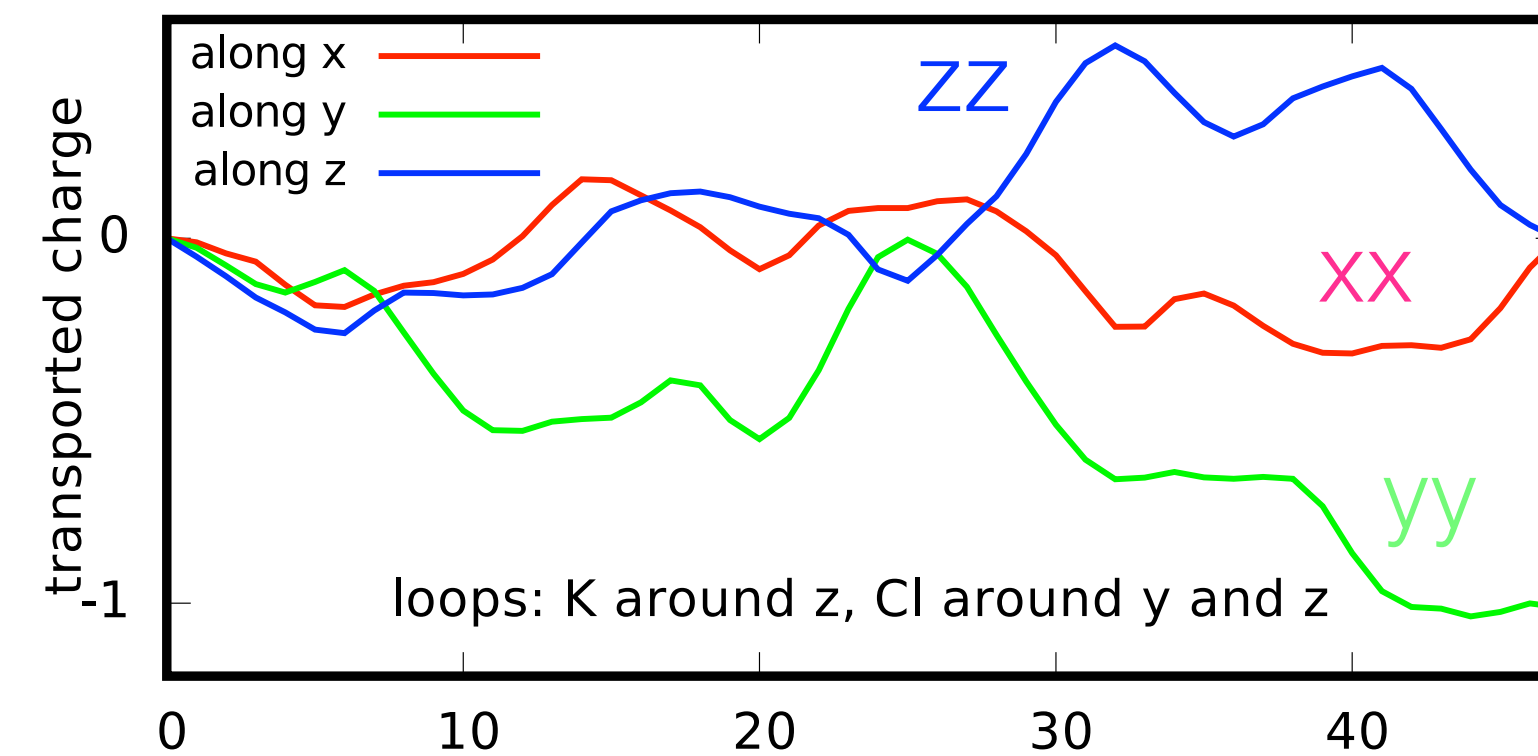
$Q_x = -0.000(6); \quad Q_y = 0.000(2); \quad Q_z = 1.00(18)$



a numerical experiment on molten KCl

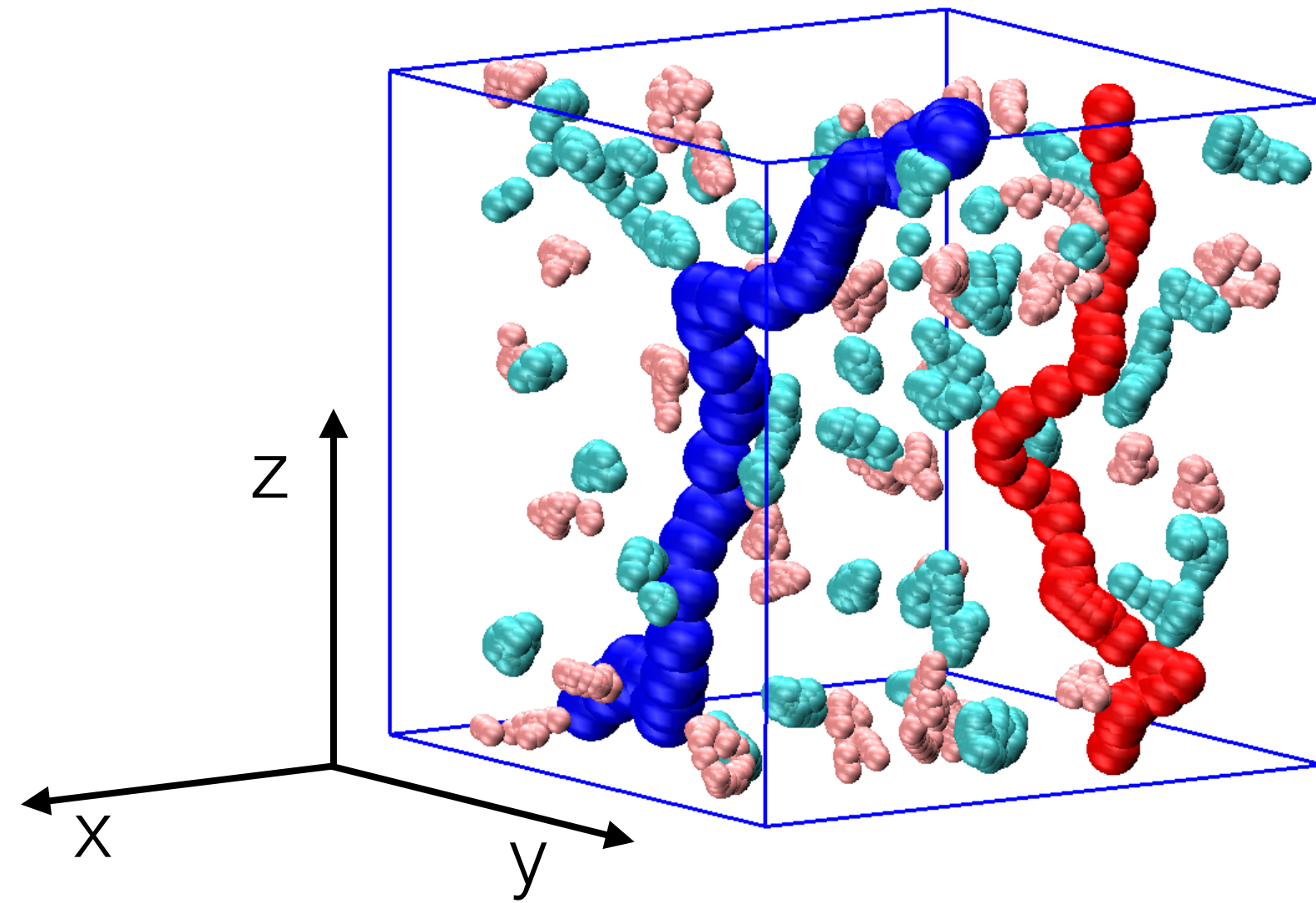


$$\begin{aligned} Q_z[\text{Cl}] &= -1 & Q_y[\text{Cl}] &= -1 \\ Q_z[\text{K}] &= 1 & Q_z[\text{K}] &= 0 \end{aligned}$$

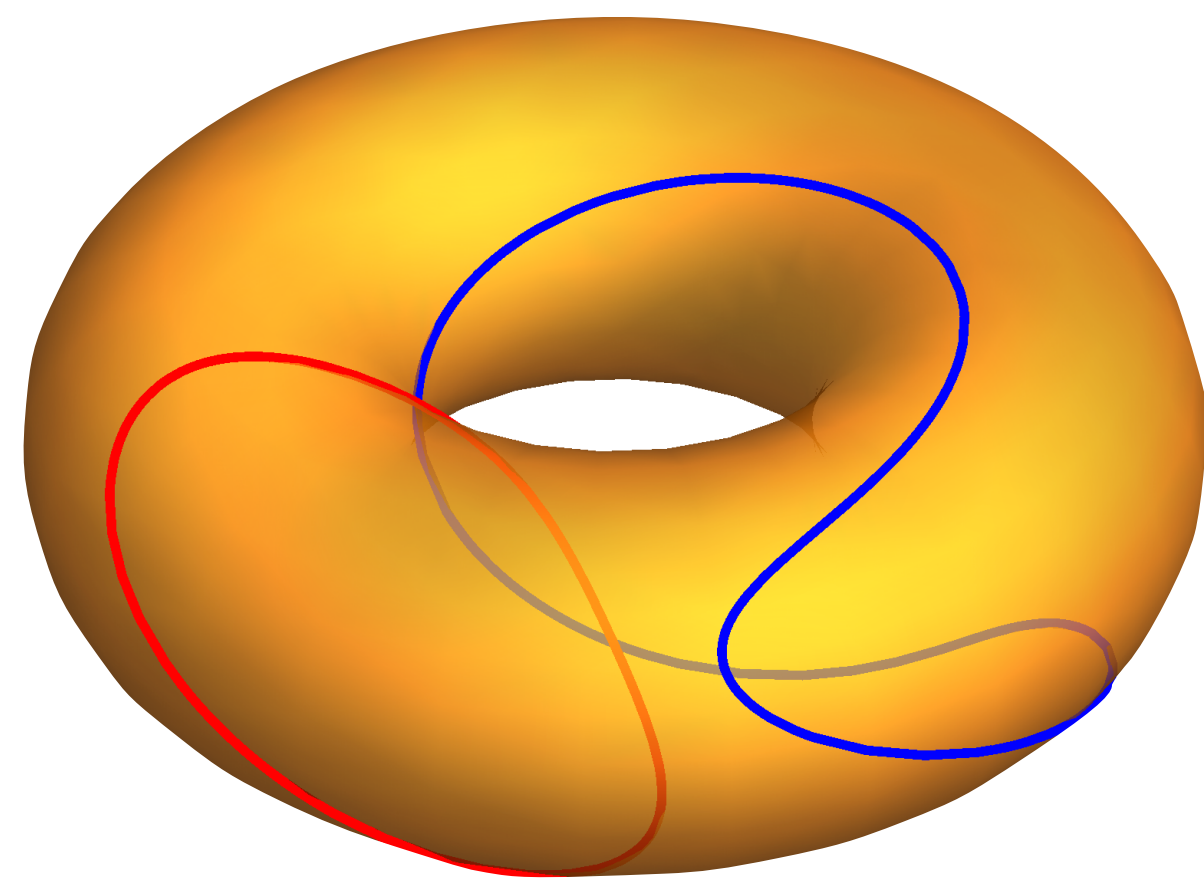
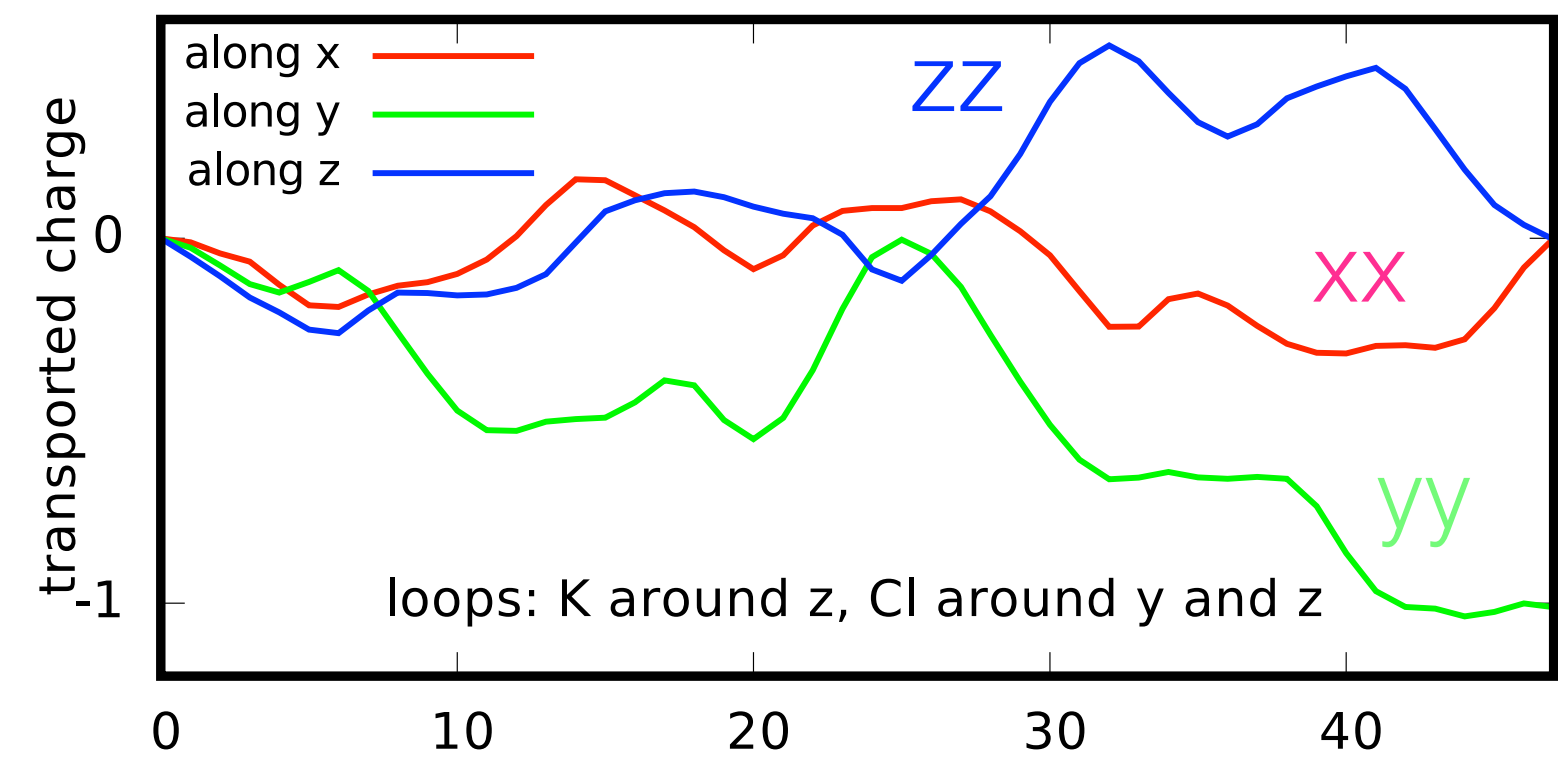


the charges transported by K and Cl
around z cancel exactly

a numerical experiment on molten KCl

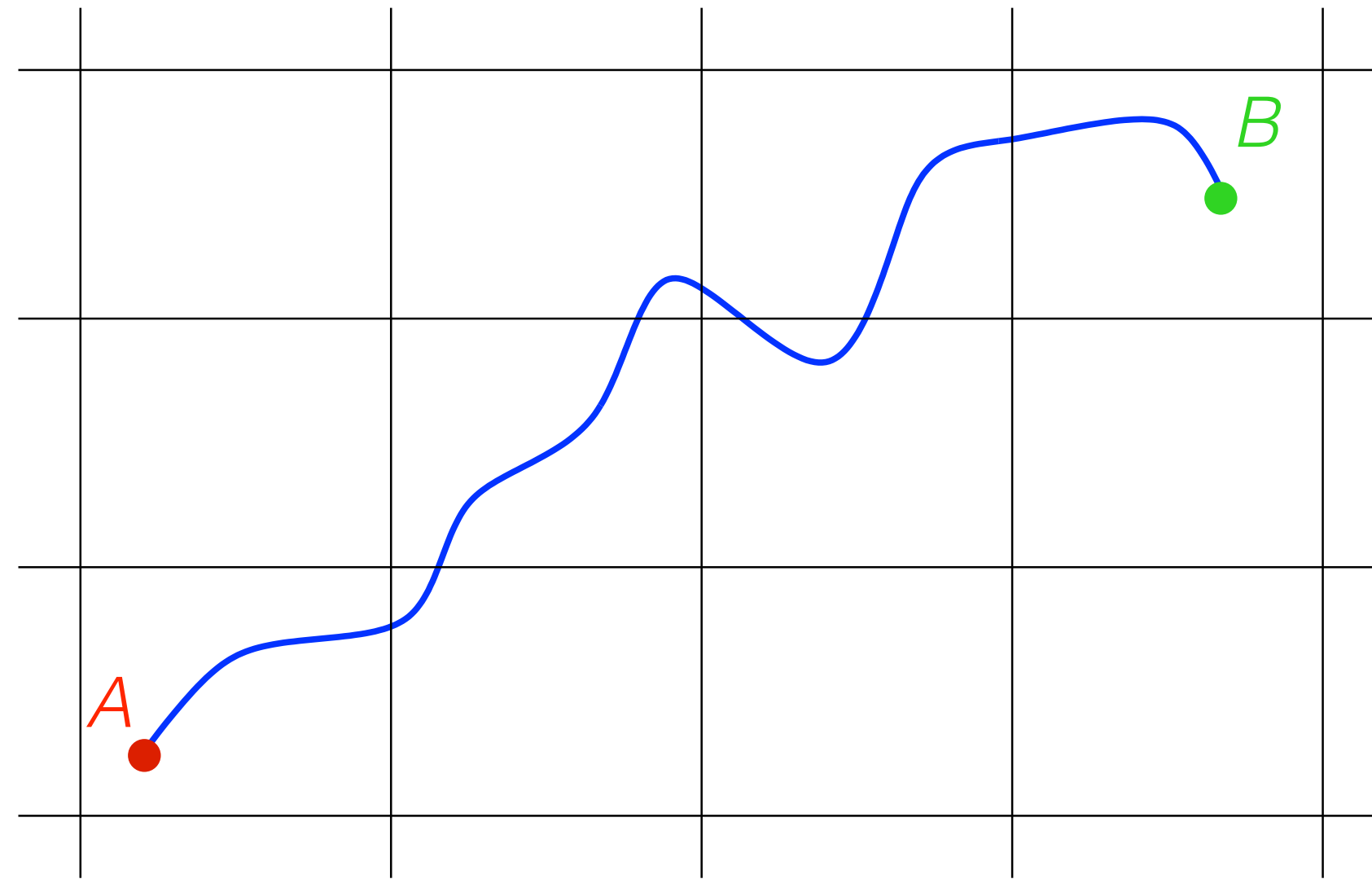


$$\begin{array}{ll} Q_z[\text{Cl}] = -1 & Q_y[\text{Cl}] = -1 \\ Q_z[\text{K}] = 1 & Q_z[\text{K}] = 0 \end{array}$$



the charges transported by K and Cl
around z cancel exactly

gauge invariance of charge transport

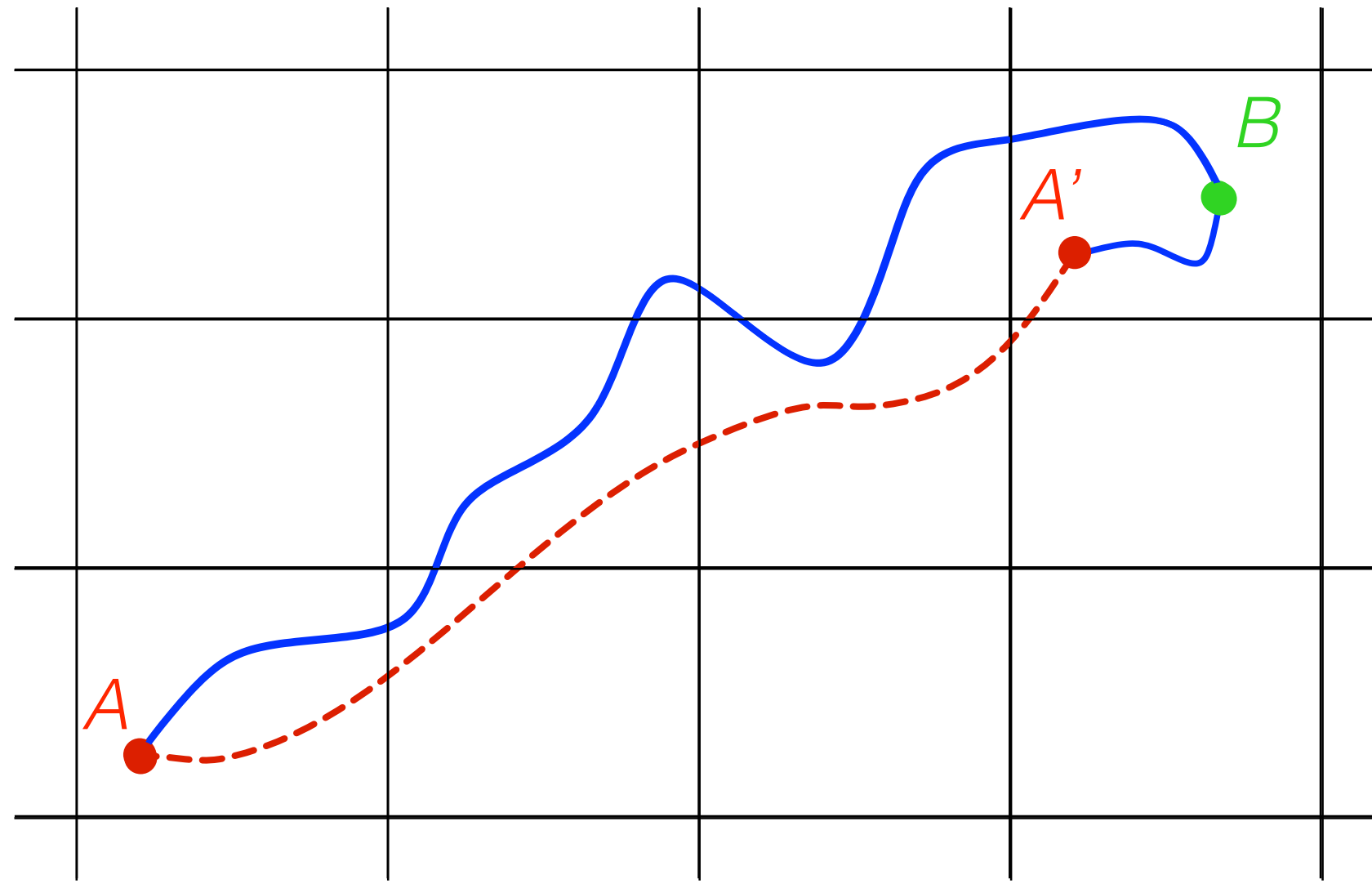


$$\sigma \propto \lim_{t \rightarrow \infty} \frac{1}{2t} \text{var} [\mu_{AB}(t)]$$
$$\mu_{AB}(t) = \int_0^t J(t') dt'$$

gauge invariance of charge transport

$$\hat{H}(B) \neq \hat{H}(A)$$

$$\hat{H}(A') = \hat{H}(A)$$



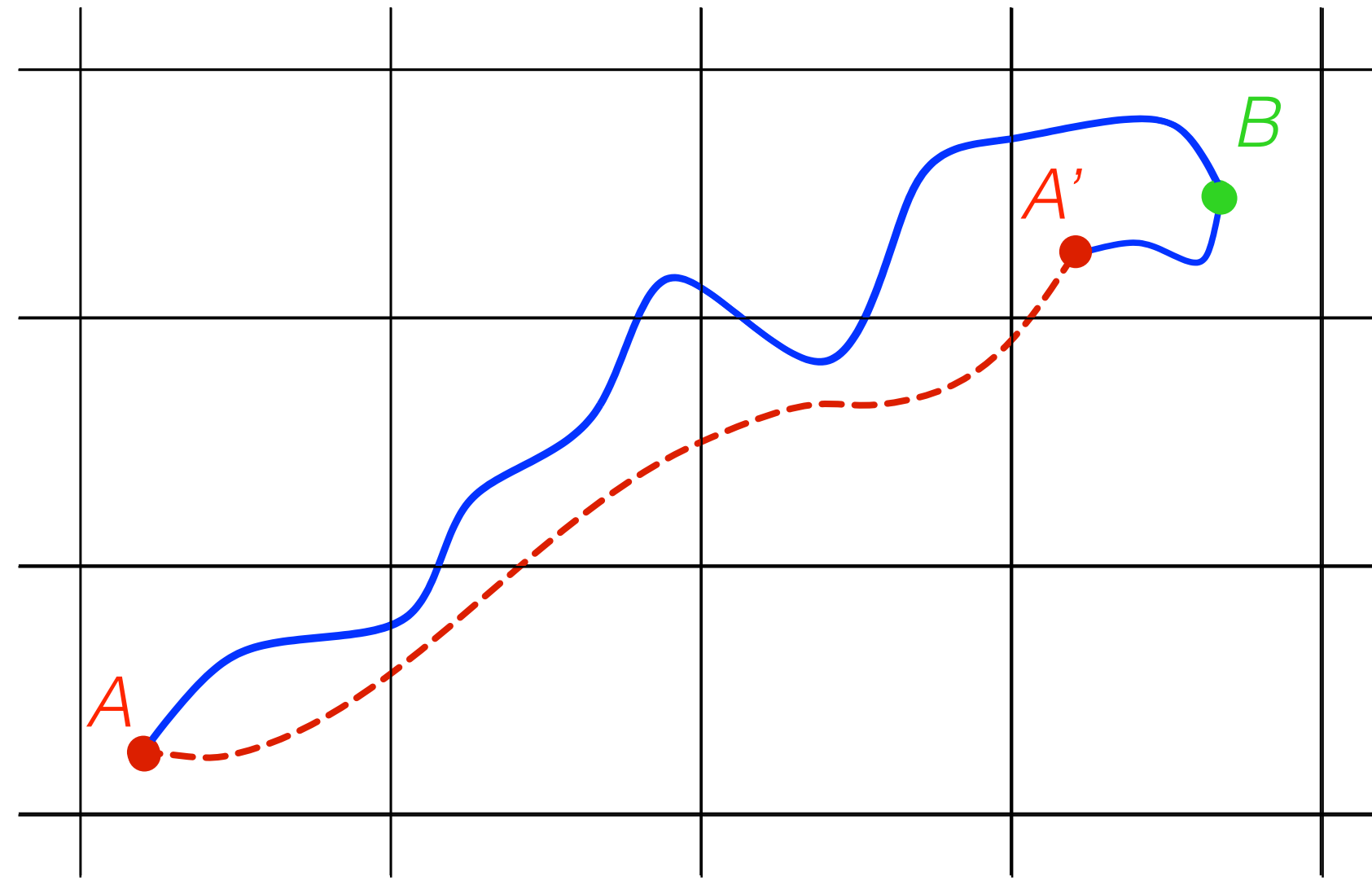
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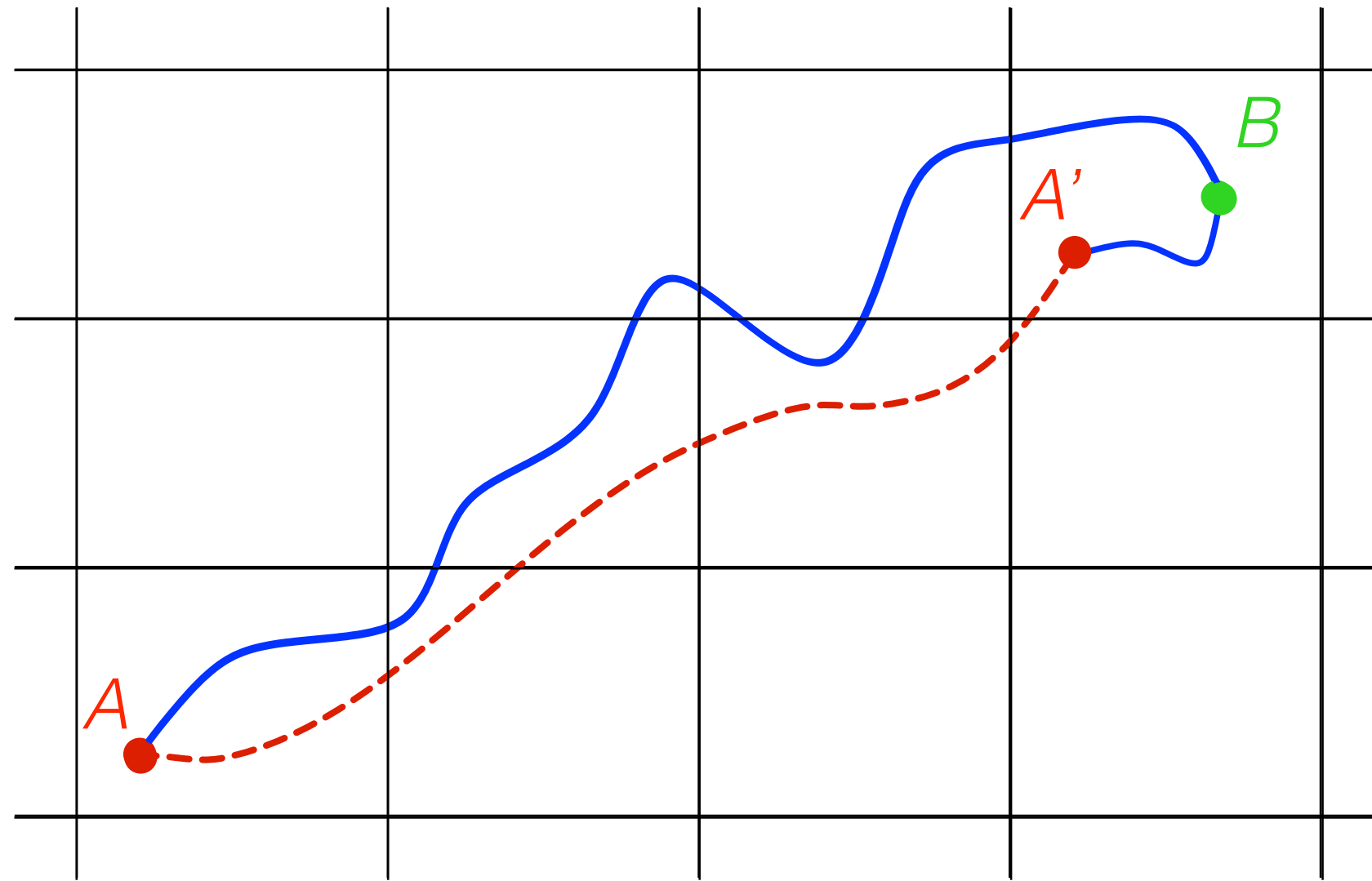


$$\sigma \propto \lim_{t \rightarrow \infty} \frac{1}{2t} \text{var}[\mu_{AB}(t)]$$

$$\begin{aligned} \mu_{AB}(t) &= \int_0^t J(t') dt' \\ &= \mu_{AA'} + \mu_{A'B} \end{aligned}$$

gauge invariance of charge transport

$$\hat{H}(B) \neq \hat{H}(A)$$
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$$\sigma \propto \lim_{t \rightarrow \infty} \left(\frac{1}{2t} \text{var} [\mu_{AB}(t)] \right)$$
$$\mu_{AB}(t) = \int_0^t J(t') dt'$$
$$= \mu_{AA'} + \cancel{\mu_{A'B}}$$

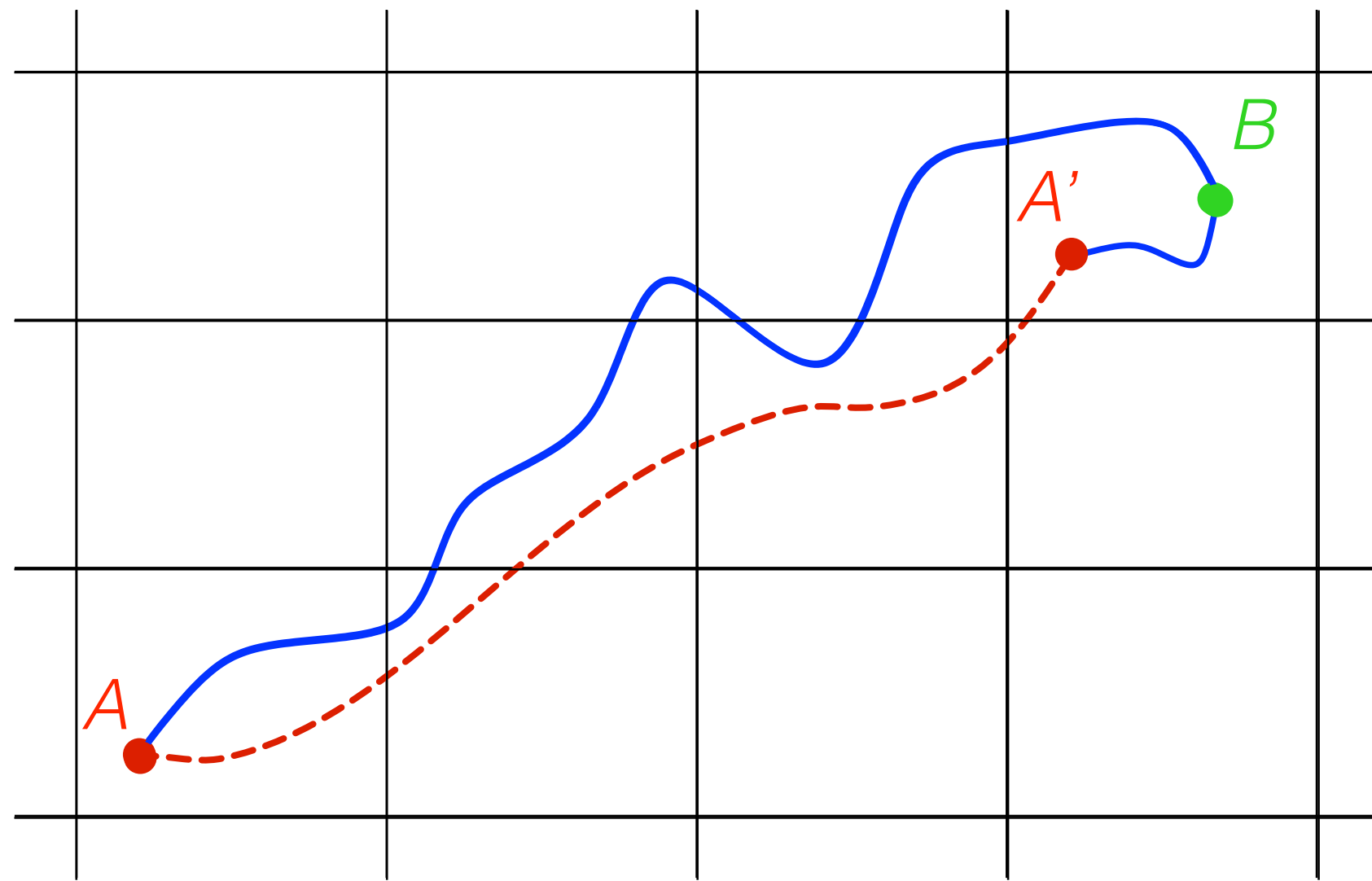
$$\mu_{AB}(t) = \mu_{AA'} + \mathcal{O}(1)$$



gauge invariance of charge transport

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$$\mu_{AB}(t) = \int_0^t J(t') dt'$$

$$= \mu_{AA'} + \cancel{\mu_{A'B}}$$

$$J_\alpha = \sum_{i\beta} Z_{i\alpha\beta}^* v_{i\beta}$$

$$J'_\alpha = \sum_i q_{S(i)} v_{i\alpha}$$

$$\mu_{AB}(t) = \mu_{AA'} + \mathcal{O}(1)$$

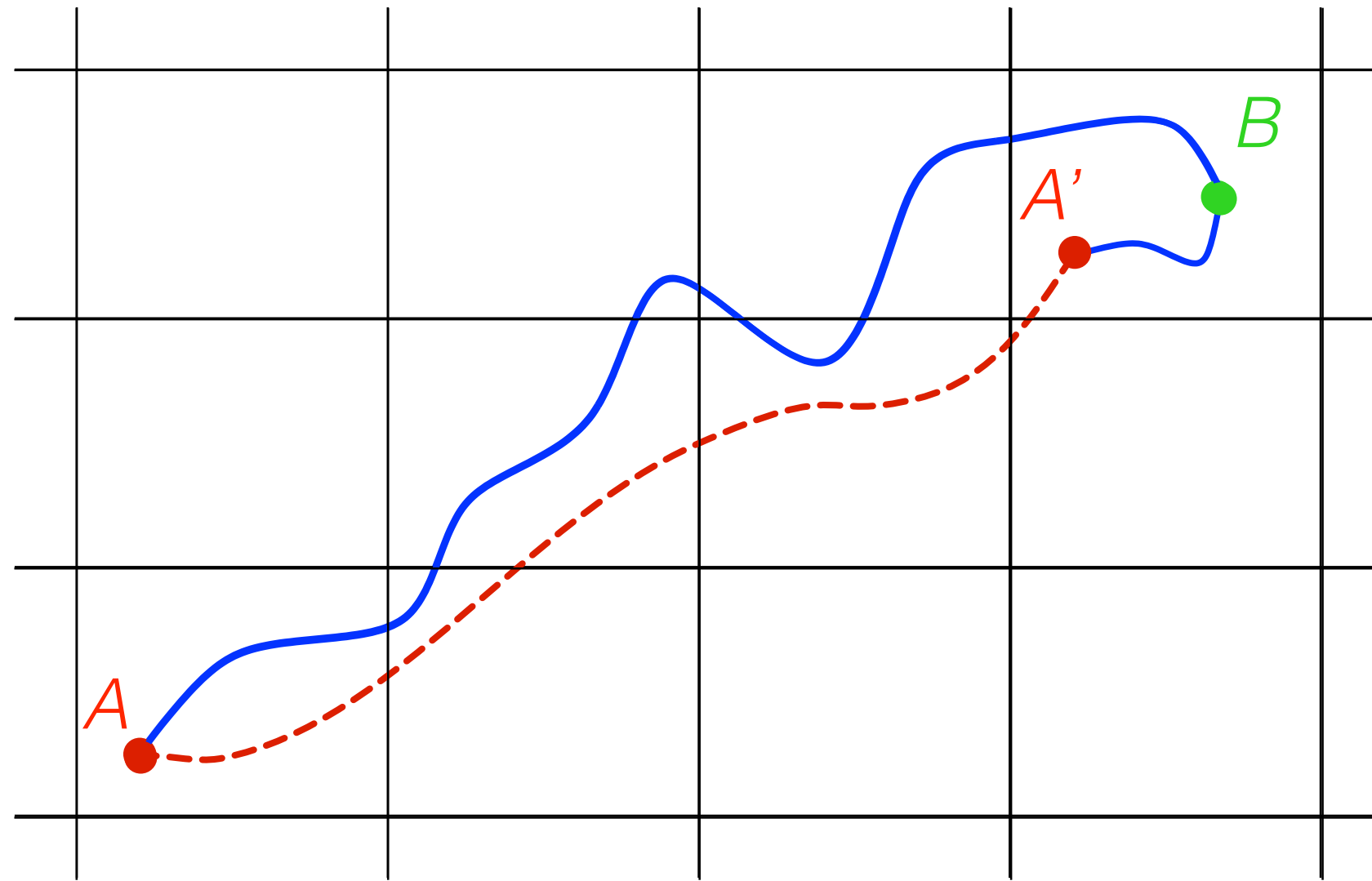
$$= \mu'_{AA'} + \mathcal{O}(1) \quad (\text{Thouless})$$



gauge invariance of charge transport

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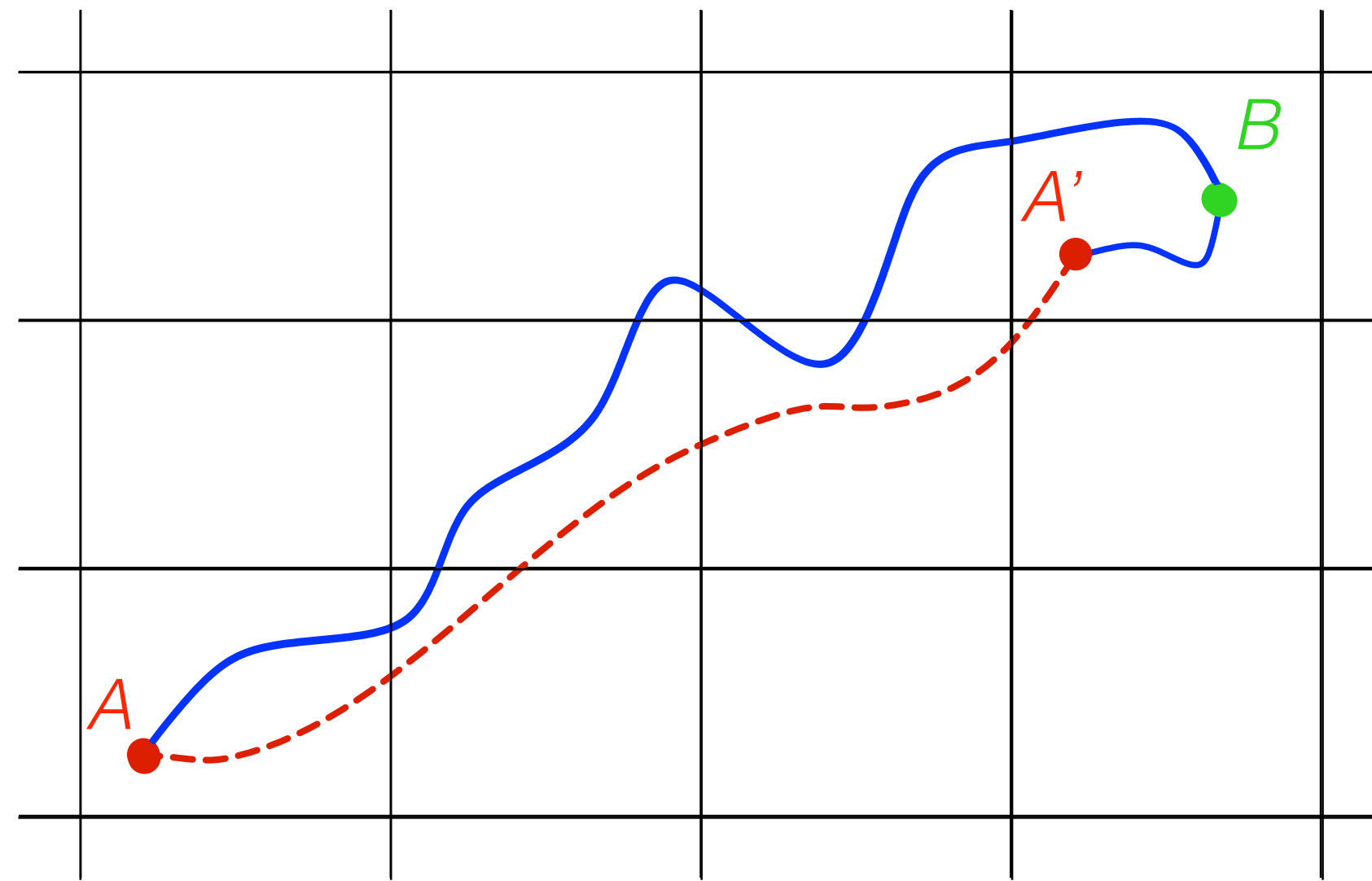
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gauge invariance of charge transport

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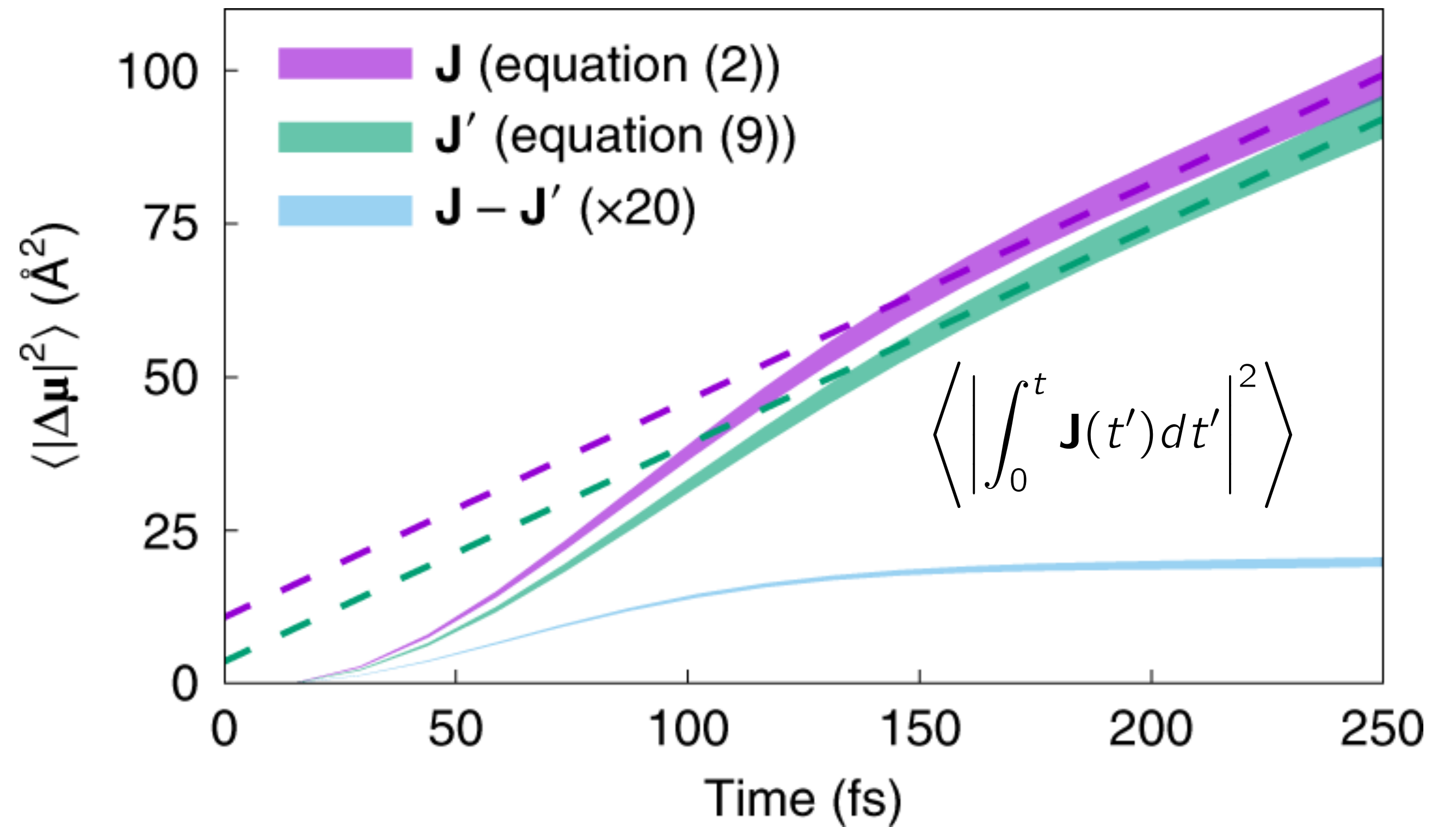
$$\sigma = \sigma'$$



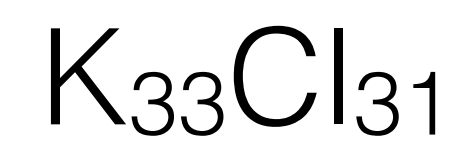
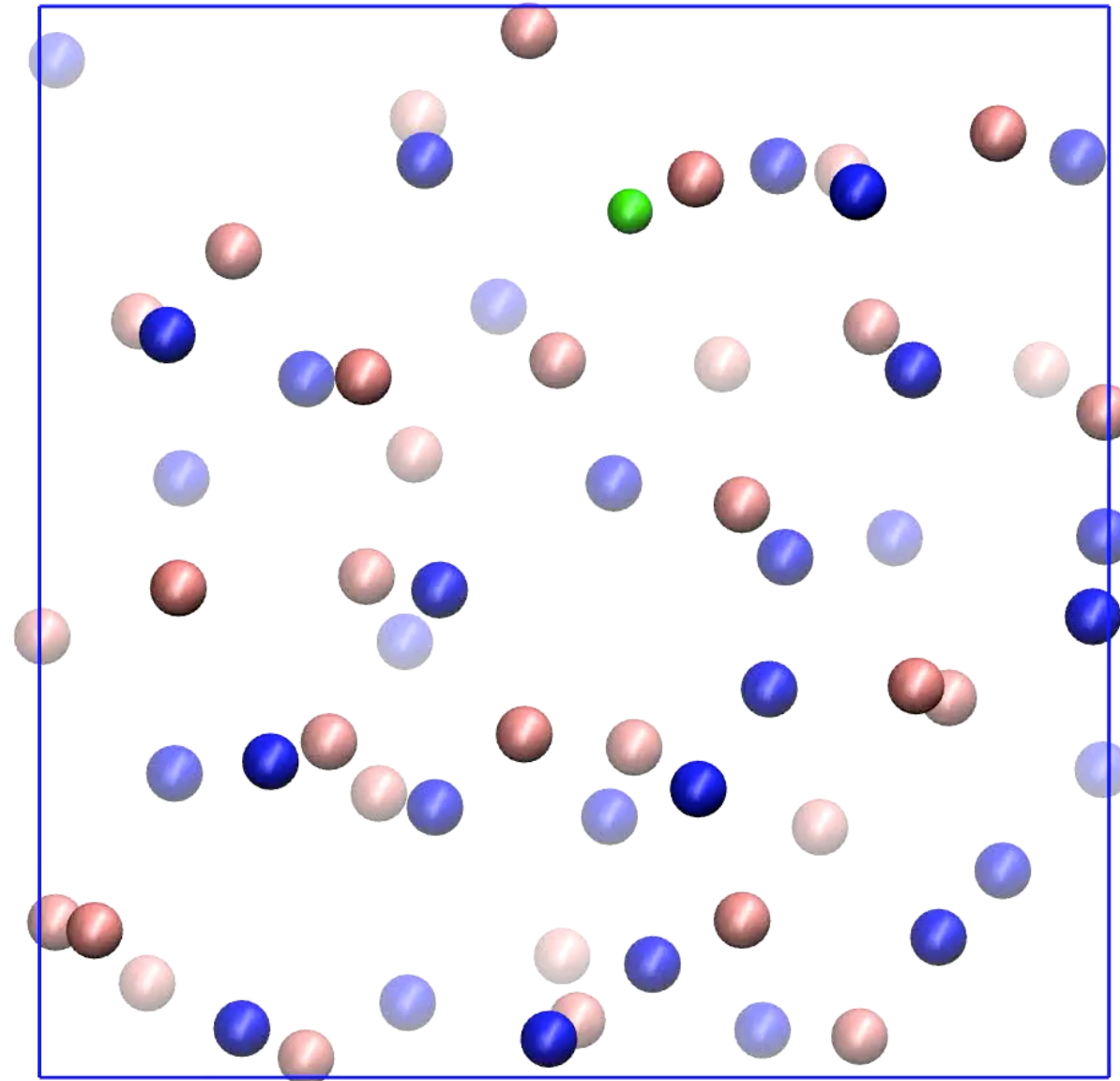
currents from atomic oxidation numbers

$$(2) \quad J_{\alpha} = \sum_{i\beta} Z_{i\alpha\beta}^* V_{i\beta}$$

$$(9) \quad J'_{\alpha} = \sum_i q_{S(i)} V_{i\alpha}$$

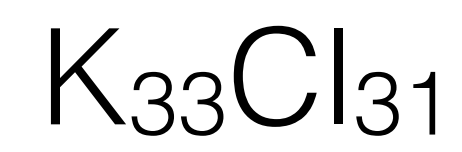
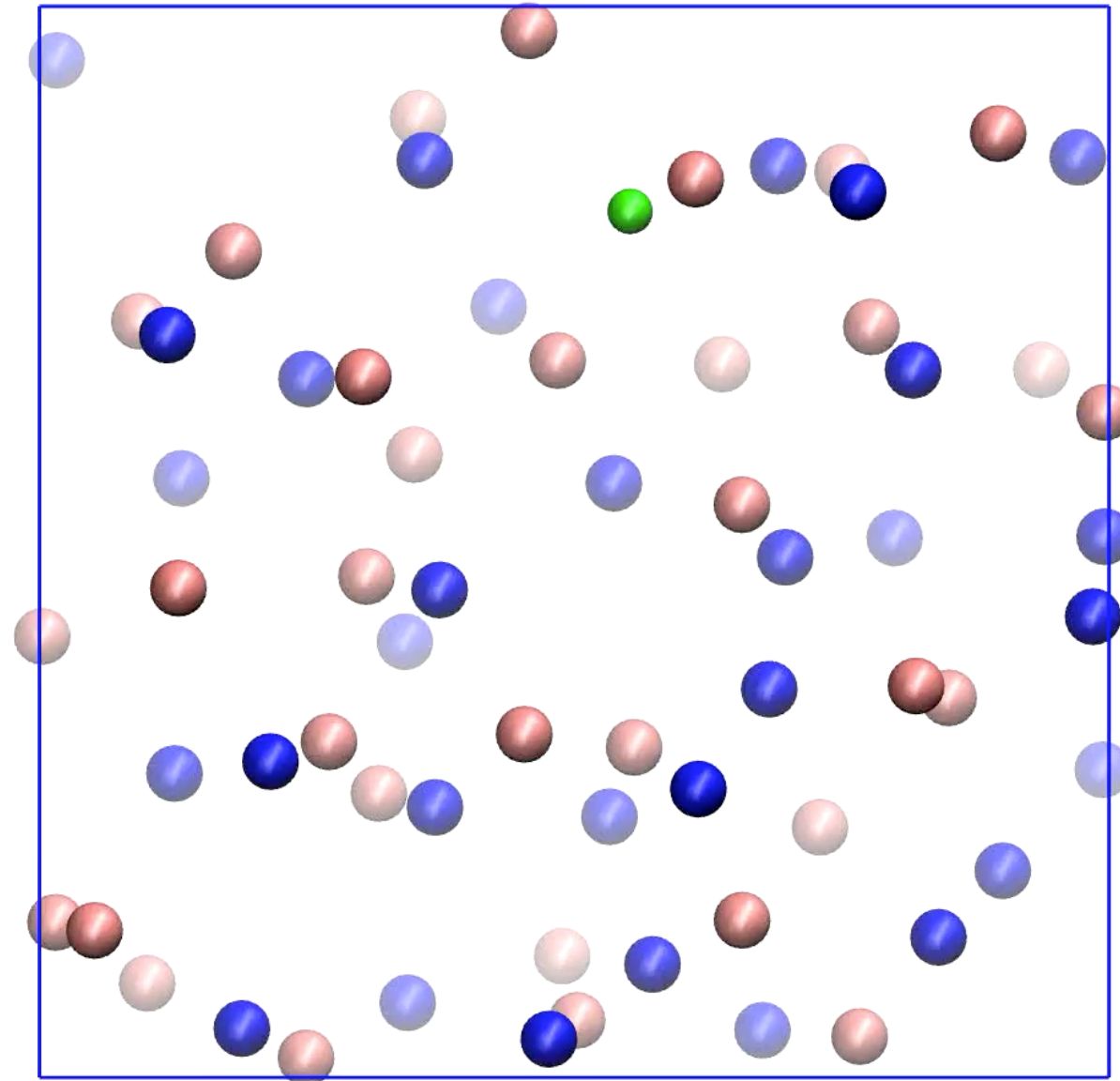


non-stoichiometric melts



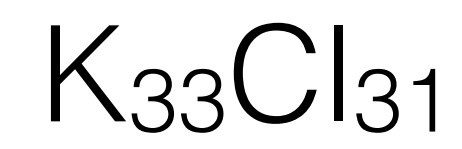
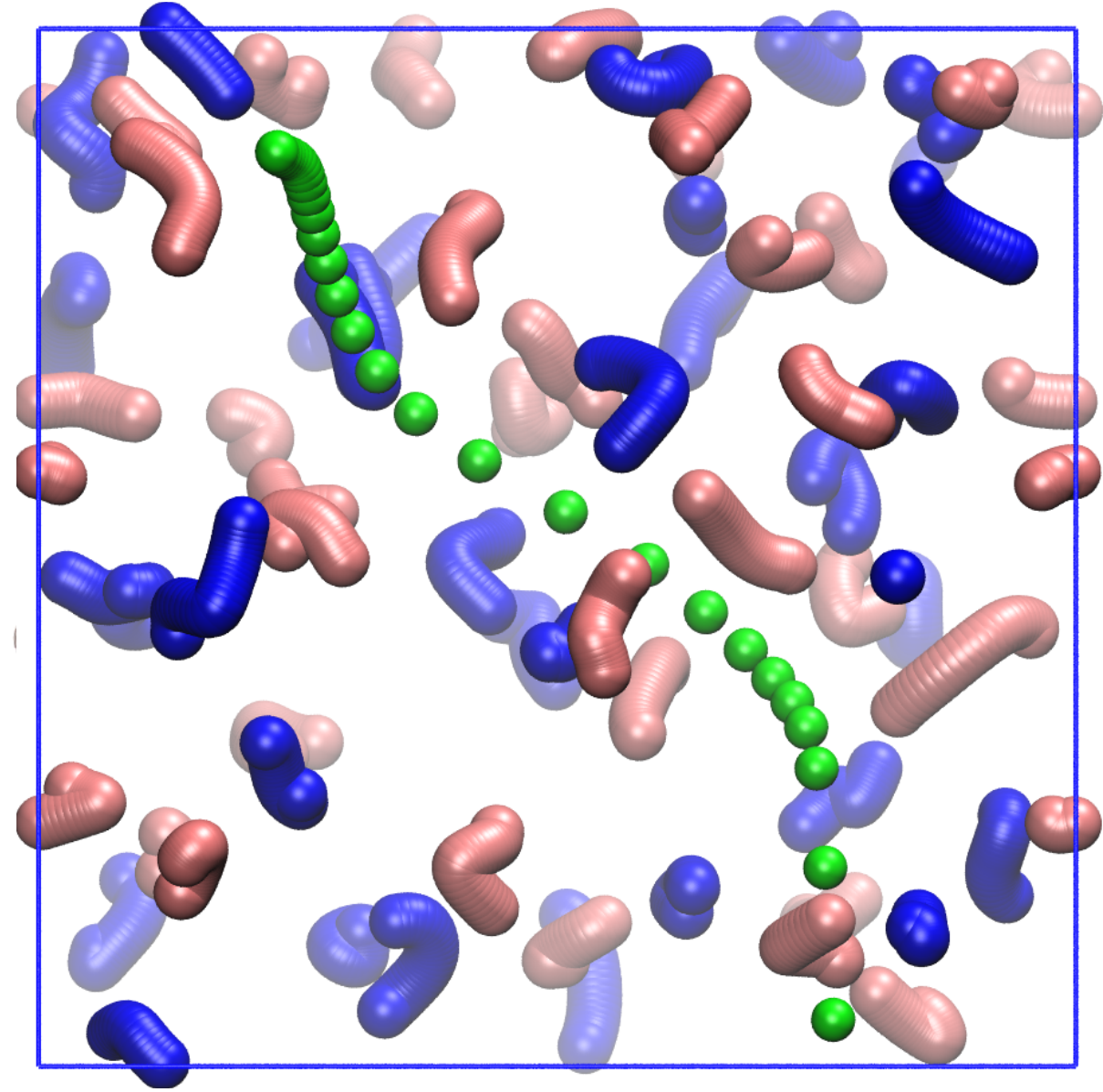
$$x \approx 0.06$$

non-stoichiometric melts



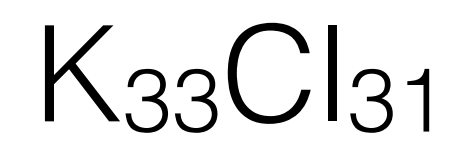
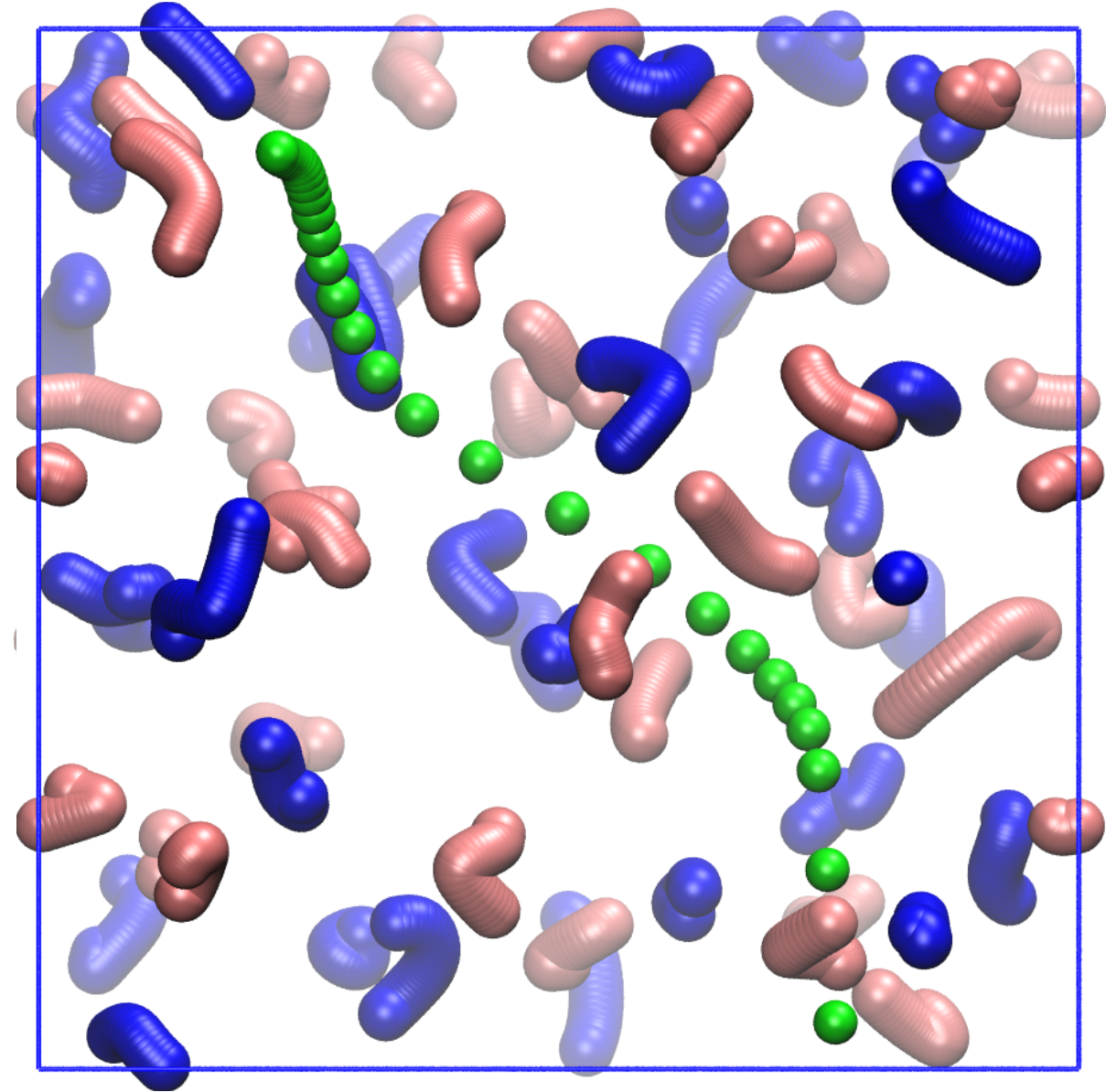
$$x \approx 0.06$$

non-stoichiometric melts

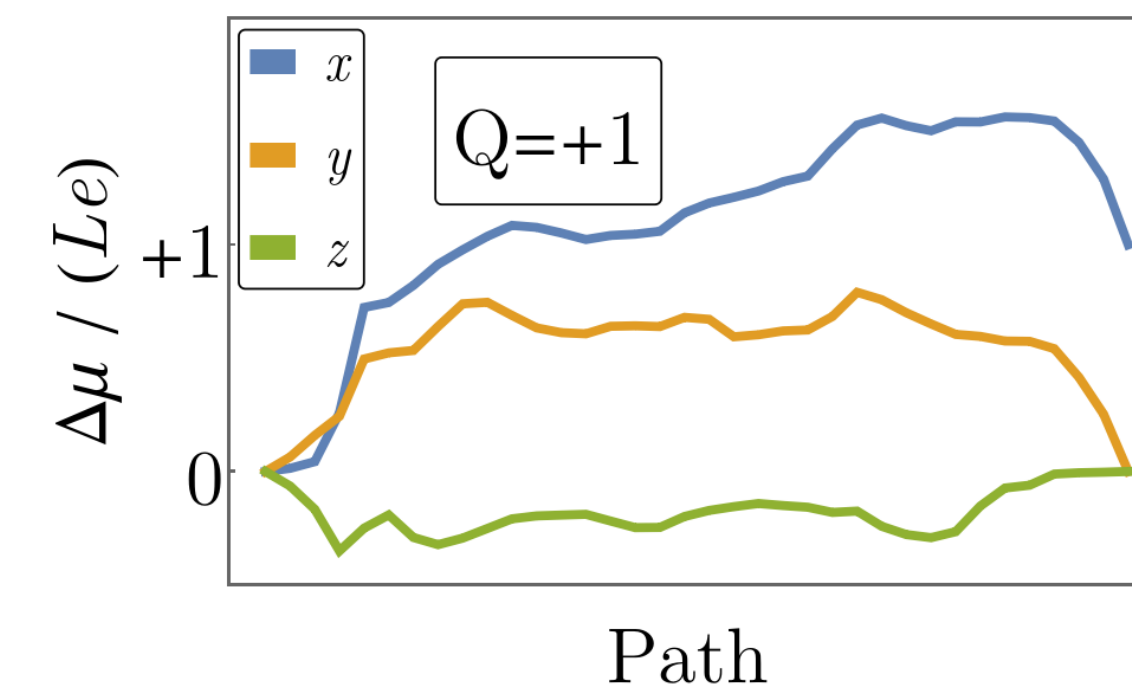
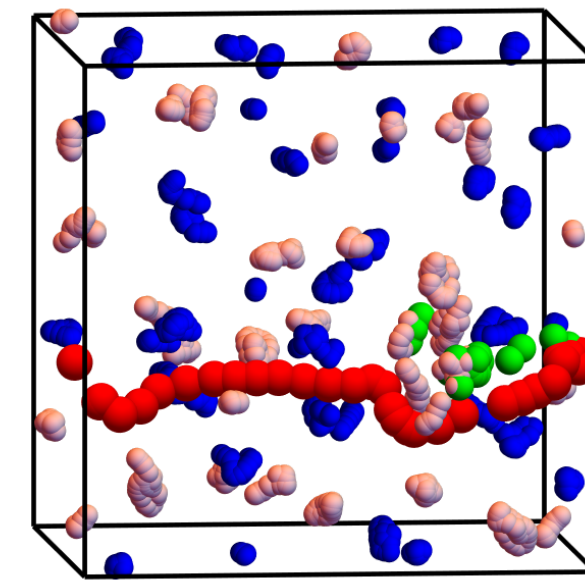


$$x \approx 0.06$$

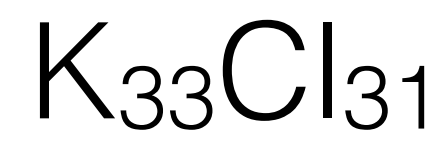
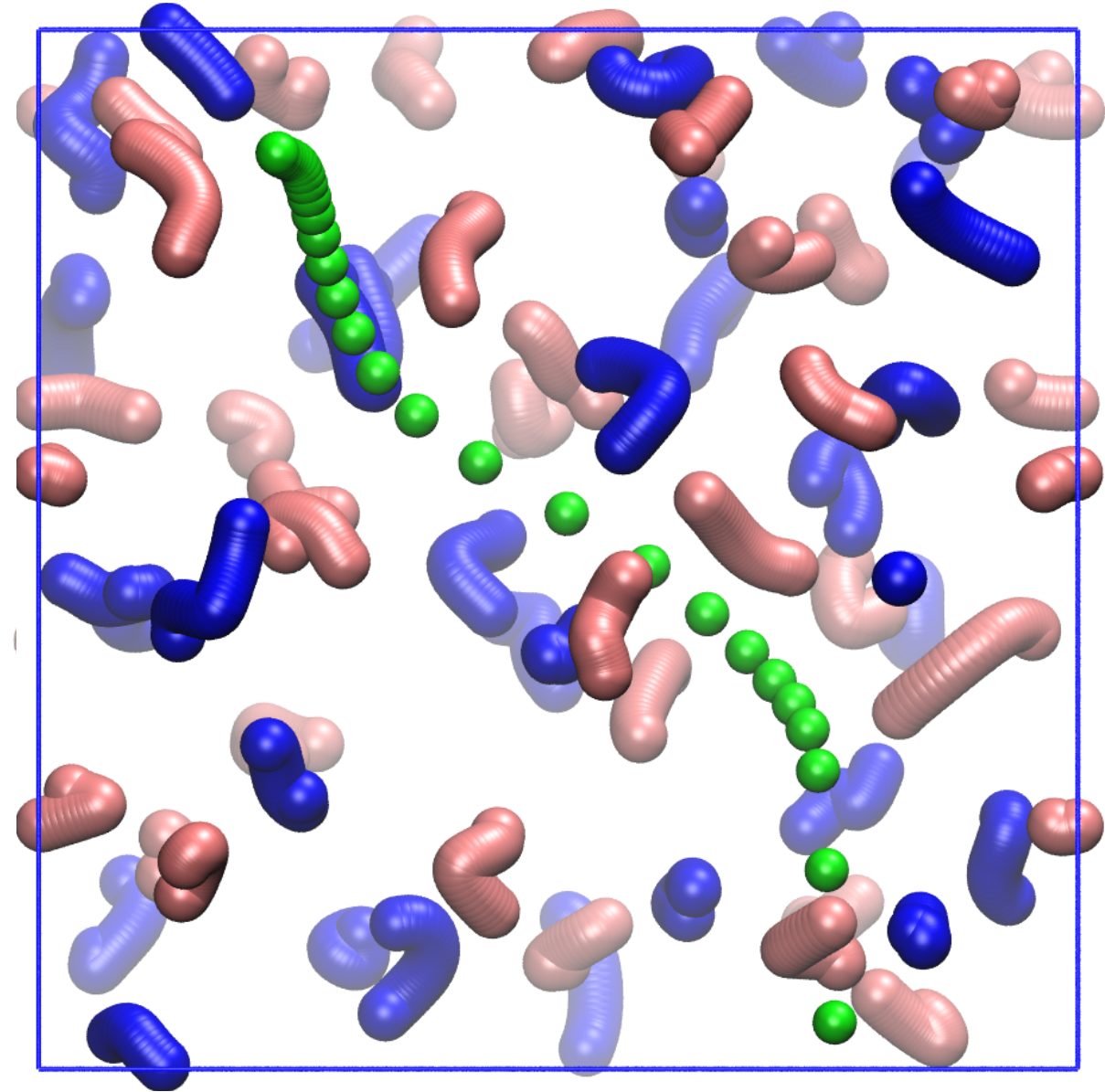
non-stoichiometric melts



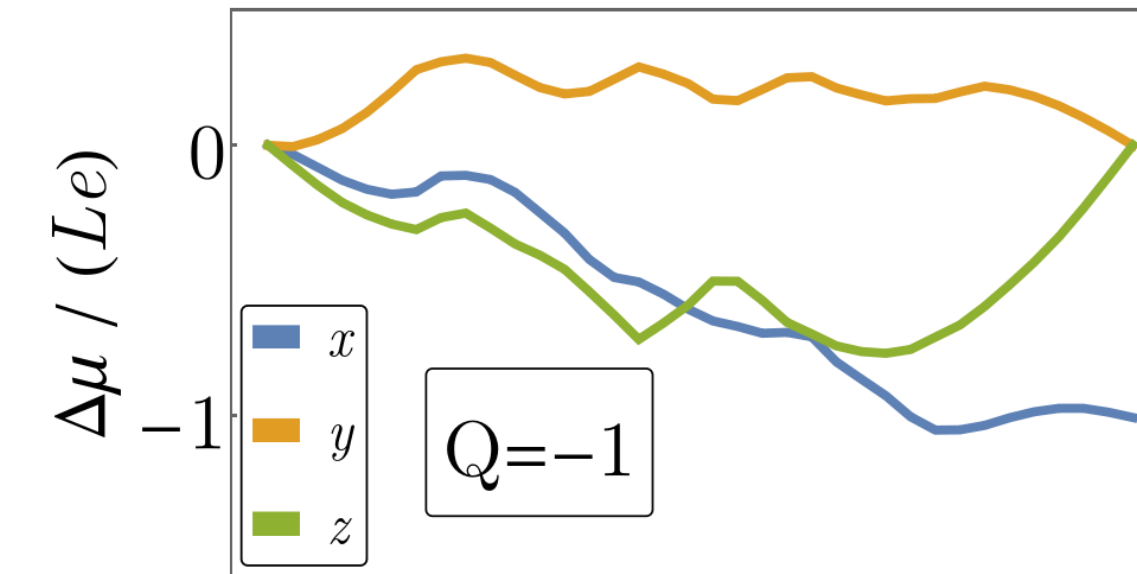
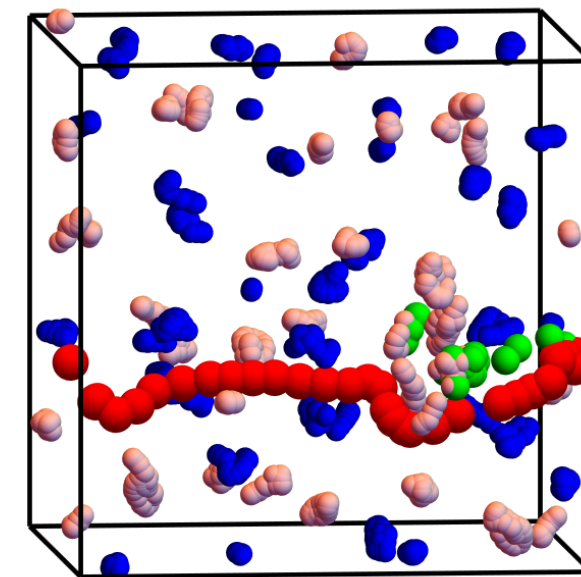
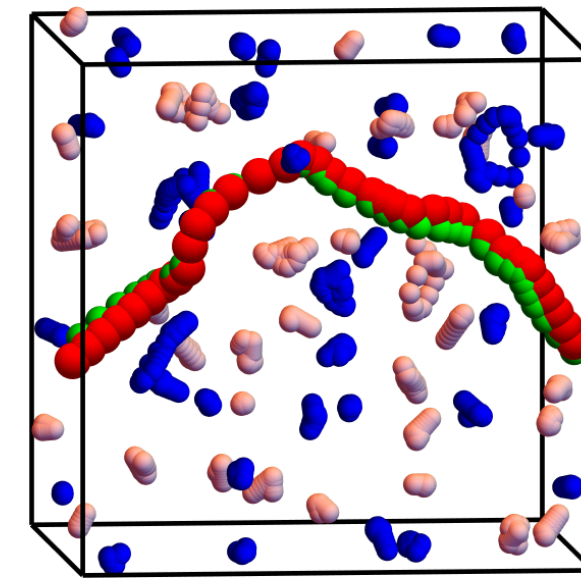
$$x \approx 0.06$$



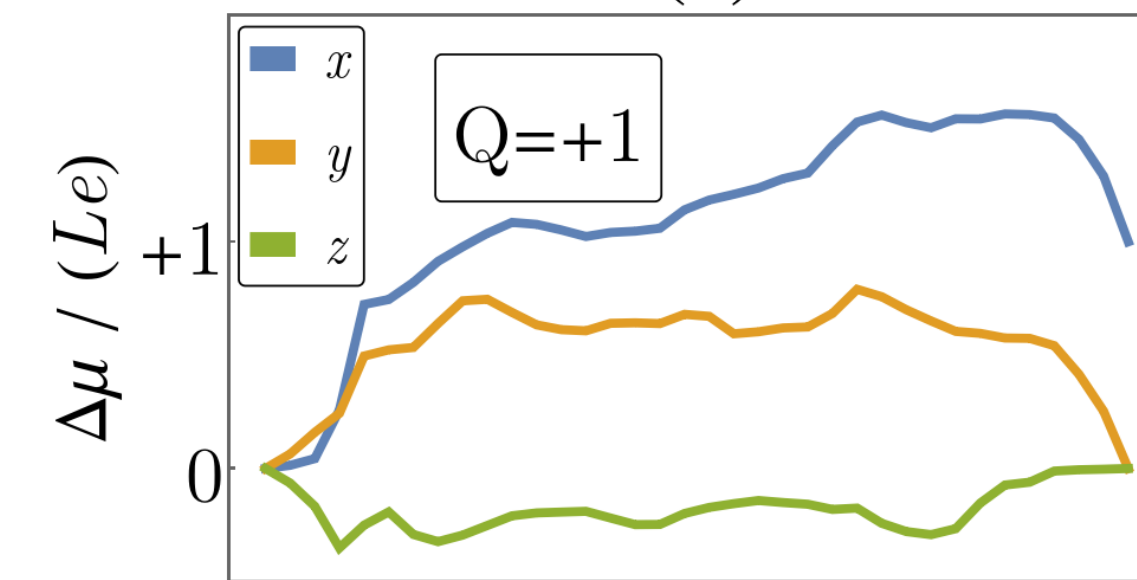
non-stoichiometric melts



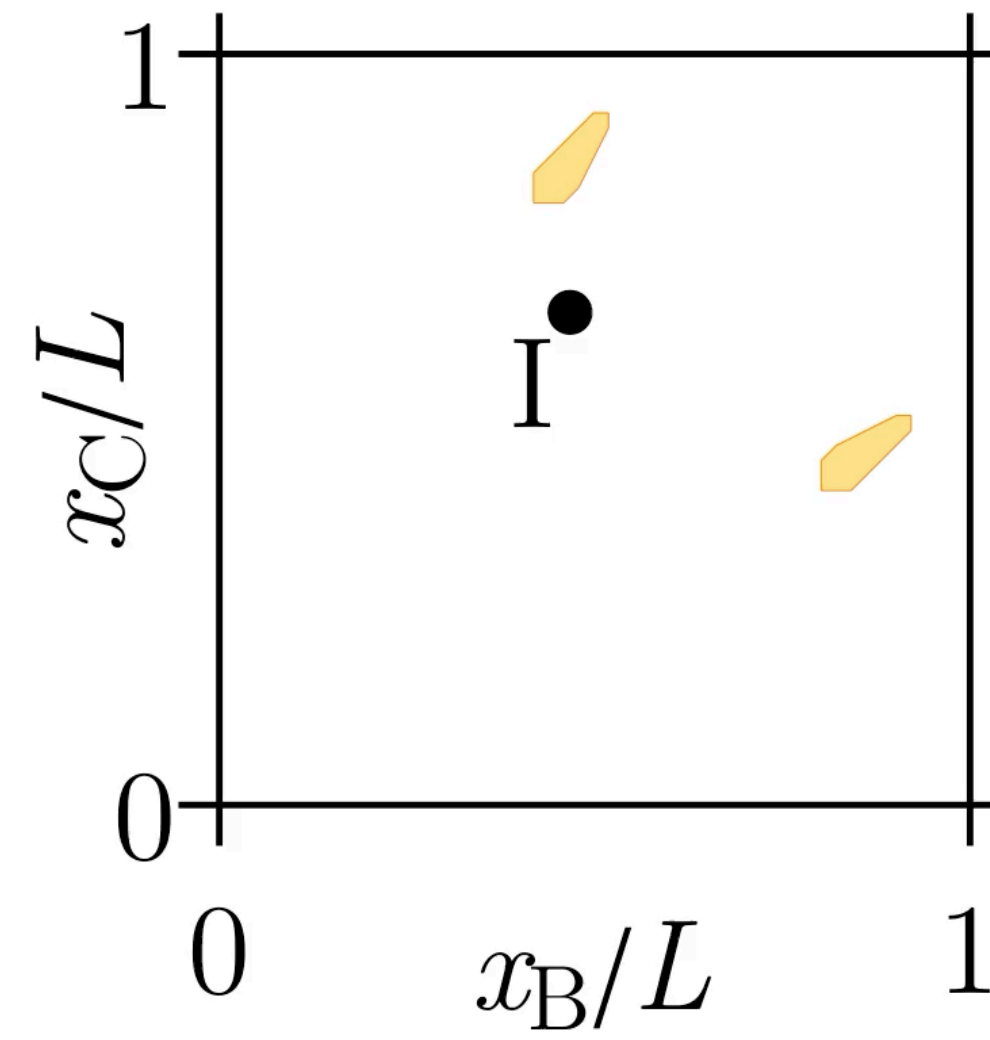
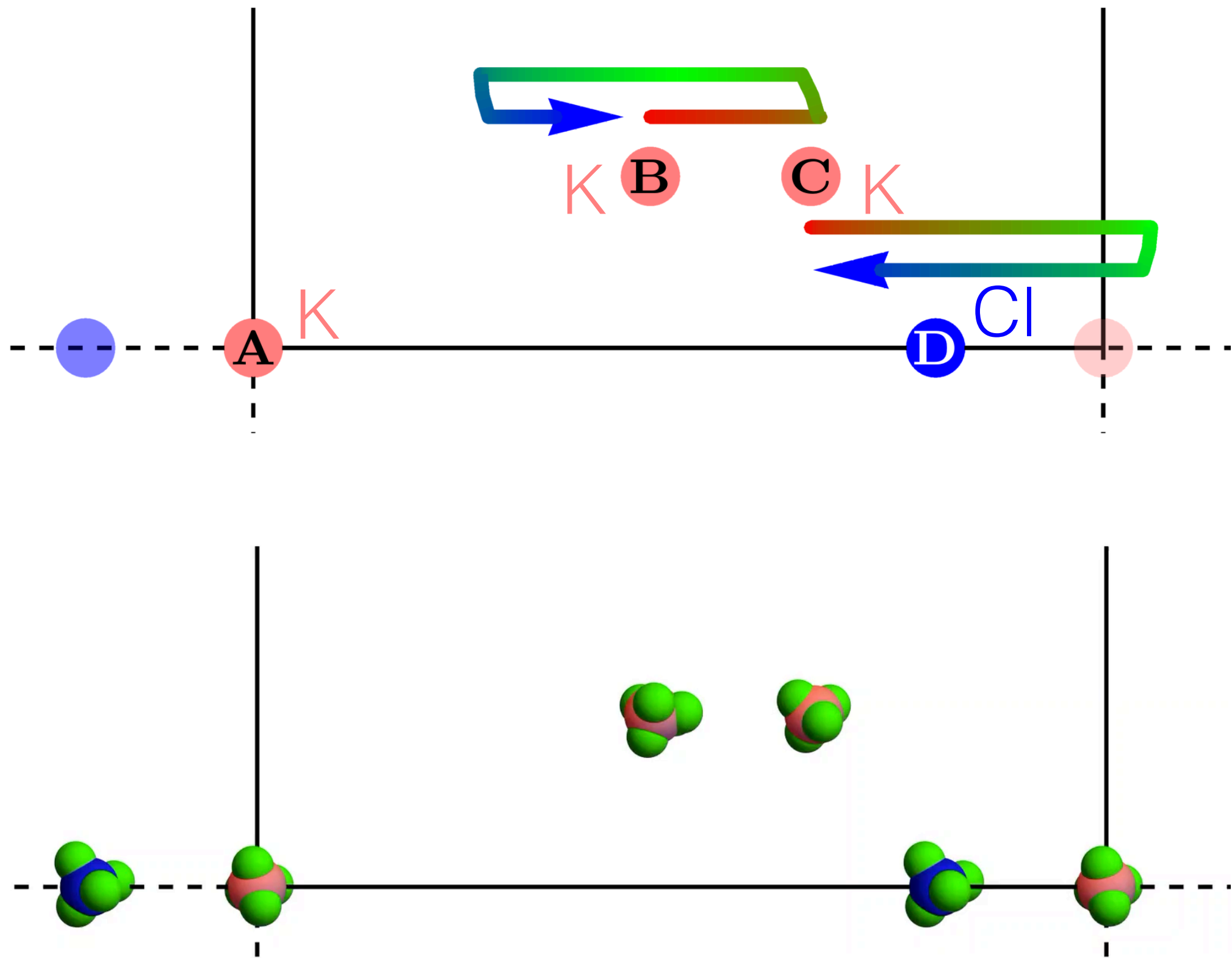
$$x \approx 0.06$$



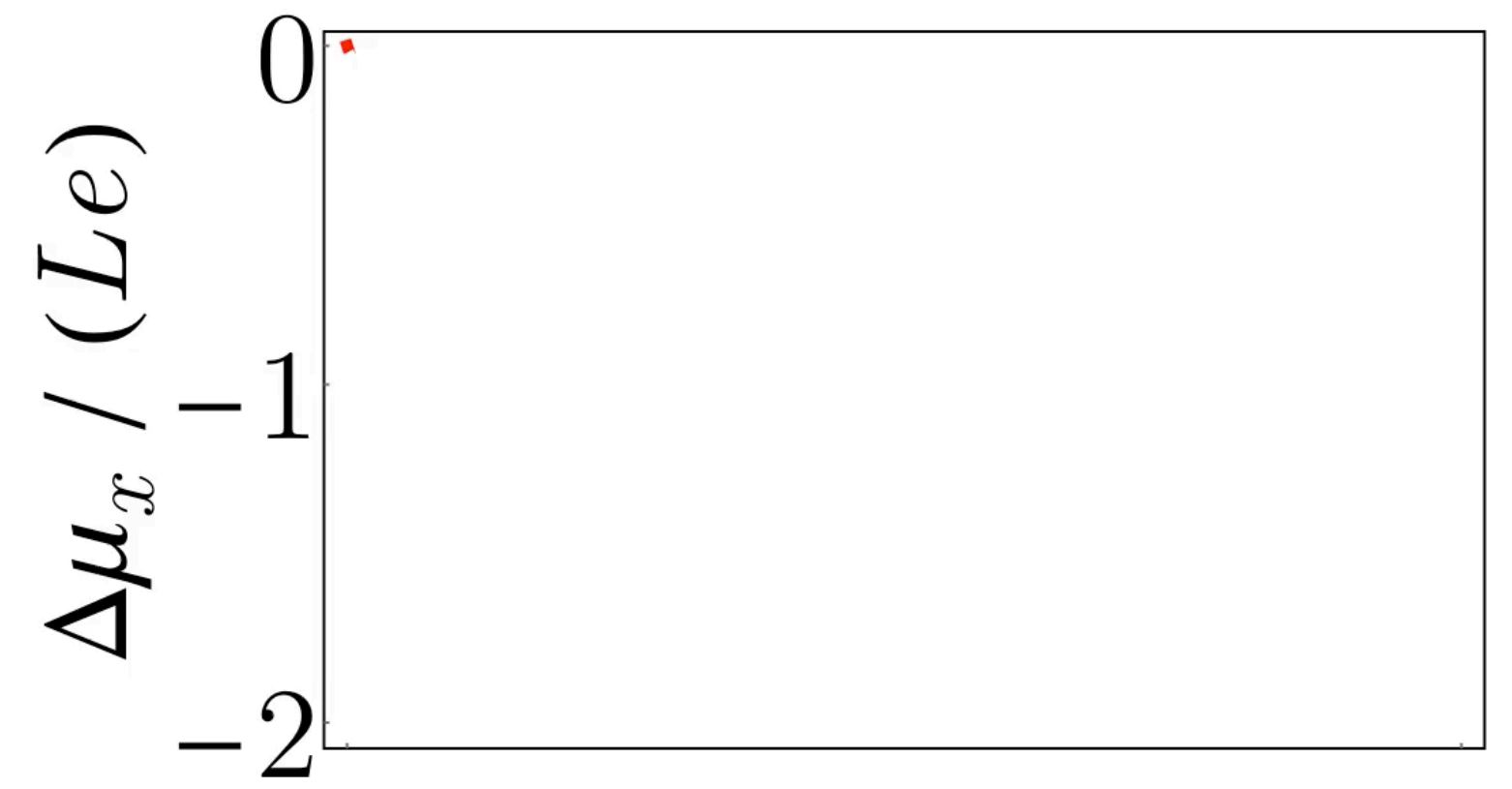
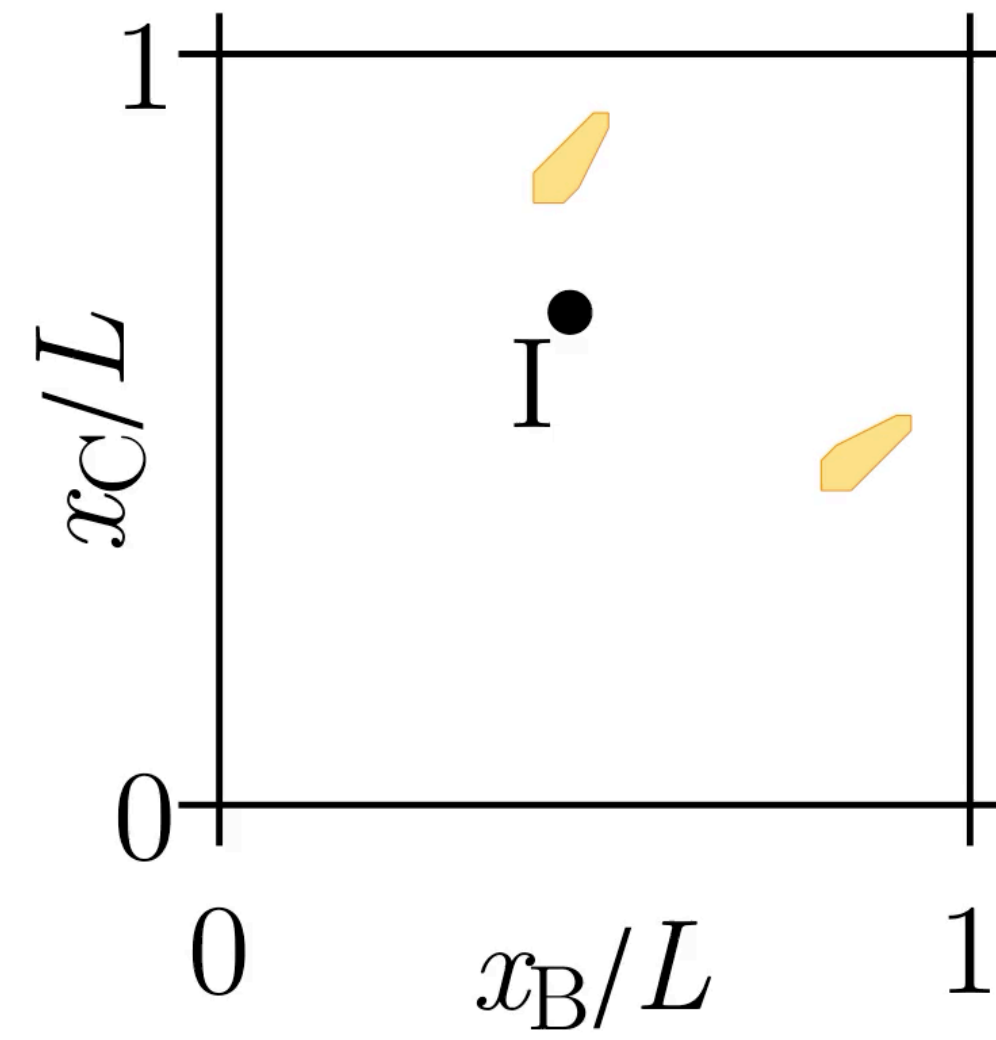
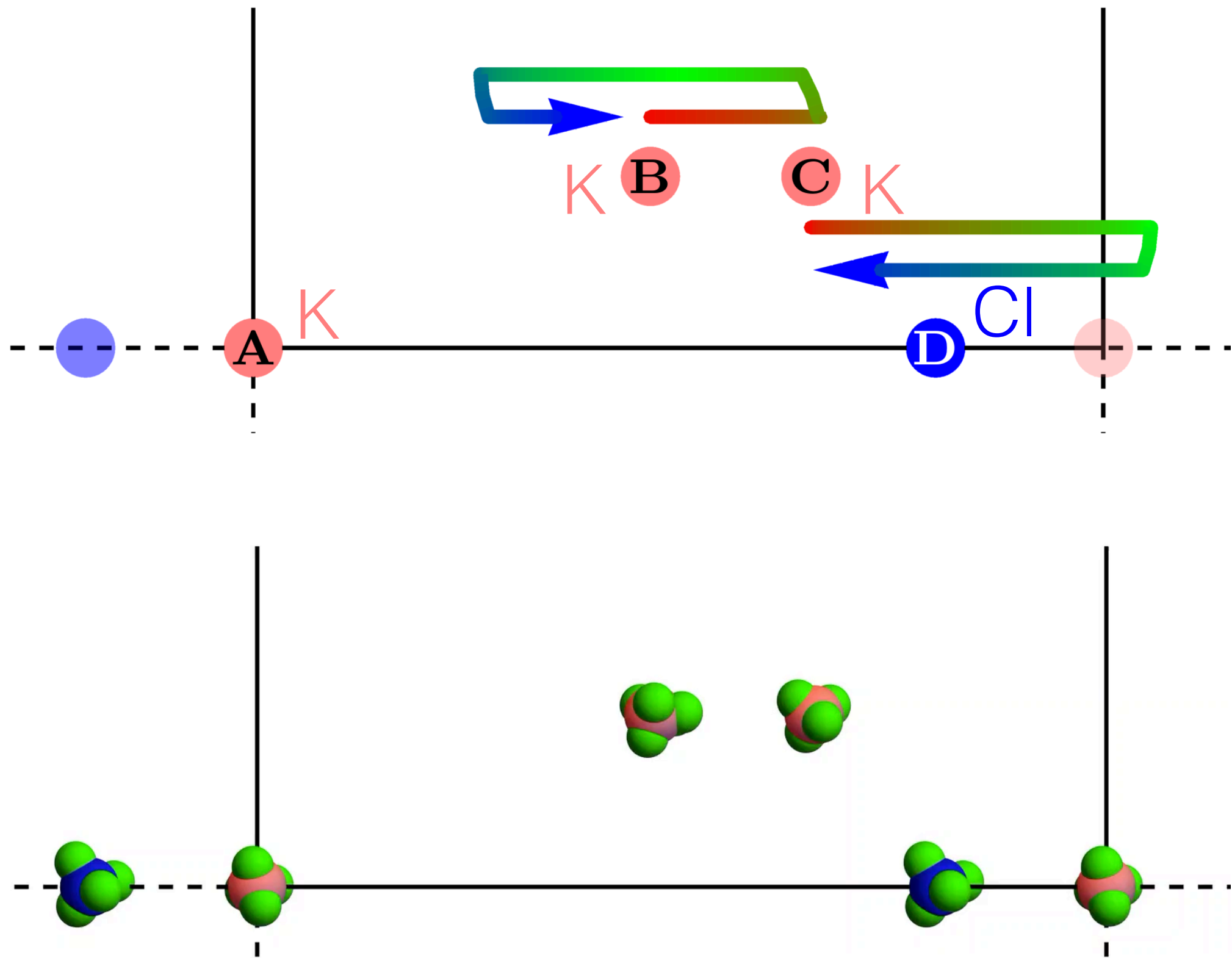
Path
(b)



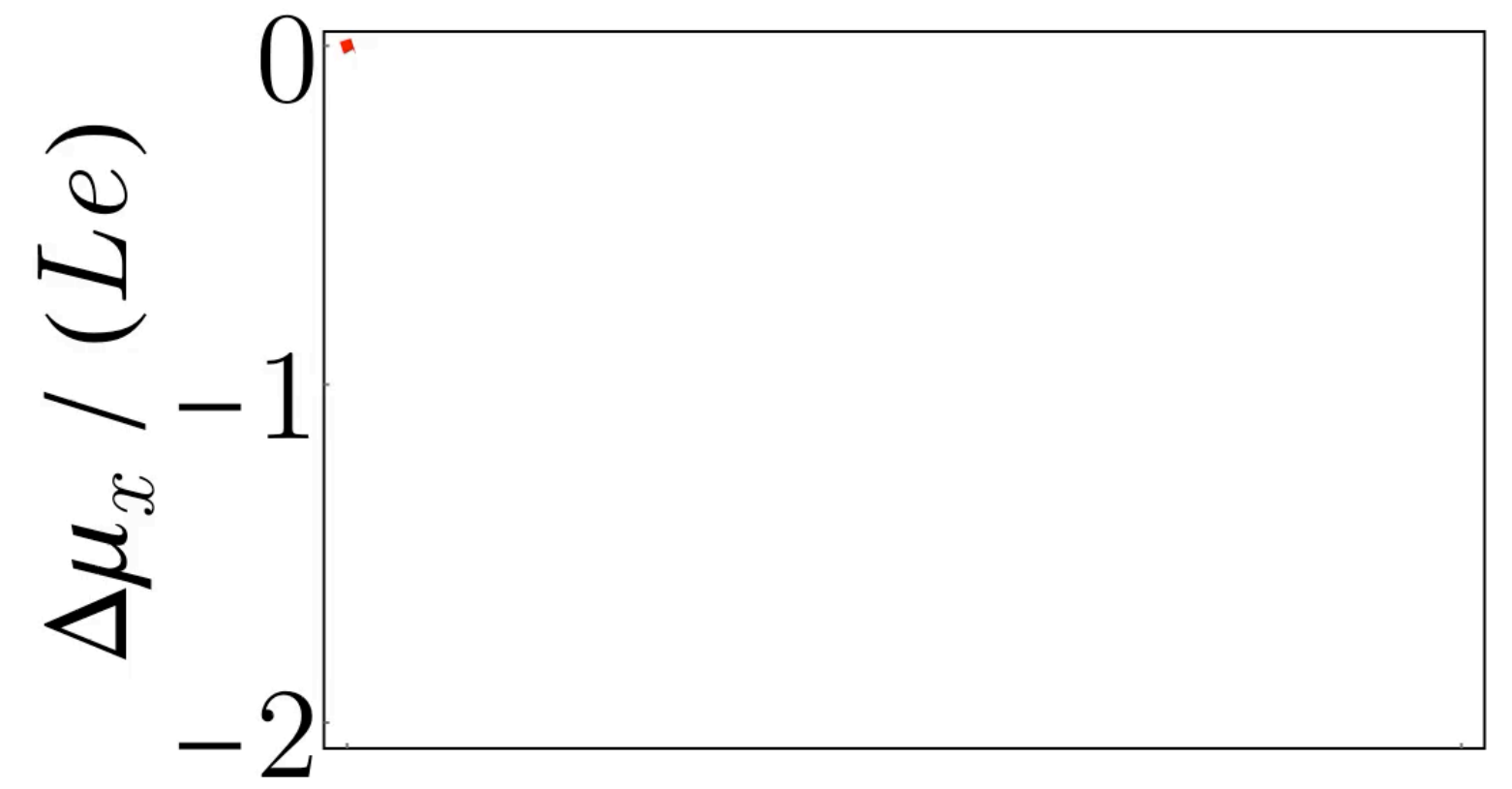
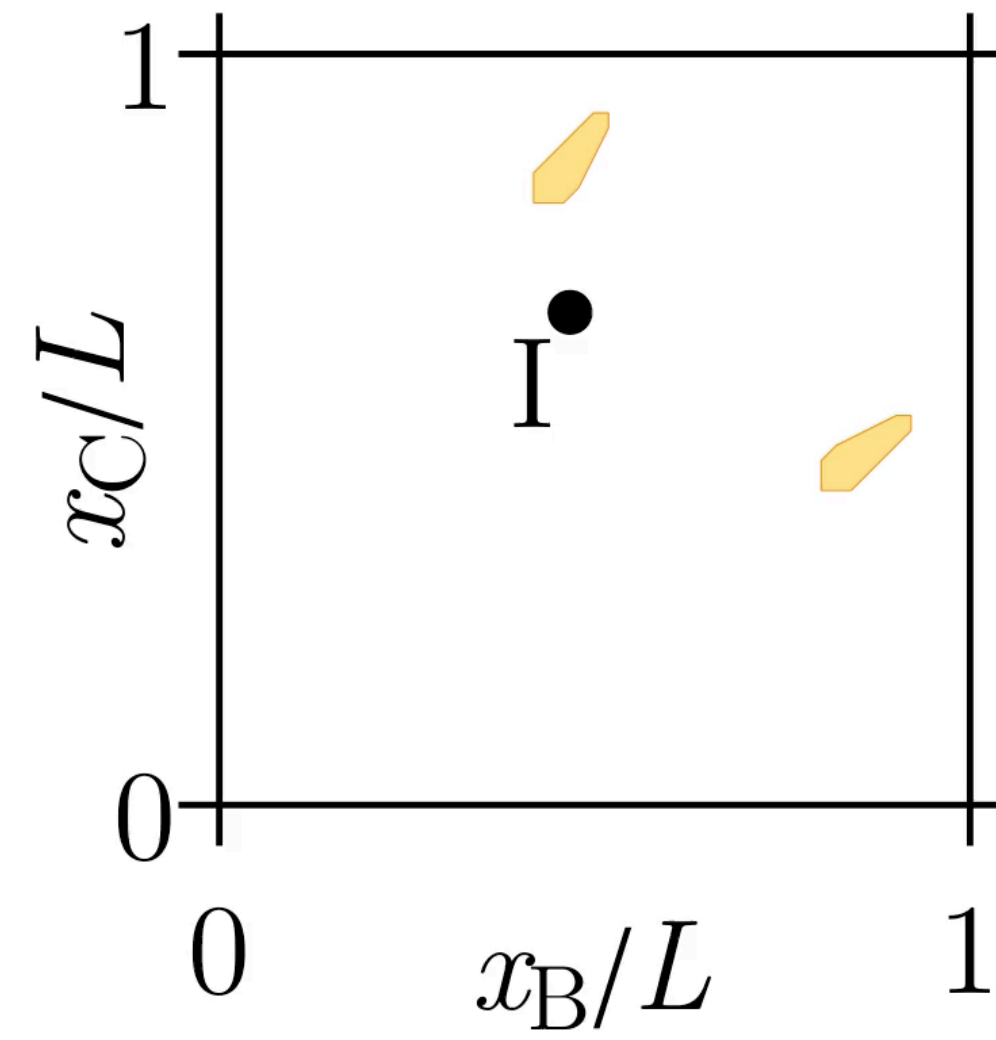
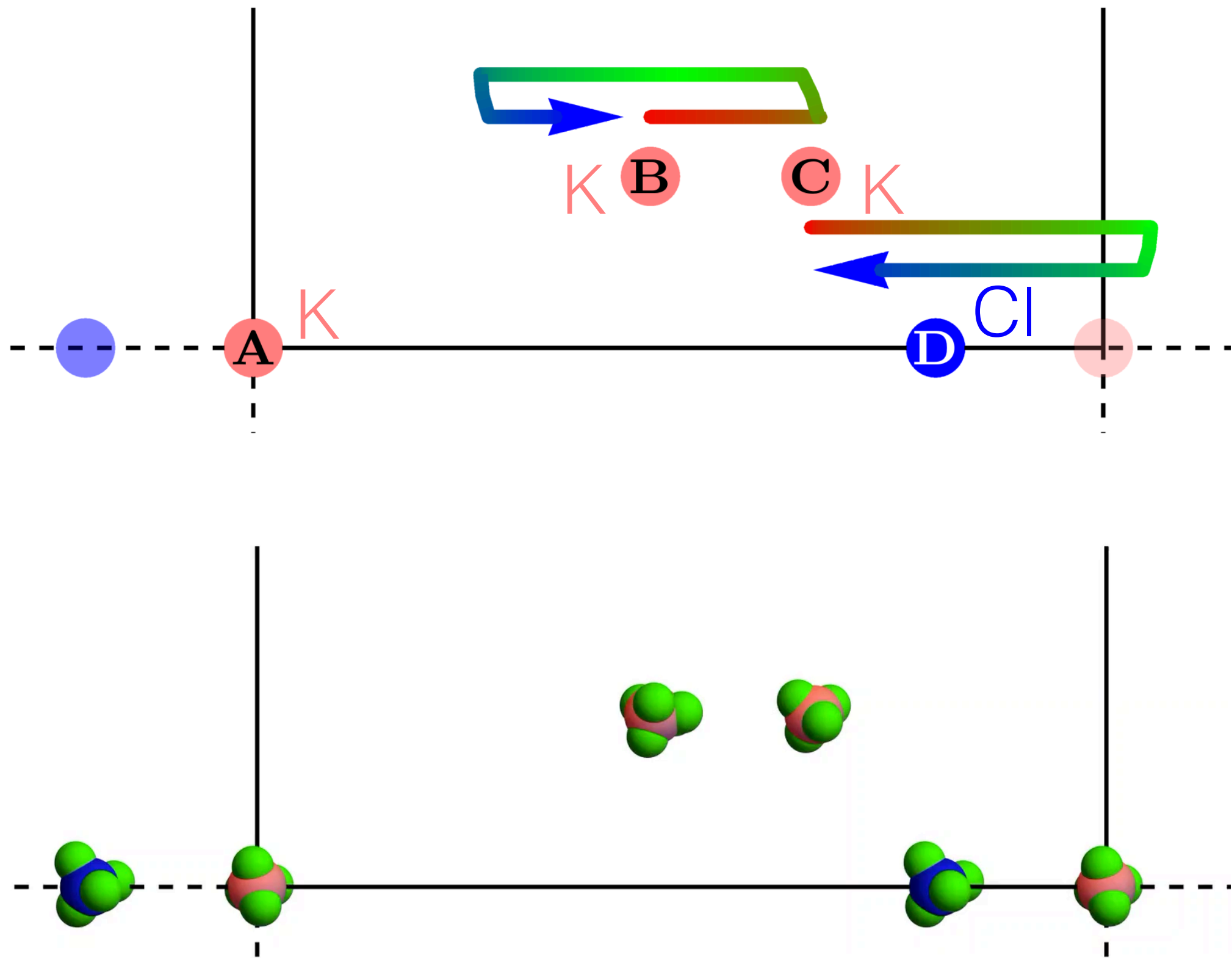
non-trivial particle transport



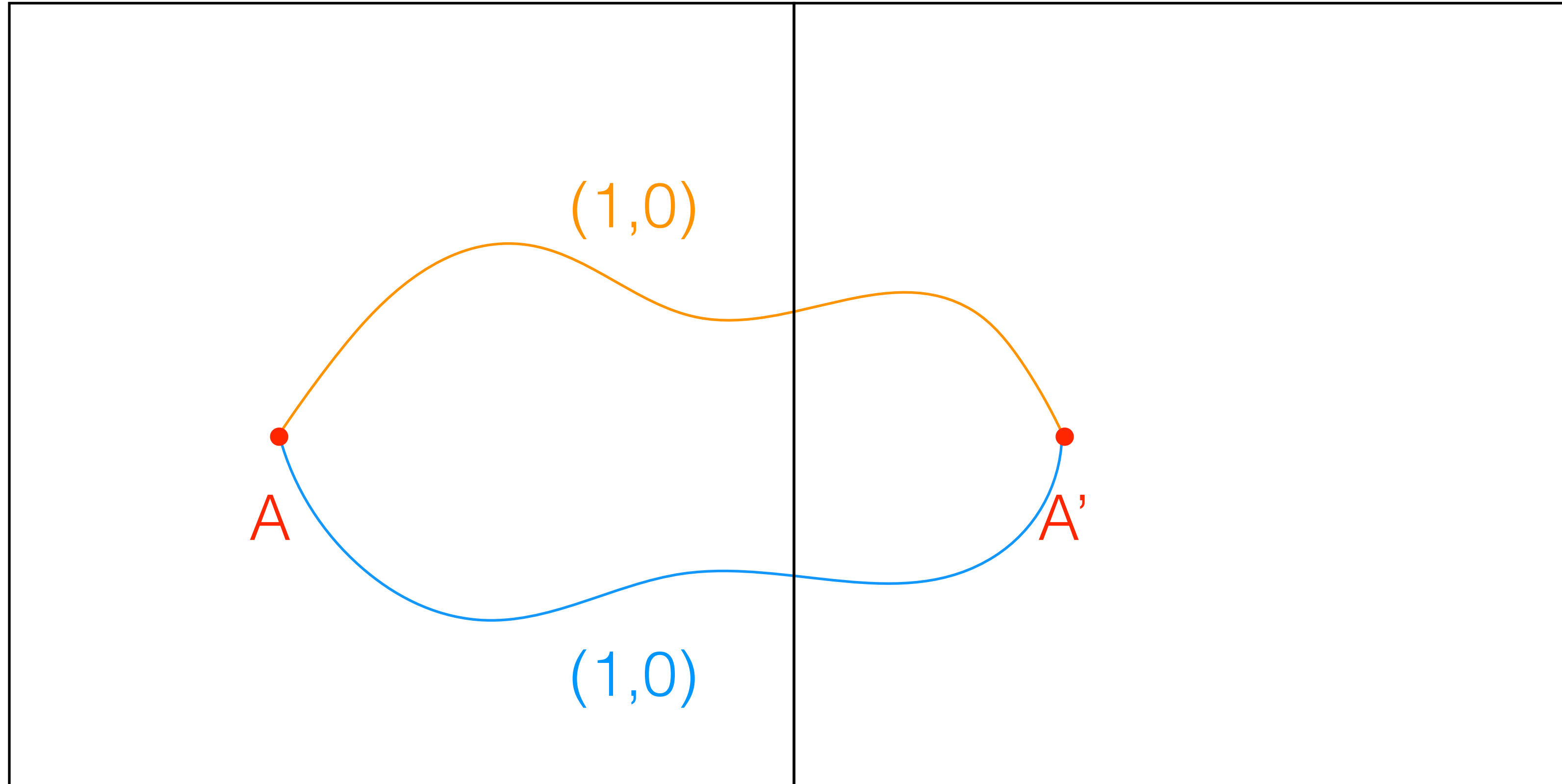
non-trivial particle transport



non-trivial particle transport

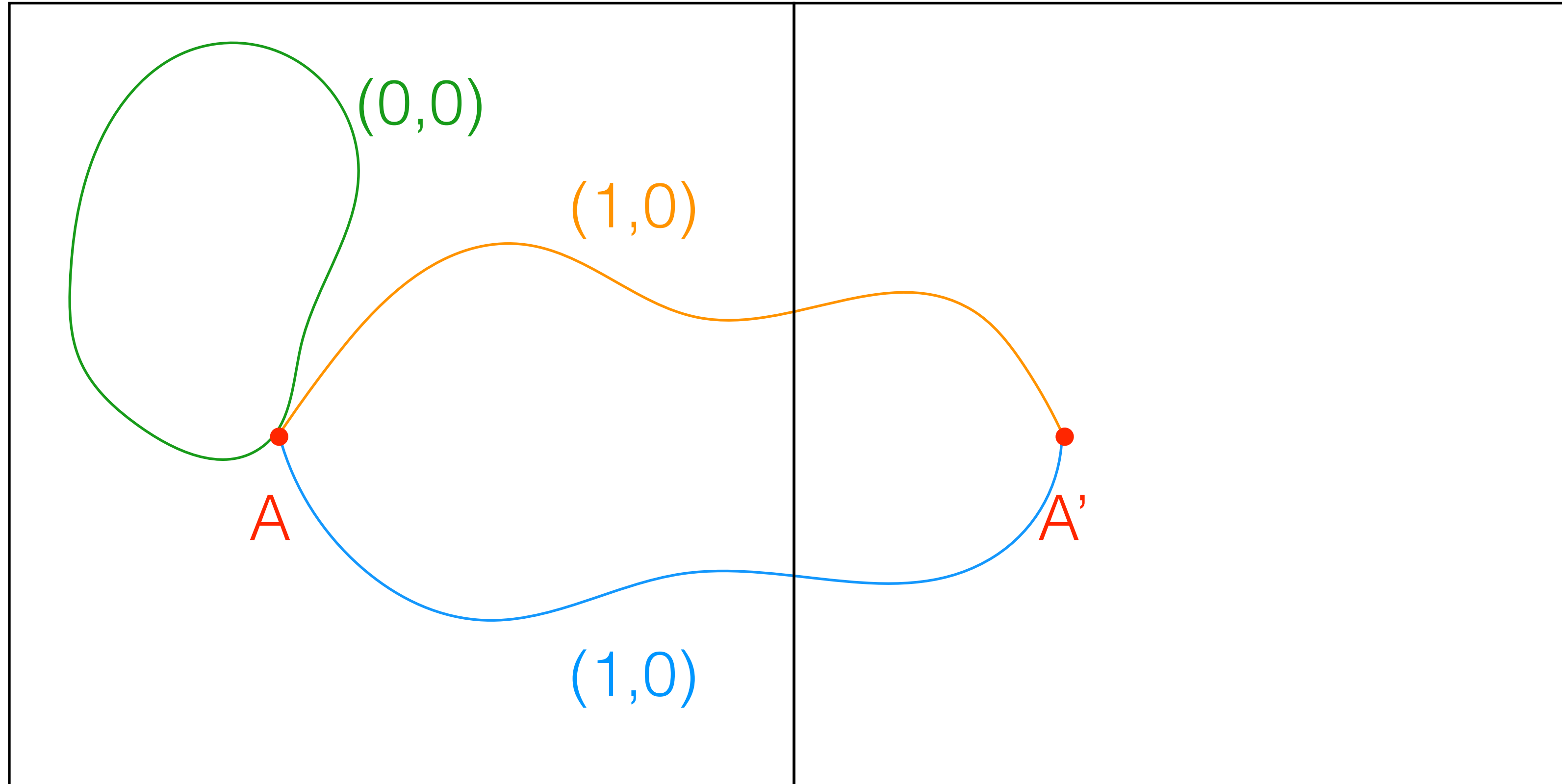


breach of strong adiabaticity



$$\mu = \mu^*$$

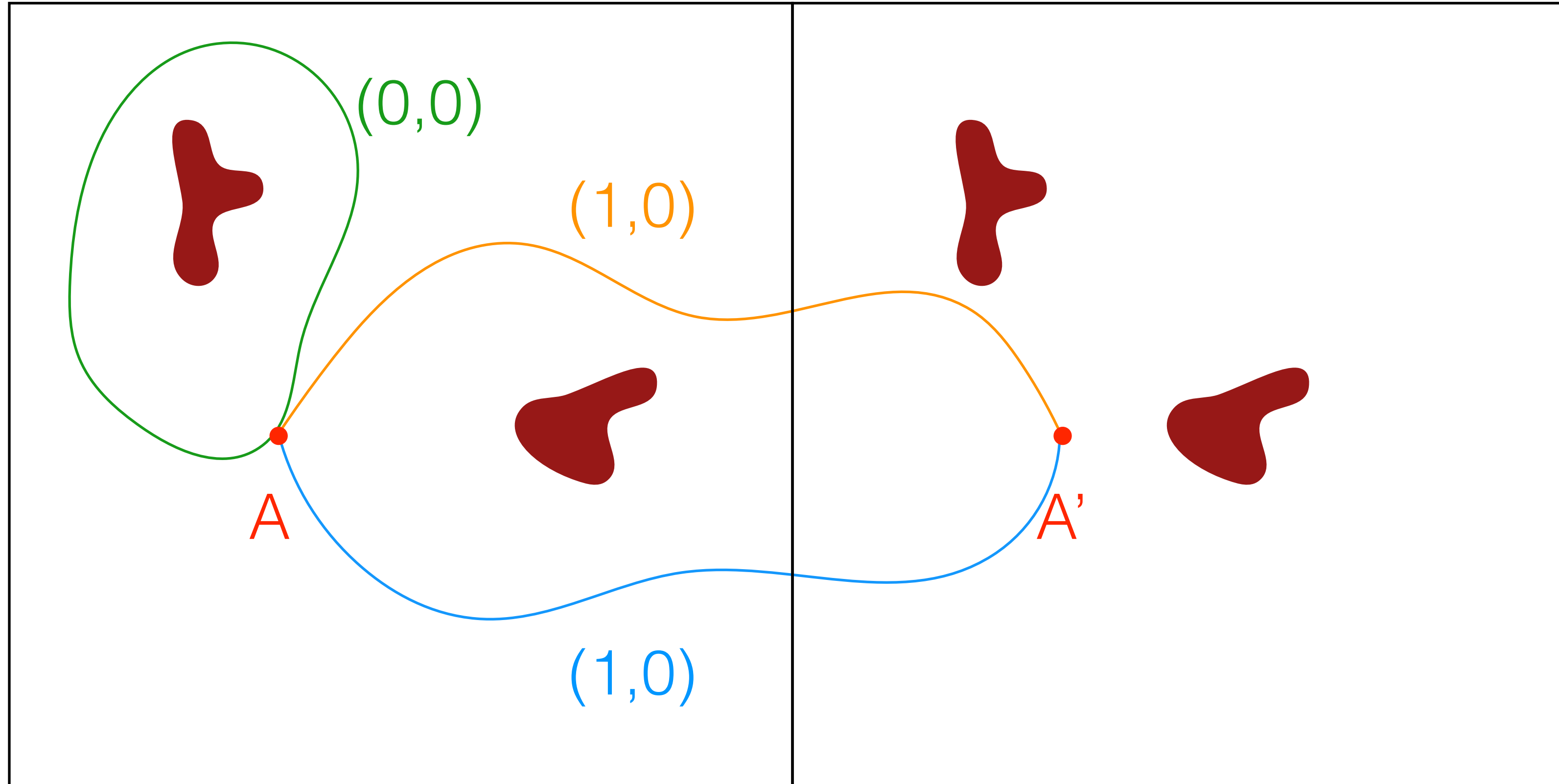
breach of strong adiabaticity



$$\mu = \mu^*$$

$$\mu = 0$$

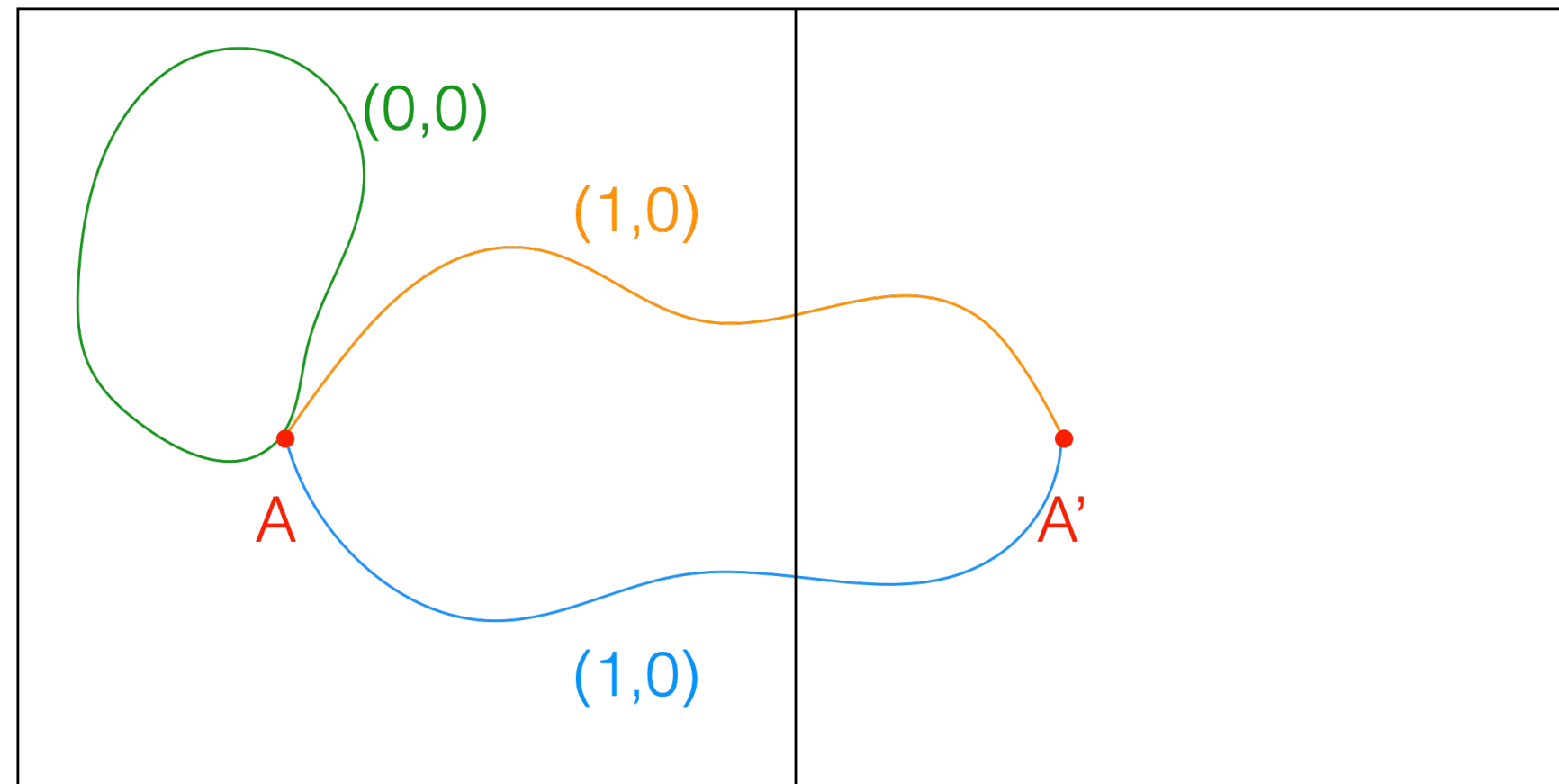
breach of strong adiabaticity



$$\mu \neq \mu^*$$

$$\mu \neq 0$$

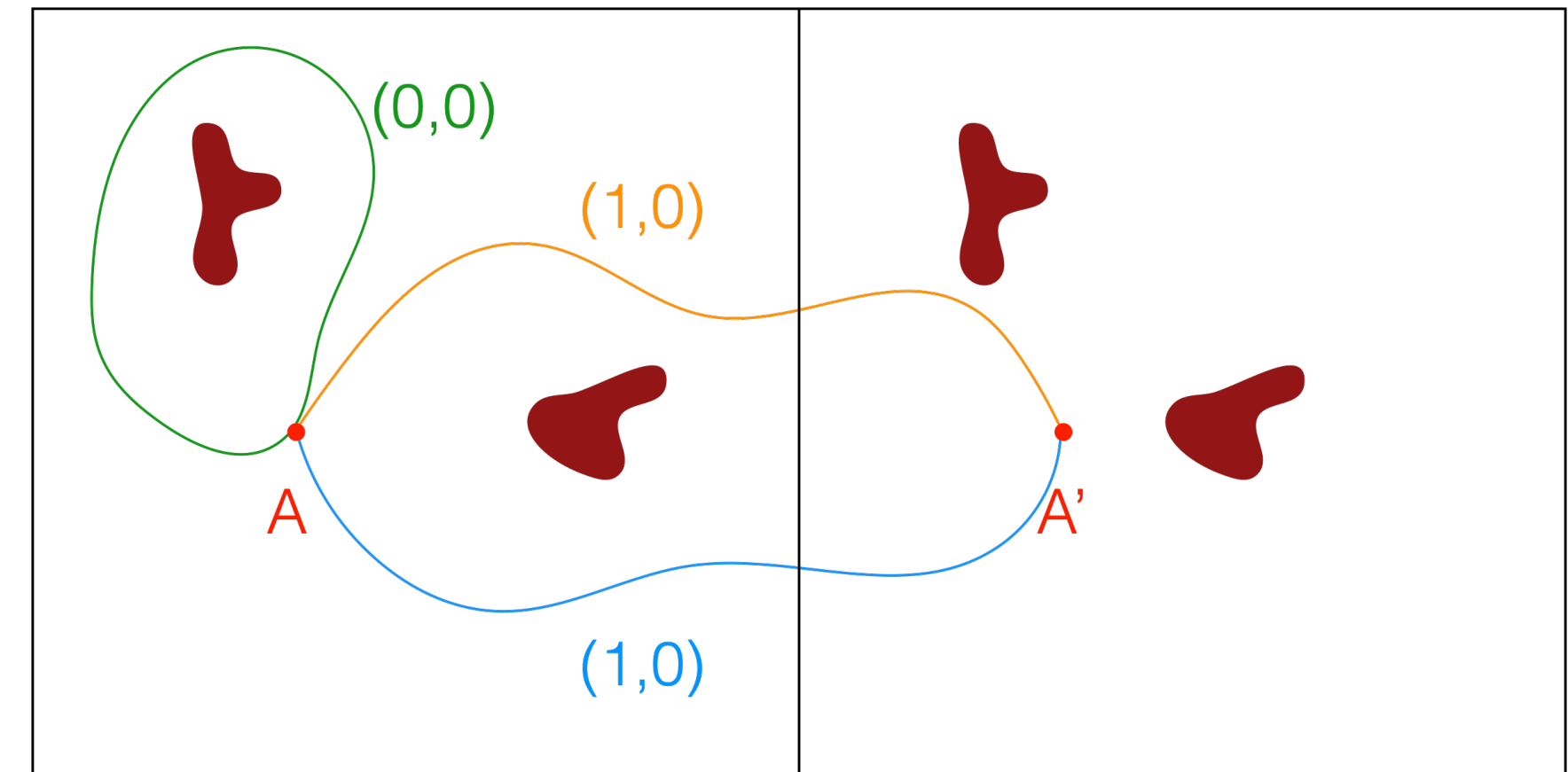
strongly adiabatic transport



$$\begin{aligned} \mu &= \mu^* \\ \mu &= 0 \end{aligned}$$



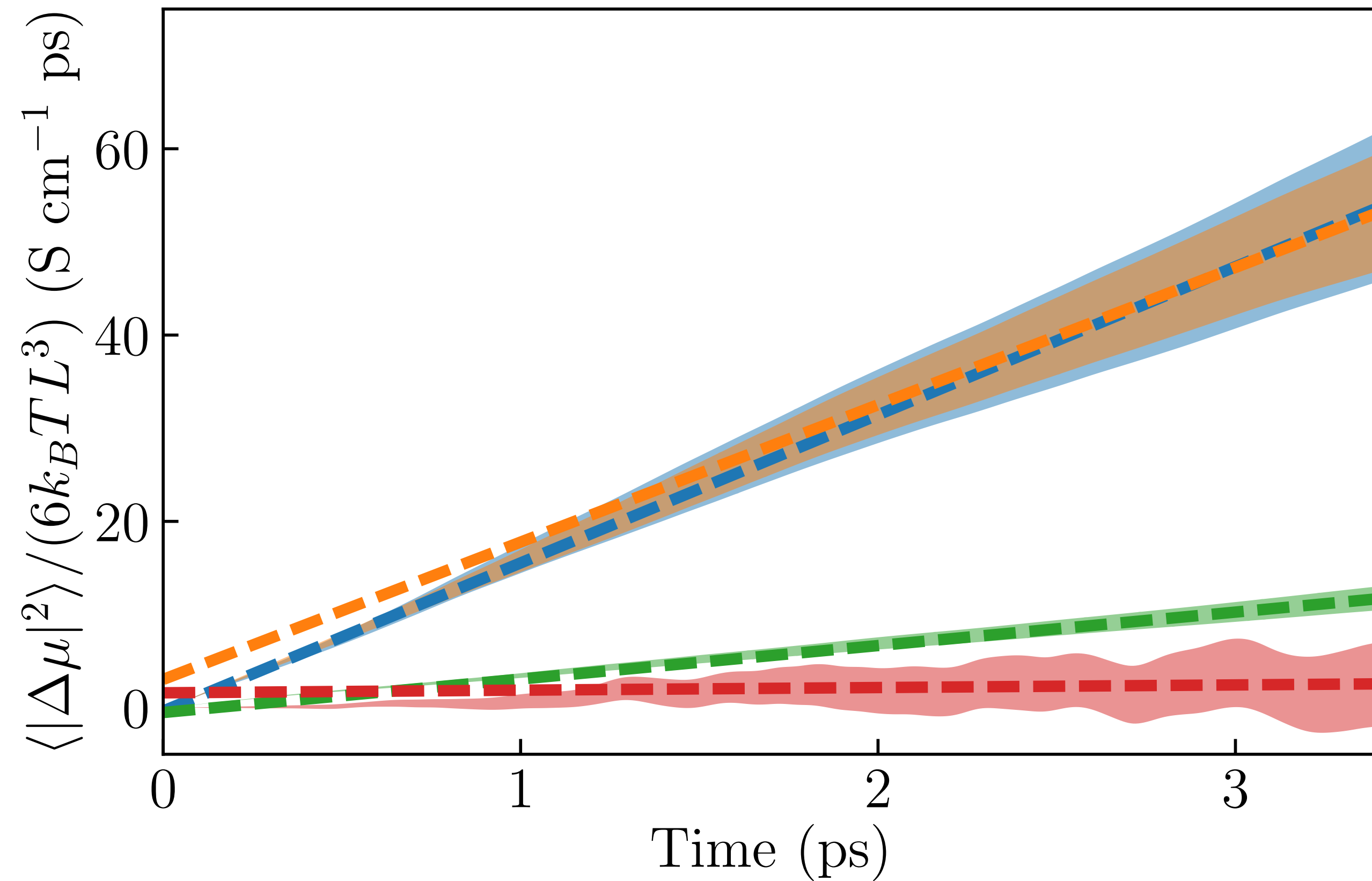
weakly adiabatic transport



$$\begin{aligned} \mu &\neq \mu^* \\ \mu &\neq 0 \end{aligned}$$



not trivial weakly adiabatic conductivity



$$\Delta\boldsymbol{\mu} = e \int_0^t \mathbf{J}(t') dt'$$

$$J_\alpha(t) = \sum_{i\beta} Z_{i\alpha\beta}^*(t) v_{i\beta}(t)$$

$$J_\alpha(t) = \sum_i q_{S(i)} v_{i\alpha}(t) - 2v_\alpha^{lp}(t)$$

cross term



conclusions

- topological quantisation of adiabatic charge transport allows for a rigorous definition of the atomic oxidation states;
- gauge invariance and quantisation of charge transport make the electric conductivity of stoichiometry electrolytes depend on the formal oxidation numbers of the ionic species, via the Green-Kubo formula;
- breach of strong adiabaticity in non-stoichiometric electrolytes triggers an anomalous transport regime, intermediate between metallic and ionic, whereby charge may be transported without any concurrent mass displacement.



thanks to:



Federico Grasselli



Paolo Pegolo

now both @EPFL



Microscopic theory and quantum simulation of atomic heat transport

Aris Marcolongo¹, Paolo Umari² and Stefano Baroni^{1*}

Topological quantization and gauge invariance of charge transport in liquid insulators

Federico Grasselli¹ and Stefano Baroni^{1,2*}

PHYSICAL REVIEW X

Oxidation States, Thouless' Pumps, and Nontrivial Ionic Transport in Nonstoichiometric Electrolytes

Paolo Pegolo, Federico Grasselli, and Stefano Baroni
Phys. Rev. X **10**, 041031 – Published 12 November 2020

Review | [Open Access](#) |

Topology, Oxidation States, and Charge Transport in Ionic Conductors

Paolo Pegolo , Stefano Baroni , Federico Grasselli

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